

# Cosmic Background Neutrino Decay Search – COBAND Experiment –

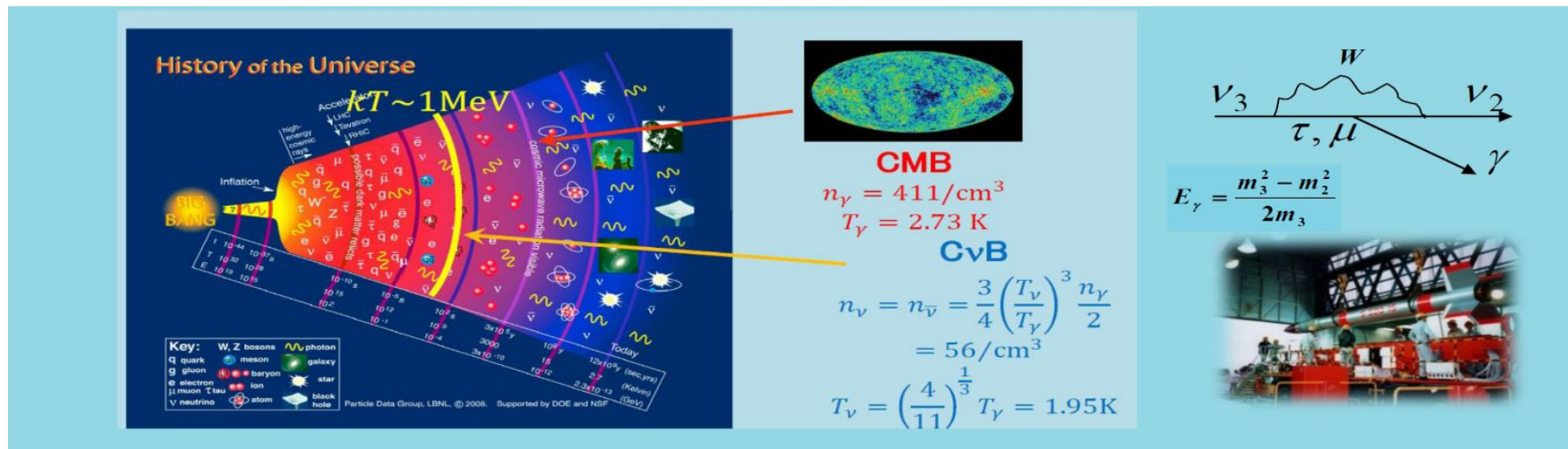
## Continuous Spectral Measurement in Far-Infrared Region using STJ

Shin-Hong Kim (University of Tsukuba, TCHoU)  
for COBAND collaboration

at Center for the Underground Physics (CUP) at the IBS



# COBAND (COsmic Background Neutrino Decay) Collaboration

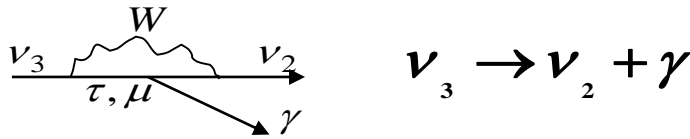


S.H. Kim, Y. Takeuchi, K. Takemasa, K. Nagata, K. Kasahara, S. Yagi, R. Wakasa, R. Senzaki,  
 K. Moriuchi, C. Asano, T. Iida (University of Tsukuba),  
 S. Matsuura (Kwansei Gakuin University),  
 H. Ikeda, T. Wada, K. Nagase, S. Baba (JAXA),  
 Y. Arai, I. Kurachi, M. Hazumi (KEK),  
 T. Yoshida, M. Sakai, T. Nakamura (University of Fukui),  
 Y. Kato (Kindai University),  
 K. Kiuchi, S. Mima (RIKEN),  
 H. Ishino, H. Kibayashi (Okayama University),  
 S. Shiki, G. Fujii, M. Ukibe, M. Ohkubo (AIST),  
 S. Kawahito (Shizuoka University),  
 E. Ramberg, P. Ruvinov, D. Sergatskov (Fermilab),  
 S.B. Kim (Seoul National University)

# Motivation of Search for Cosmic Background Neutrino Decay

- Why three generations ? Why such mass hierarchy with large mass differences ?  
Only neutrino masses are not measured . To determine the neutrino mass itself is an important subject.

Neutrino decay observation can determine the neutrino mass.

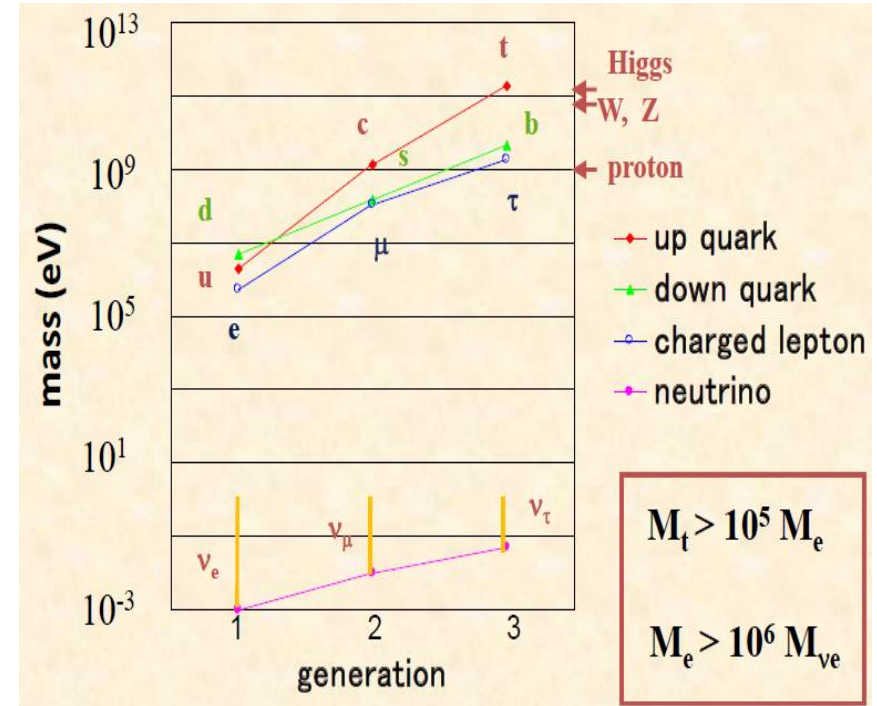


$$E_\gamma = \frac{m_3^2 - m_2^2}{2m_3} = \frac{\Delta m_{23}^2}{2m_3}$$

Using  $\Delta m_{23}^2 = (2.43 \pm 0.09) \times 10^{-3} \text{ eV}^2$

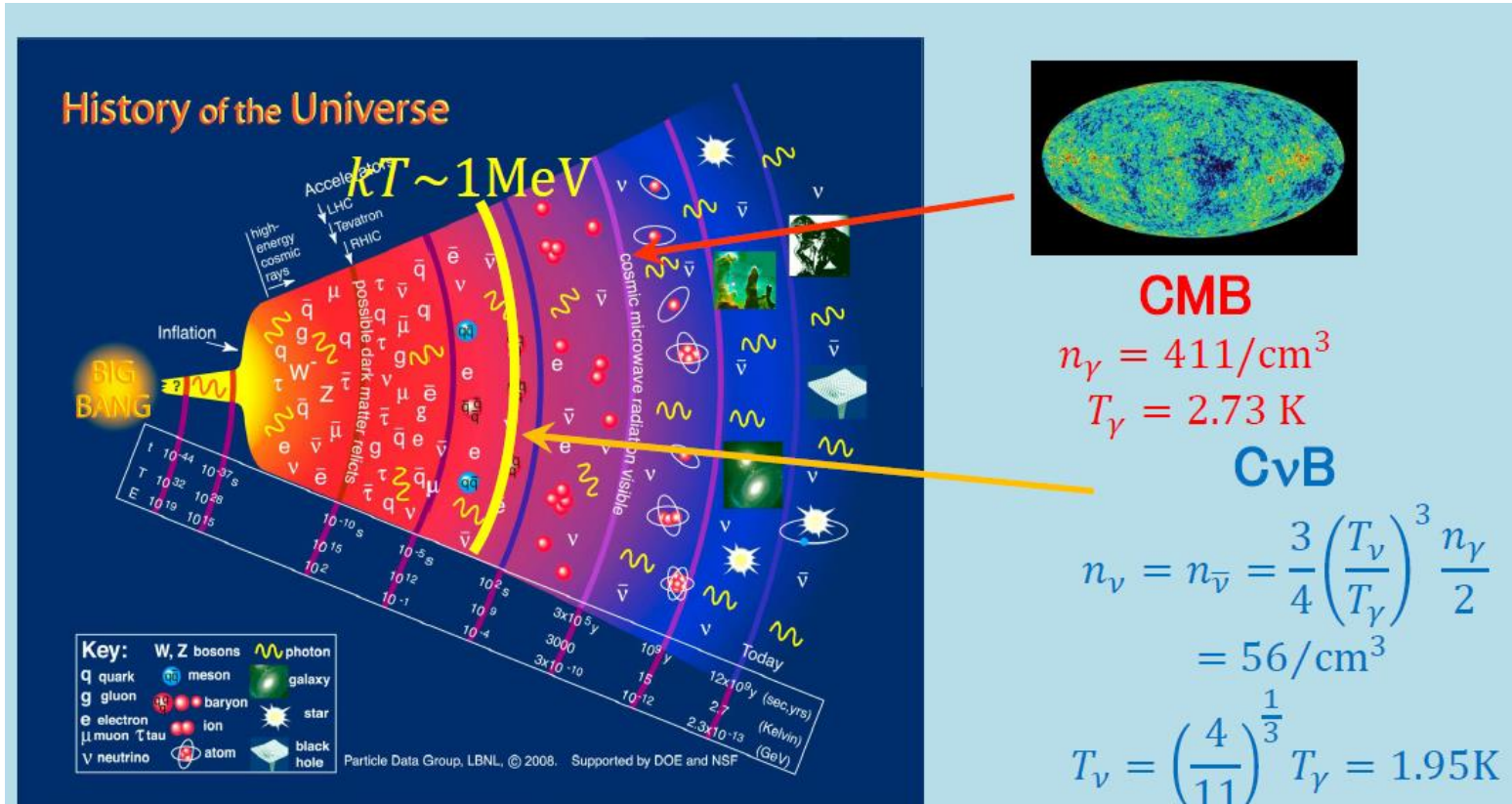
$E_\gamma = 10 \sim 25 \text{ meV}$  at  $\nu_3$  rest frame.

(Far - Infrared region  $\lambda = 50 \sim 125 \mu$ )



- As the neutrino lifetime is very long, we need use cosmic background neutrino to observe the neutrino decay. To observe this decay of the cosmic background neutrino means a discovery of the cosmic background neutrino predicted by cosmology.
- Left-Right symmetric model predicts the neutrino lifetime larger than  $10^{17}$  year while the standard model predicts  $2 \times 10^{43}$  year.  
Measured neutrino lifetime limit  $\tau > 3 \times 10^{12}$  year.

# Big-Bang Cosmology and Cosmic Background Neutrino (CvB)

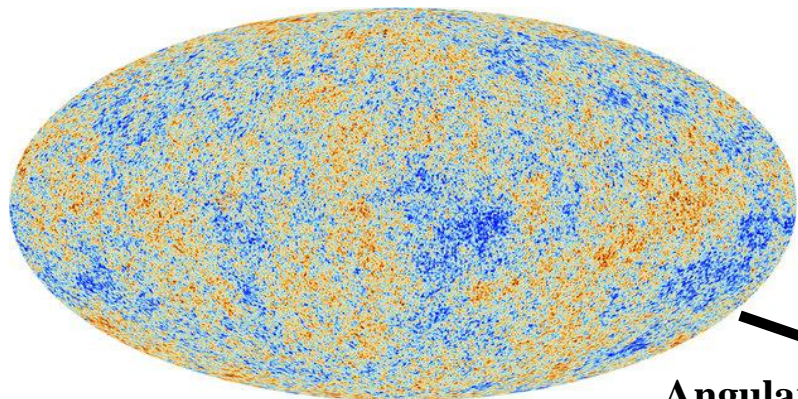


- A few seconds after Big Bang → Cosmic Background Neutrino (CvB) became free.
- 300,000 years after Big Bang → Cosmic Microwave Background (CMB) became free.



# Large Scale Structure of the Universe and CvB

Temperature fluctuation of CMB  
observed by Planck satellite



Angular  
correlation  
at high  
temperature

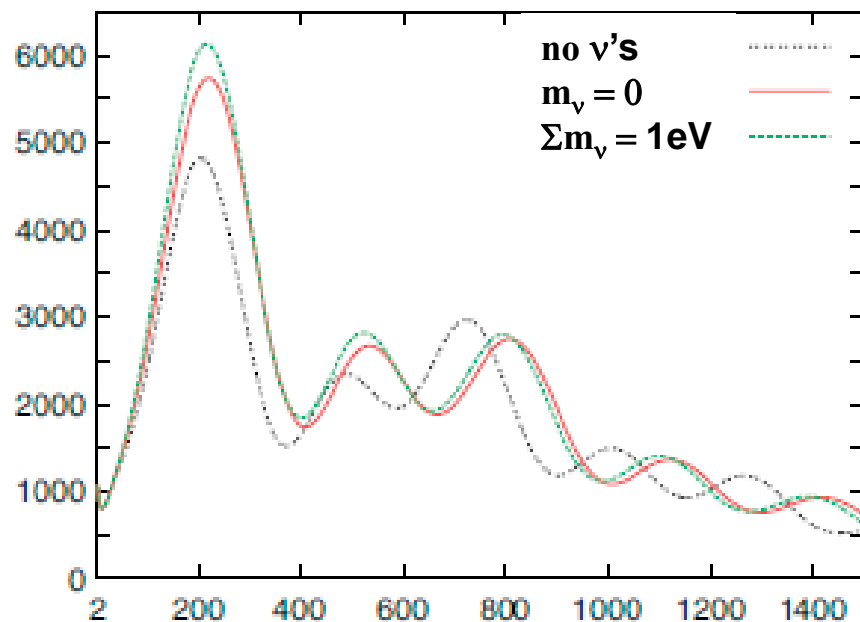
Decision Factor of Temperature  
Fluctuation:

Dark matter, Dark energy, **Neutrino**



Matter density  
or Large scale structure of the Universe

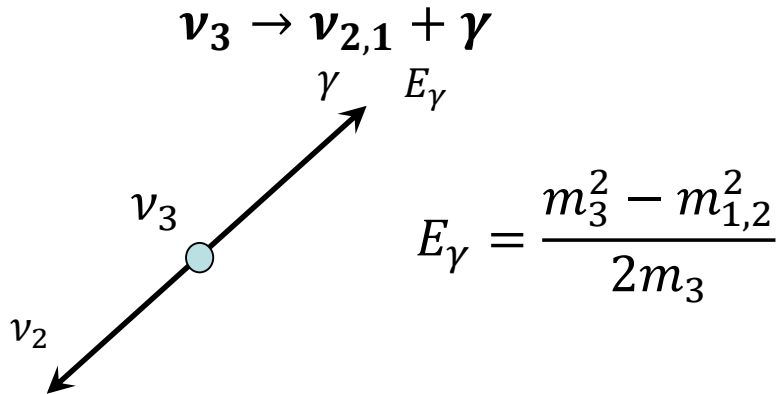
Theoretical prediction on Power  
Spectra of the angular correlation of  
the CMB temperature fluctuation  
(depending on the neutrino mass)



Large Angular scale ( $180^\circ / l$ ) Small  
High Temperature = High Galaxy matter density  
CvB exists → large angular correlation at high angle  
higher neutrino mass → larger angular correlation at high angle

# Photon Energy from Neutrino Decay

## Neutrino decay

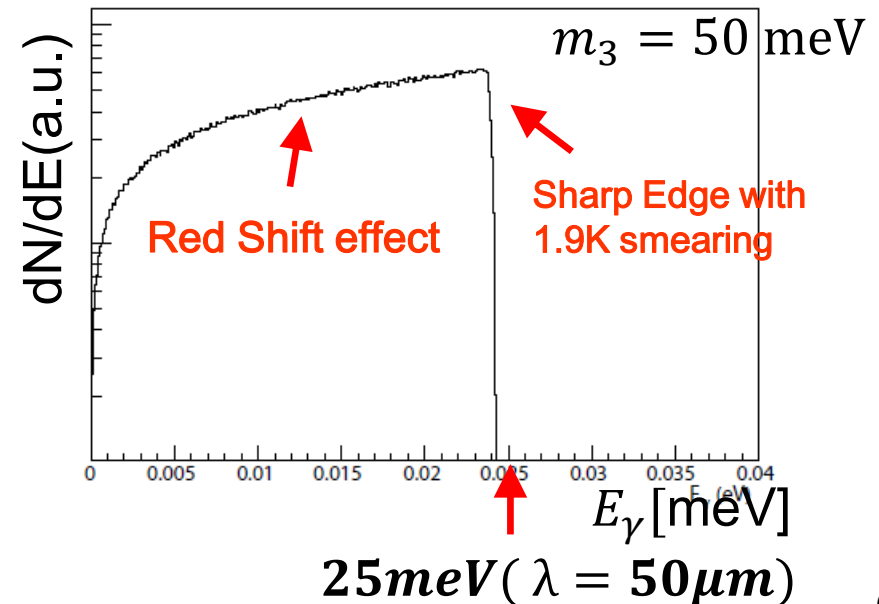
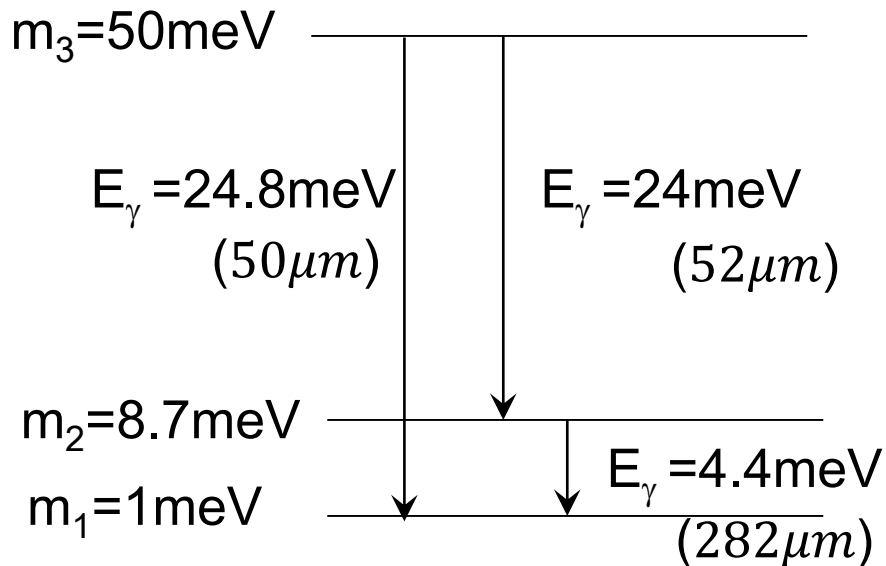


- Neutrino oscillation results
  - $|\Delta m_{23}^2| = |m_3^2 - m_2^2| \sim 2.4 \times 10^{-3} \text{ eV}^2$
  - $\Delta m_{12}^2 \sim 7.65 \times 10^{-5} \text{ eV}^2$
- Space observatory results  
( Planck+WP+highL+BAO )
  - $\Sigma m_i < 0.23 \text{ eV}$

→  $50 \text{ meV} < m_3 < 87 \text{ meV}$

$E_\gamma = 14 \sim 24 \text{ meV}$  ( $\lambda_\gamma = 51 \sim 89 \mu\text{m}$ )

## Photon energy distribution $\nu_3 \rightarrow \nu_2 + \gamma$



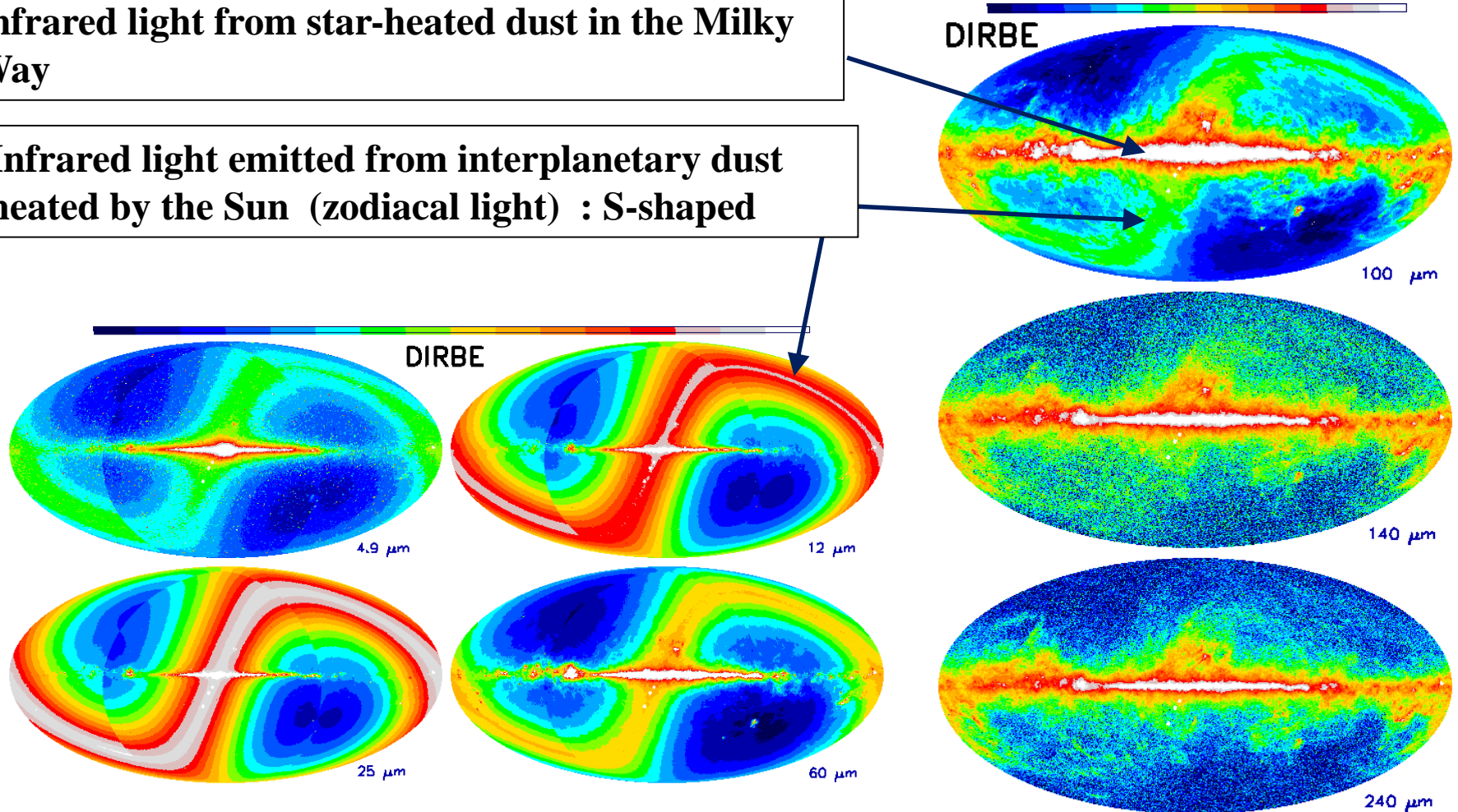
# **Cosmic Infrared Background Radiation**

# Cosmic Infrared Background Observation by COBE

**January 1998 (Press Release):** COBE (Cosmic Background Explorer) announced the first observation of the cosmic infrared background (CIB)

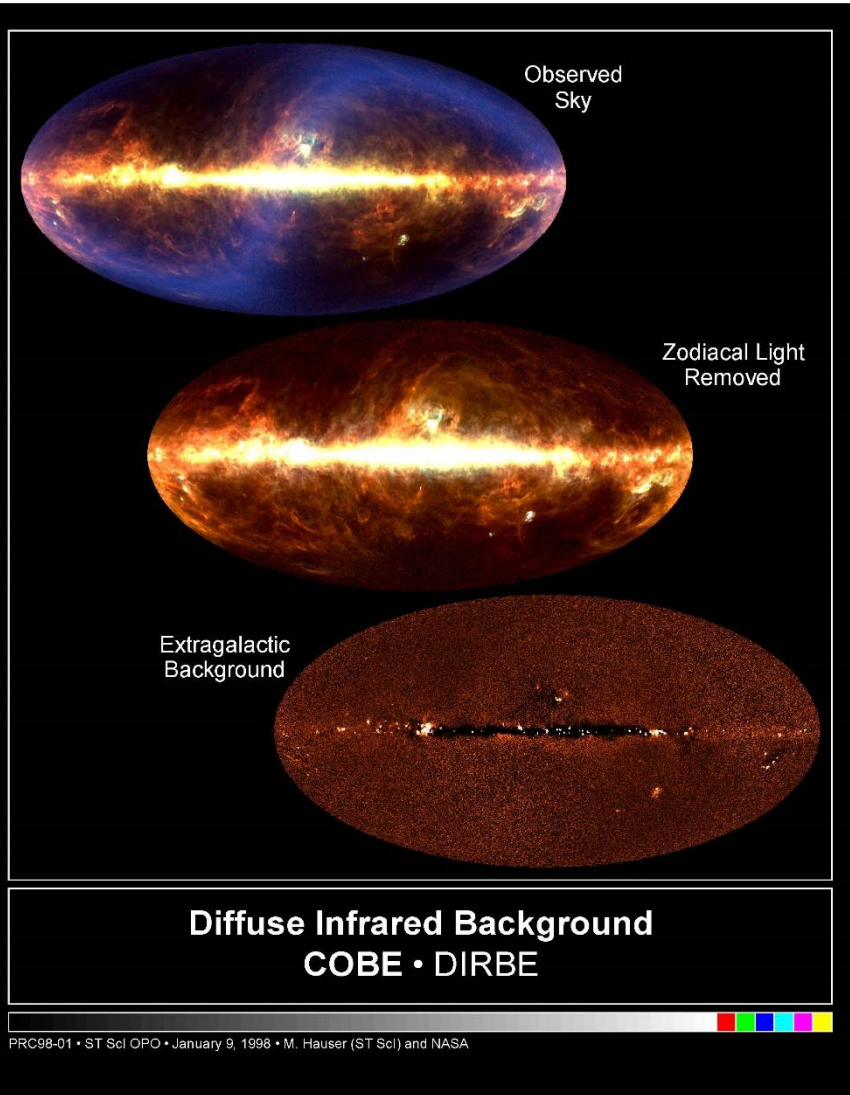
**Infrared light from star-heated dust in the Milky Way**

**Infrared light emitted from interplanetary dust heated by the Sun (zodiacal light) : S-shaped**





# Cosmic Infrared Background Observation by COBE



**(Top)**

**Observed Sky at  $\lambda = 60 \mu\text{m}$ (blue),  $100 \mu\text{m}$ (green) and  $240 \mu\text{m}$ (red).**

**Horizontal band corresponds to Infrared light from interstellar gas in the plane of our Milky Way Galaxy. The blue S-shaped region corresponds to Infrared light from interplanetary dust in the solar system called zodiacal emission.**

**(Middle)**

**After removing the zodiacal emission.**

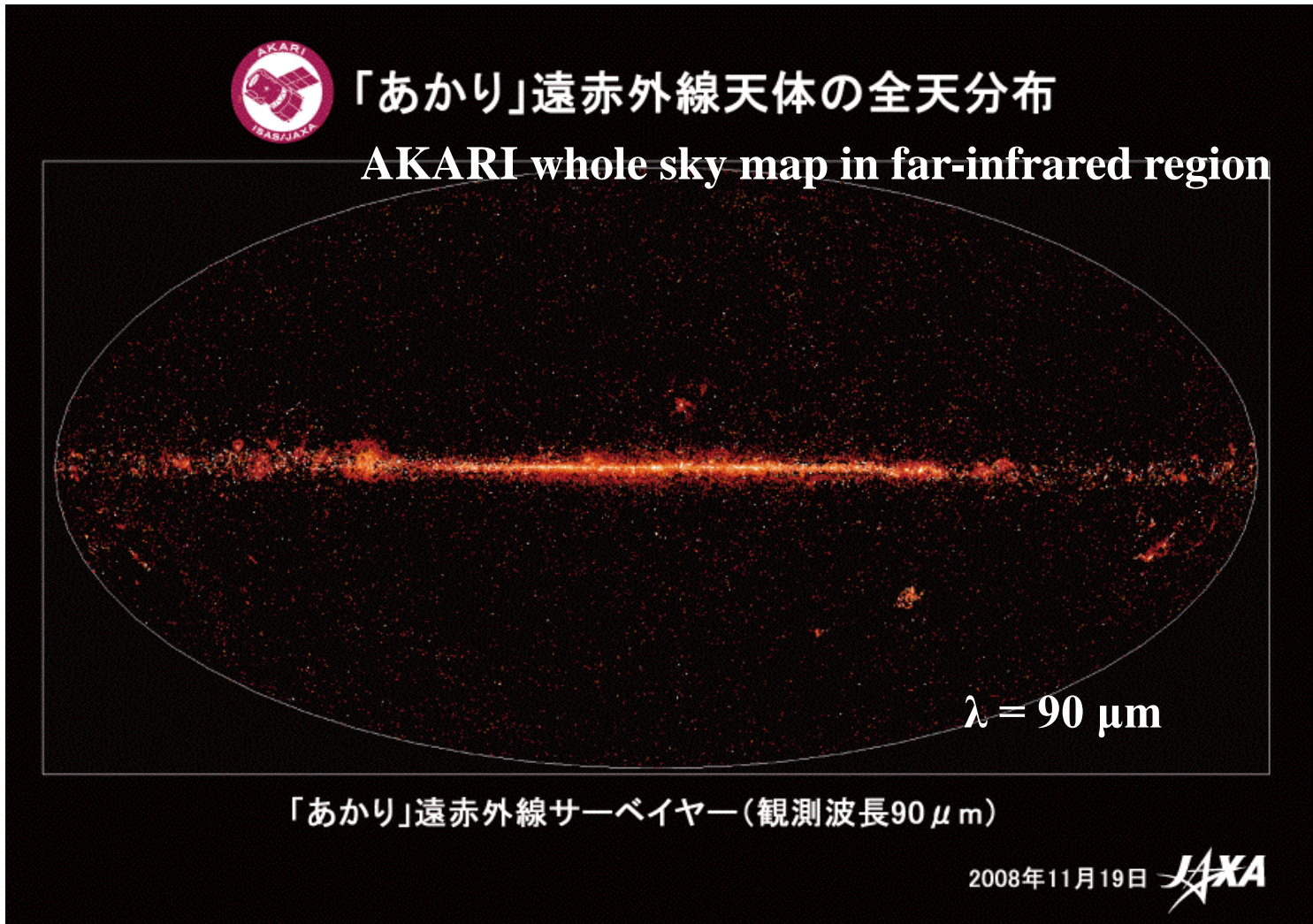
**(Bottom)**

**After removing the infrared light from our solar system and Galaxy at  $\lambda = 240 \mu\text{m}$ .**

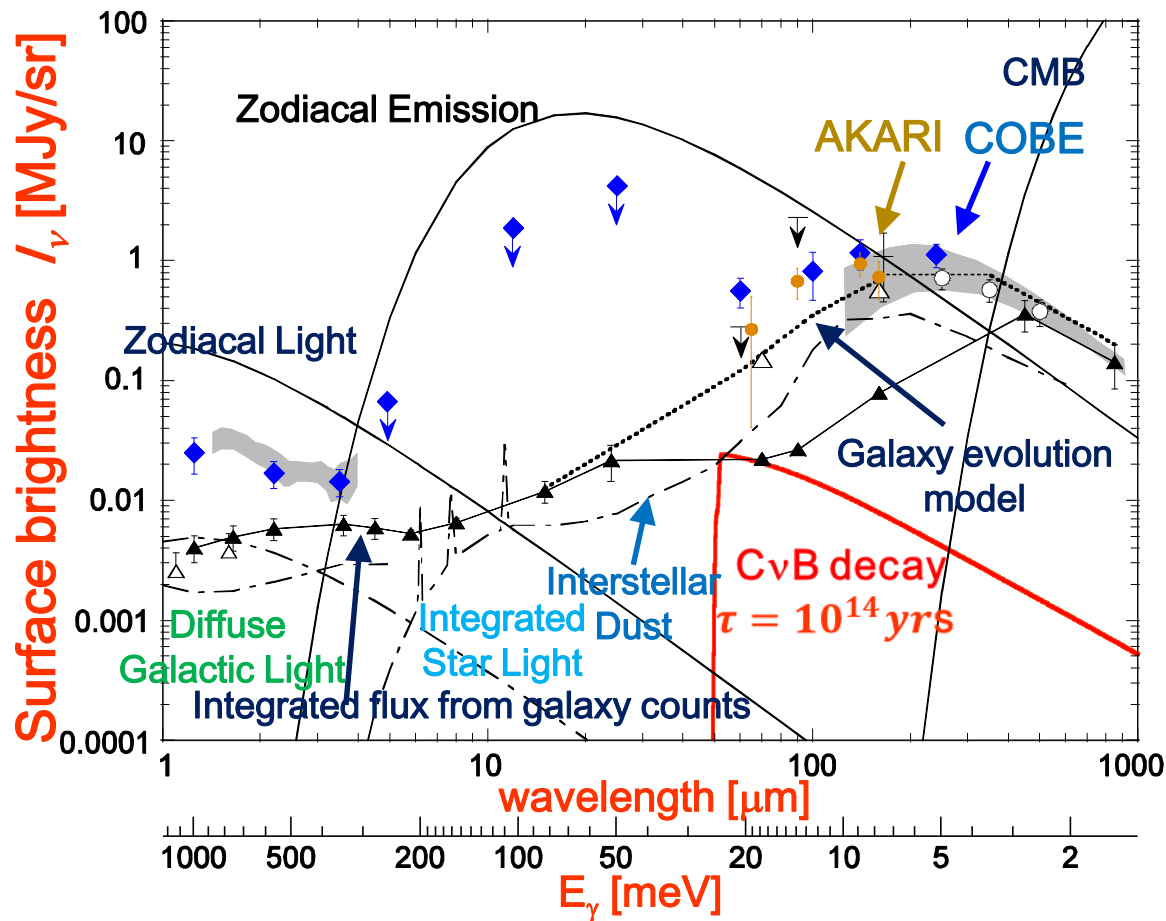
**COBE announced that they observed the extragalactic residual infrared light called “Cosmic Infrared Background” at the 240 and 140 micrometer wavelength.**

# Cosmic Far-Infrared Background Observation by AKARI

November 2008 : AKARI ( JAXA Infrared Imaging Satellite) observed 64,000 stars and galaxies at wavelengths of 65, 90, 140 and 160  $\mu\text{m}$ . Made a whole sky map. They extracted CIB map from this observation.



# Photon Energy Spectrum from Outer Space



CIB  
measurements  
(● AKARI,  
◆ COBE)

There was an excess of the CIB measured by COBE and AKARI over the prediction by galaxy evolution model.

Sources at FIR: dusts of far galaxies or far blackhole, or CvB decay ?



# Our paper published in JPSJ on Jan. 18<sup>th</sup>, 2012

Journal of the Physical Society of Japan **81** (2012) 024101

FULL PAPERS

DOI: [10.1143/JPSJ.81.024101](https://doi.org/10.1143/JPSJ.81.024101)

## Search for Radiative Decays of Cosmic Background Neutrino using Cosmic Infrared Background Energy Spectrum

Shin-Hong KIM\*, Ken-ichi TAKEMASA, Yuji TAKEUCHI, and Shuji MATSUURA<sup>1</sup>

*Graduate School of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Ibaraki 305-8571, Japan*

<sup>1</sup>*Institute of Space and Astronautical Science, JAXA, Sagami-hara 252-5210, Japan*

(Received September 8, 2011; revised November 22, 2011; accepted December 1, 2011; published online January 18, 2012)

We propose to search for the neutrino radiative decay by fitting a photon energy spectrum of the cosmic infrared background to a sum of the photon energy spectrum from the neutrino radiative decay and a continuum. By comparing the present cosmic infrared background energy spectrum observed by AKARI and Spitzer to the photon energy spectrum expected from neutrino radiative decay with a maximum likelihood method, we obtained a lifetime lower limit of  $3.1 \times 10^{12}$  to  $3.8 \times 10^{12}$  years at 95% confidence level for the third generation neutrino  $\nu_3$  in the  $\nu_3$  mass range between 50 and 150 meV/ $c^2$  under the present constraints by the neutrino oscillation measurements. In the left-right symmetric model, the minimum lifetime of  $\nu_3$  is predicted to be  $1.5 \times 10^{17}$  years for  $m_3$  of 50 meV/ $c^2$ . We studied the feasibility of the observation of the neutrino radiative decay with a lifetime of  $1.5 \times 10^{17}$  years, by measuring a continuous energy spectrum of the cosmic infrared background.

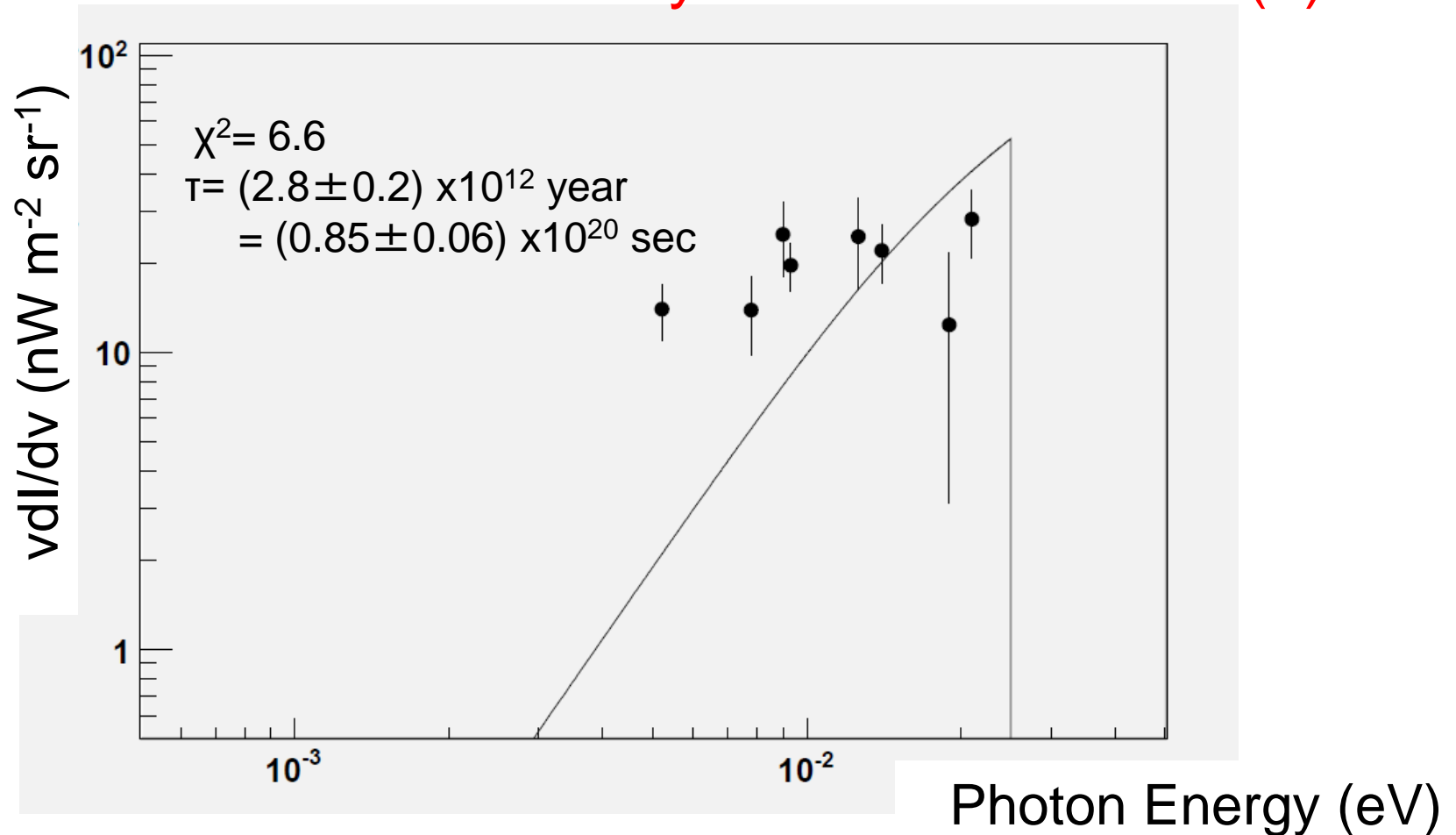
KEYWORDS: neutrino radiative decay, neutrino mass, cosmic background neutrino, cosmic infrared background, COBE, AKARI, Spitzer

Search Region:  $\lambda = 35 \sim 250 \mu\text{m}$  ( $E_\gamma = 35 \sim 5 \text{ meV}$ )

In Rocket experiment,  $\lambda = 40 \sim 80 \mu\text{m}$  ( $E_\gamma = 31 \sim 15 \text{ meV}$ )

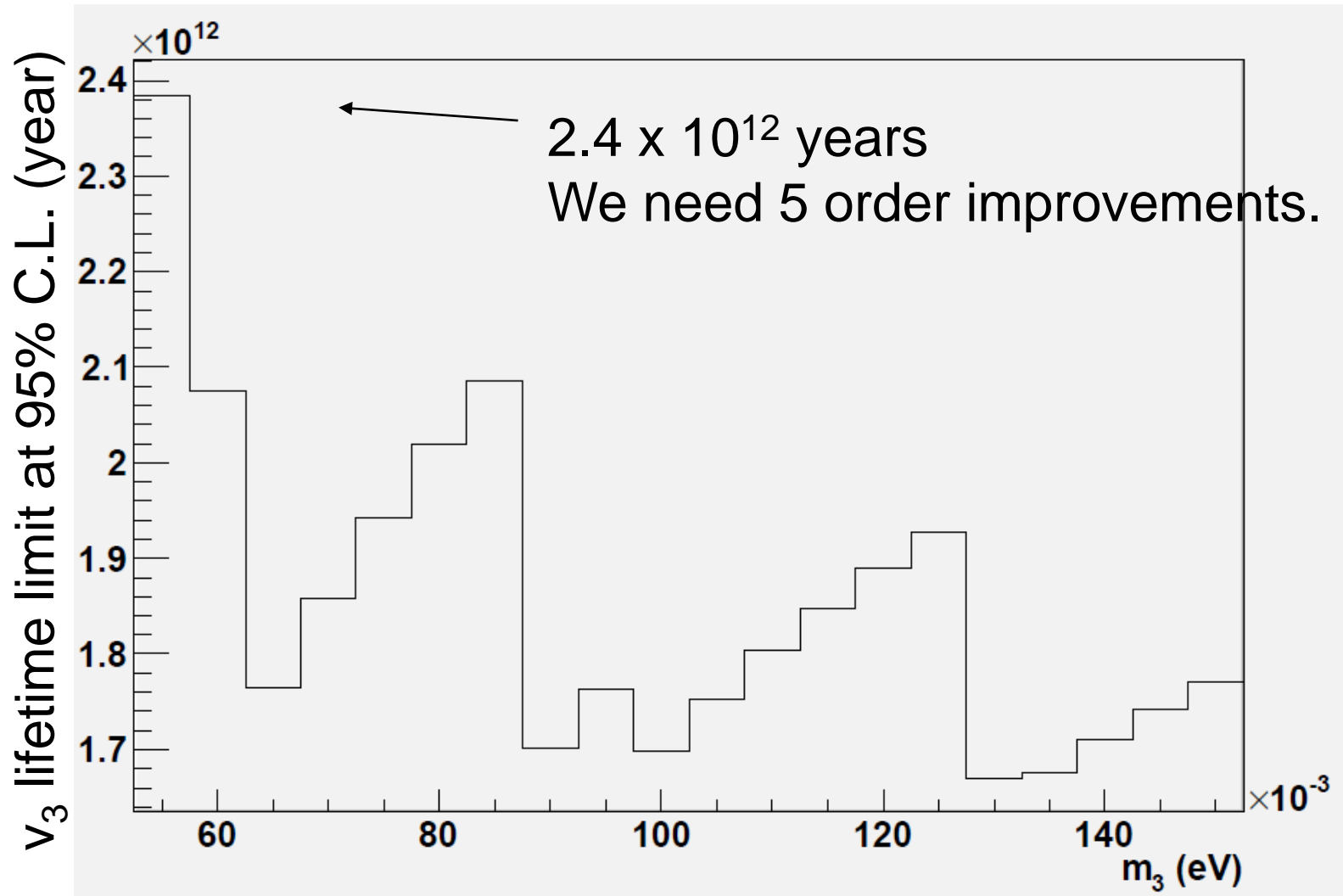


# Lower Limit of Lifetime from the Energy Spectrum Fit the CIB measured by COBE and AKARI(1)



Using the CIB at 60, 100 ( ApJ, 544, 81, 2000 ), 140, 240  $\mu\text{m}$  ( ApJ, 508, 25, 1998 ), 65, 90, 140, and 160  $\mu\text{m}$  ( arXiv:1002.3674, 2010 ), the photon energy spectrum from neutrino radiative decay gives a lifetime lower limit of  $2.4 \times 10^{12} \text{ year}$  at 95% C.L. for  $m_3 = 0.05\text{eV}$  and  $m_2 = 0.01\text{eV}$ .  
( My calculation )

# Lower Limit of Lifetime from the Energy Spectrum Fit to the CIB measured by COBE and AKARI(2)

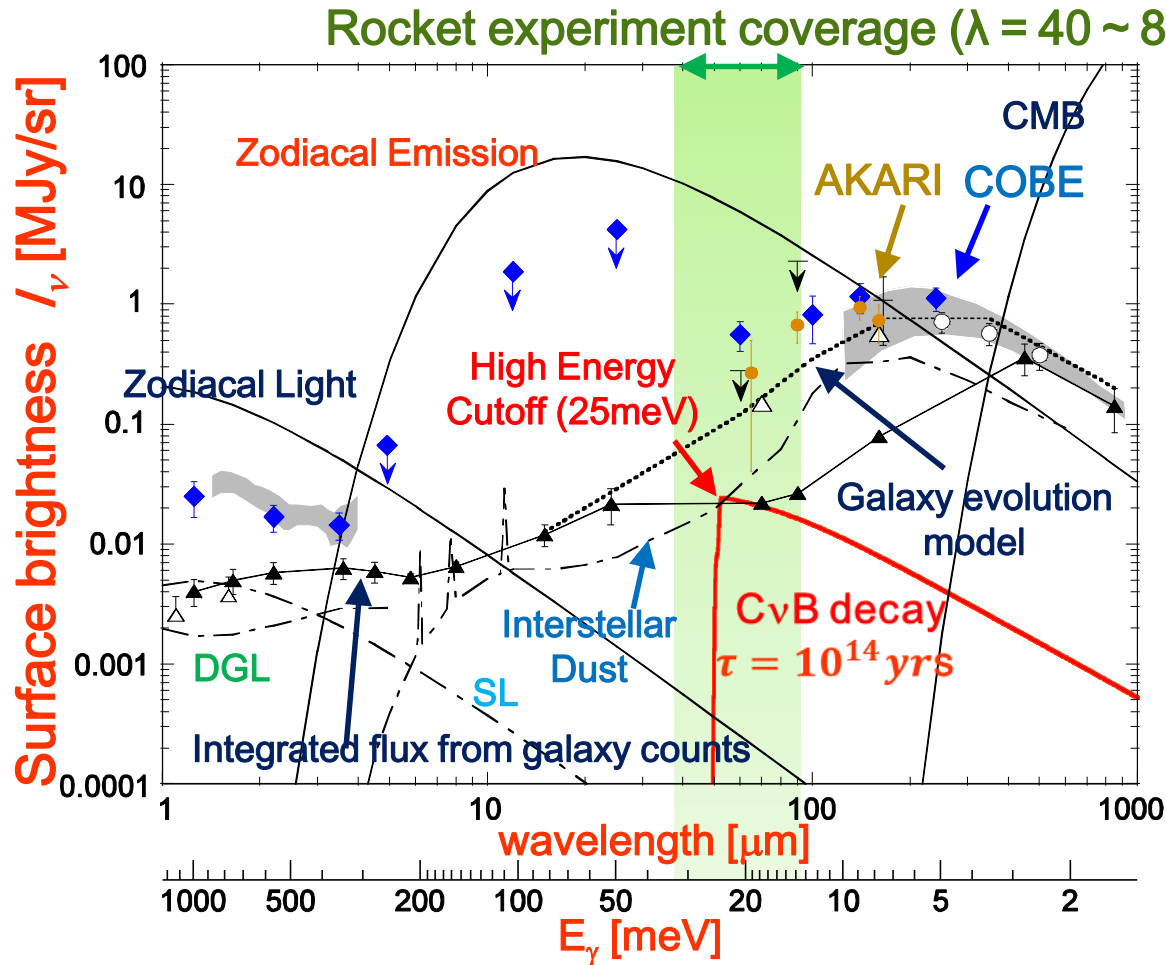


$\nu_3$  mass (eV)

$$\Delta m_{23}^2 = 0.00243 \text{ eV}^2$$

# **Cosmic Background Neutrino Decay Search (COBAND)**

# Signal of Cosmic Background Neutrino Decay and its Backgrounds



CIB  
measurements  
(● AKARI,  
◆ COBE)

By measuring the energy spectrum of the Zodiacal Emission with the CvB decay continuously, we can see the CvB decay signal as a high energy cutoff.

## Requirements for the detector

- Continuous spectrum of photon energy around  $E_\gamma \sim 25$  meV ( $\lambda = 50\mu\text{m}$ )
- Energy measurement for single photon with better than 2% resolution for  $E_\gamma = 25$  meV to identify the sharp edge in the spectrum
- Rocket and/or satellite experiment with this detector



# COBAND (COsmic BACKGROUND Neutrino Decay Search) Experiment

**Rocket Experiment** Plan: 5minutes data acquisition at 200 km height in 2021-22.  
Improve the current limit of lifetime  $\tau(\nu_3)$  by two orders of magnitude ( $\sim 10^{14}$  years).

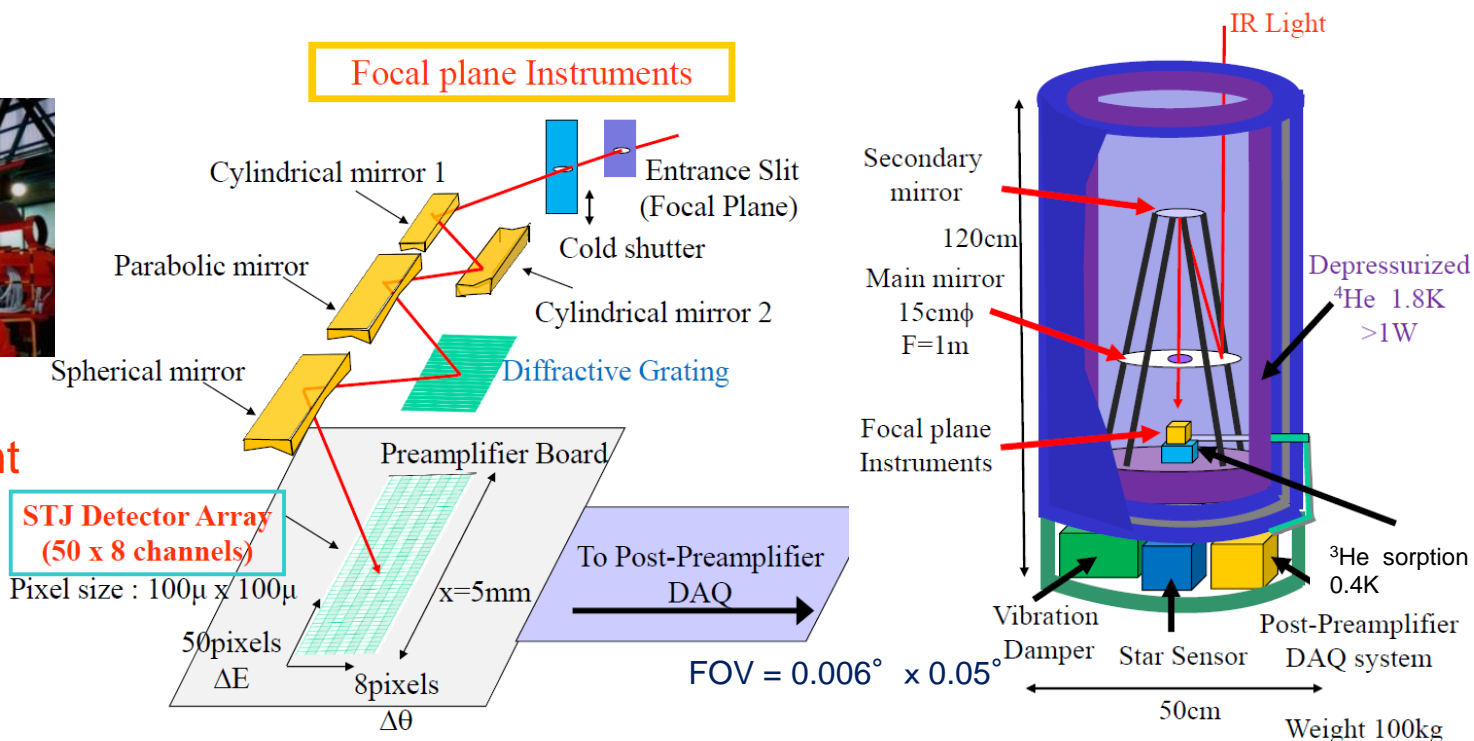
» Superconducting Tunneling Junction (STJ) detectors in development

> Array of 50 Nb/Al-STJ pixels with diffractive grating covering  $\lambda = 40 - 80 \mu\text{m}$



JAXA Rocket  
CIB Experiment

(Feb 2, 1992)



**Satellite experiment after 2025  $\rightarrow$  sensitivity of  $\tau(\nu_3) \sim 10^{17}$  year**

> STJ using Hafnium: Hf-STJ for satellite experiment ( S. H. Kim et al. JPSJ 81,024101 (2012) )

- $\Delta = 20 \mu\text{eV}$  : Superconducting gap energy for Hafnium
- $N_{\text{q.p.}} = 25\text{meV}/1.7\Delta = 735$  for 25meV photon:  $\Delta E/E < 2\%$  if Fano-factor is less than 0.3

# COBAND Rocket Experiment

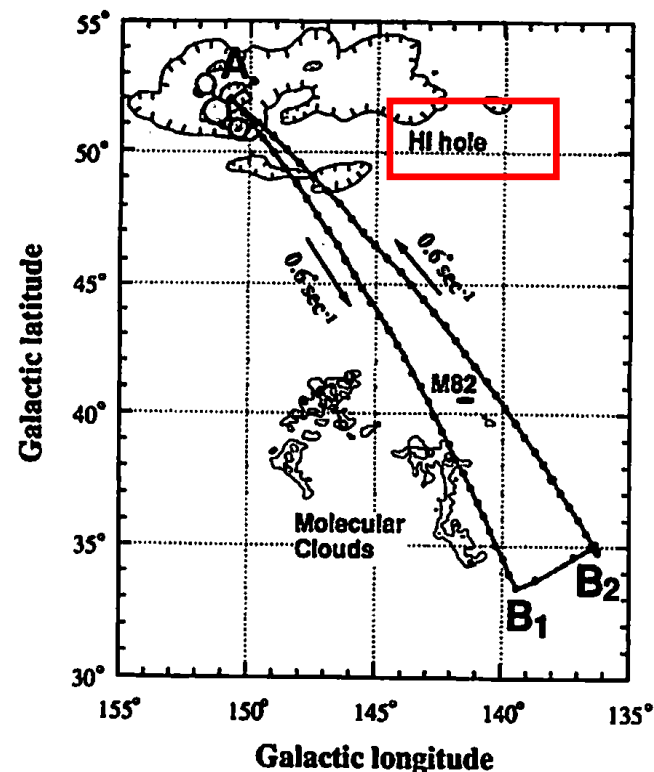
We measured CIB in the same points as S520-15 experiment measured the CIB in 1992.

## a. Pre-flight operation

X - 10h10m	start first liquid He transfer
- 4h10m	start second liquid He transfer (fill up)
	set the launcher angle to $Az=145^\circ$ , $El=85^\circ$
- 1h43m	power on (external supply)
- 40m	start pumping the cryostat tank
- 5m	close the pumping line valve
	switch to the internal power supply
- 4m	disconnect the pumping line
X	launch (1:00:00 JST, 1992 February 2)

## b. In-flight operation

X + 55s	open the nosecone covers, open the pumping line valve
+ 60s	separate the rocket motor
+ 61s	open the gas shade
+ 63s	start attitude control
+ 90s-	point at "A"
+ 130s	open the cryostat lid
+ 220s-255s	scan "A" $\rightarrow$ "B <sub>1</sub> " ( $0.6 \text{ s}^{-1}$ )
+ 255s-	point to "B <sub>2</sub> "
+ 277s-310s	scan "B <sub>2</sub> " $\rightarrow$ "A" ( $0.6 \text{ s}^{-1}$ )
+ 310s-430s	point at "A"
+ 430s-	tip down to the earth limb (recovery operation)
+ 480s	instrument jettison



## Measured Points

These are the same as S520-15.

A (Galactic Latitude 52° Galactic Longitude 151° )

# Zodiacal Emission

Thermal emission from the interplanetary dust cloud

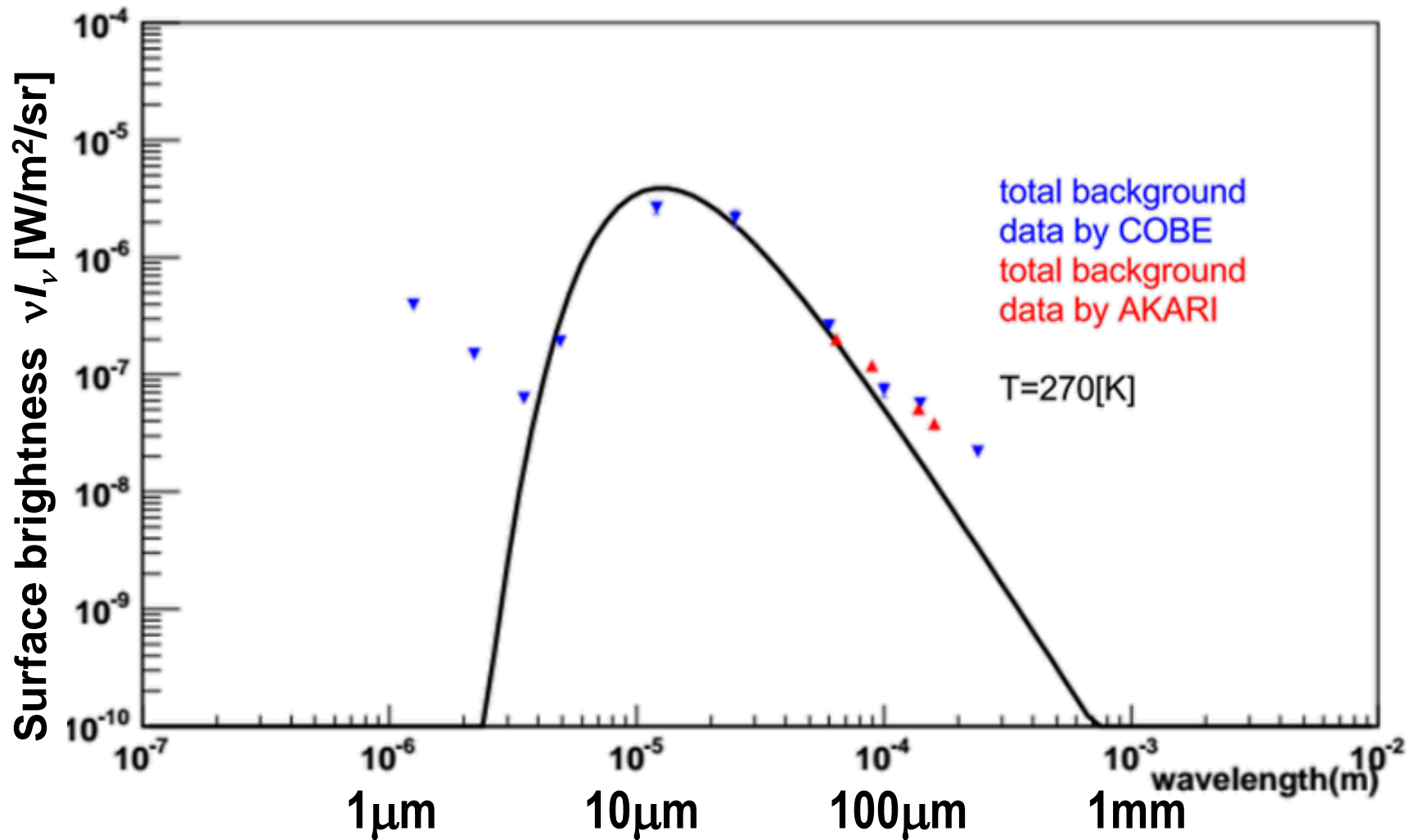
$$I_\nu = \frac{2h\nu^3}{c^2} \frac{1}{\exp(h\nu/kT) - 1} \\ \times A \left( \frac{\nu}{c} \times 10^{-5} \right)^B \text{ Wm}^{-2}\text{sr}^{-1}$$

$$T = 270K, A = 6 \times 10^{-8}, B = 0.3$$

$$h \text{ [Js]}, c \text{ [m/s]}, \lambda \text{ [m]}$$

Zodiacal Emission(ZE) is overwhelmingly dominating. Here we consider only ZE as the background.

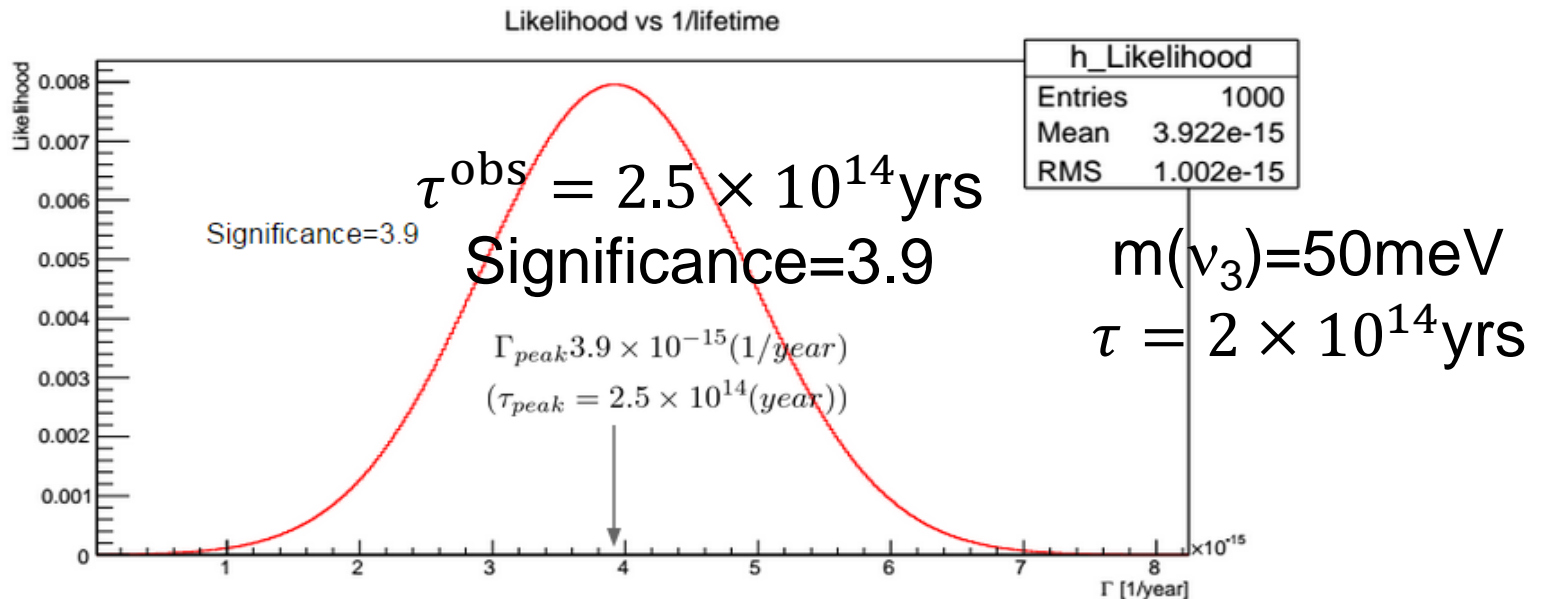
# Zodiacal Emission





# Discovery Potential for Neutrino Decay

1. Create a pseudo-data for ND and ZE expectation on assumption of  $\nu_3$  mass and lifetime:  $N_{\text{obs}}(m^{\text{true}}, \Gamma^{\text{true}})$
2. Calculate expected distribution of ND+ZE on assumption of  $\nu_3$  mass and lifetime:  $N_{\text{exp}}(m, \Gamma)$
3. Calculate likelihood of the pseudo-data as a function of the decay width for each neutrino mass
4. Obtain the most probable decay width and its significance



# Sensitivity to Neutrino Decay

Parameters in the rocket experiment simulation

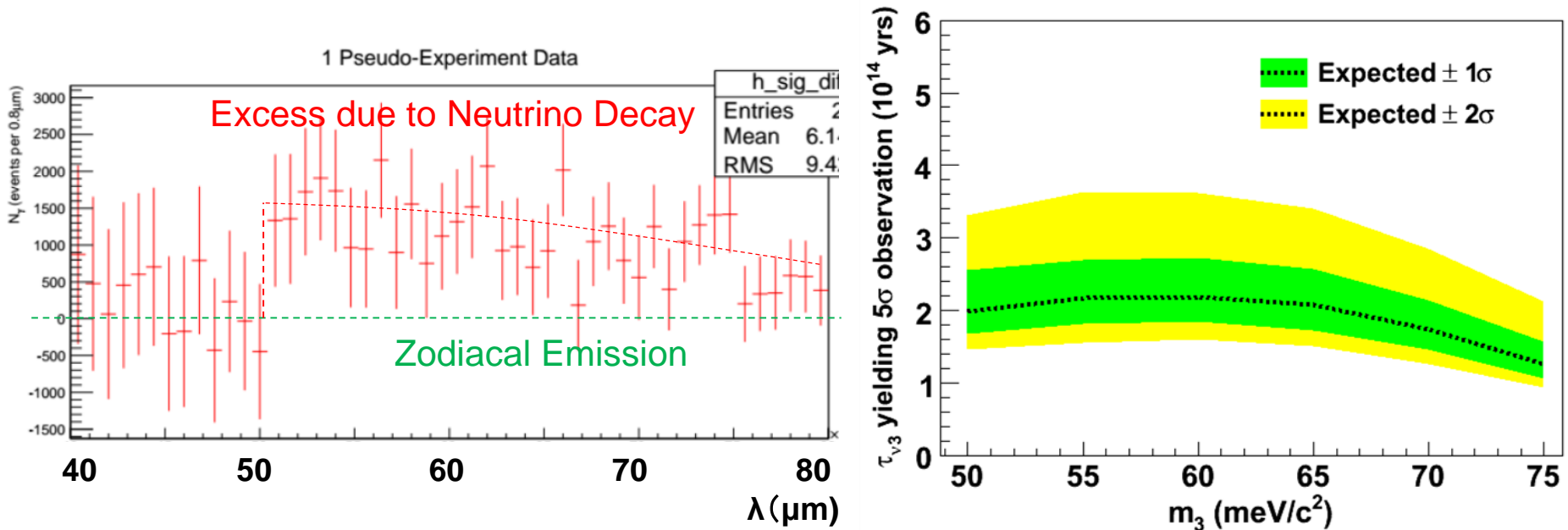
telescope diameter: 15cm

50-column ( $\lambda$ : 40 $\mu\text{m}$  – 80  $\mu\text{m}$ )  $\times$  8-row array

Viewing angle per single pixel: 100 $\mu\text{rad}$   $\times$  100 $\mu\text{rad}$

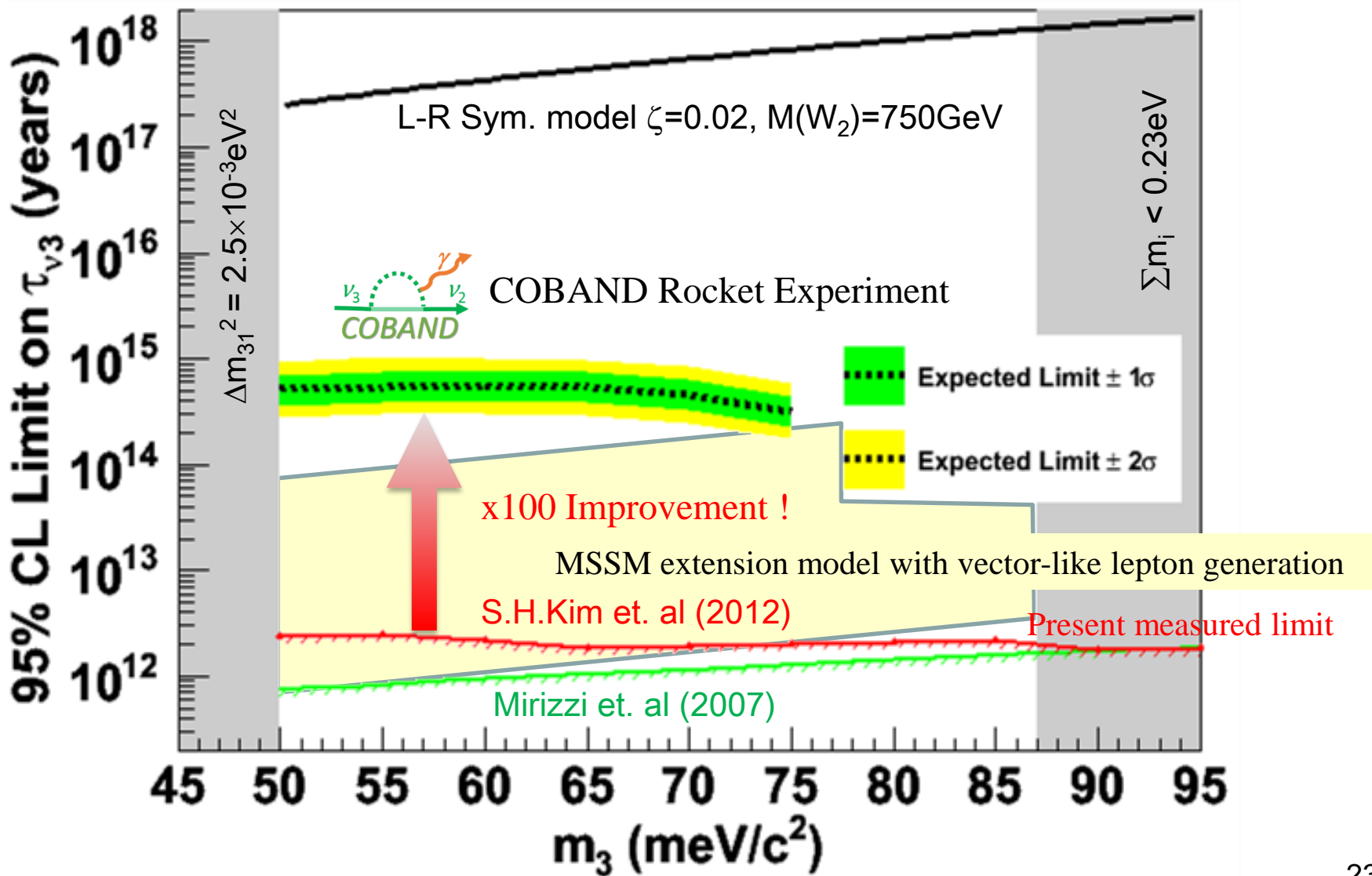
Measurement time: 200 sec.

Photon detection efficiency: 100%



- If  $\nu_3$  lifetime were  $2 \times 10^{14}$  yrs, the signal significance is at 5 $\sigma$  level

# COBAND Experiment Sensitivity to Neutrino Decay



# Requirement for the photon detector in COBAND rocket experiment

- Sensitive area of  $100\mu\text{m} \times 100\mu\text{m}$  for each pixel
- High detection efficiency for a far-infrared single-photon in  $\lambda = 40\mu\text{m} \sim 80\mu\text{m}$
- Dark count rate less than 300Hz (expected real photon rate)

$$\rightarrow \text{NEP} = \epsilon_{\gamma} \sqrt{2f_{\gamma}} \sim 1 \times 10^{-19} \text{ W} / \sqrt{\text{Hz}}$$

(Noise Equivalent Power) , where  $\epsilon_{\gamma}$  is a photon energy and  $f_{\gamma}$  is a photon rate.

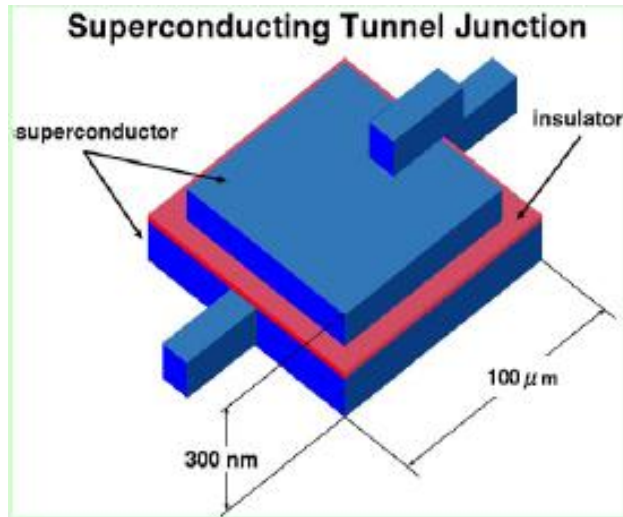
We are trying to achieve  $\text{NEP} \sim 10^{-19} \text{ W} / \sqrt{\text{Hz}}$  by using

- Superconducting Tunneling Junction detector  
(leakage current per pixel  $< 100\text{pA}$ )
- Cryogenic amplifier readout

# R&D Status of Superconducting Tunnel Junction Detector for COBAND experiment

# STJ (Superconducting Tunnel Junction) Detector

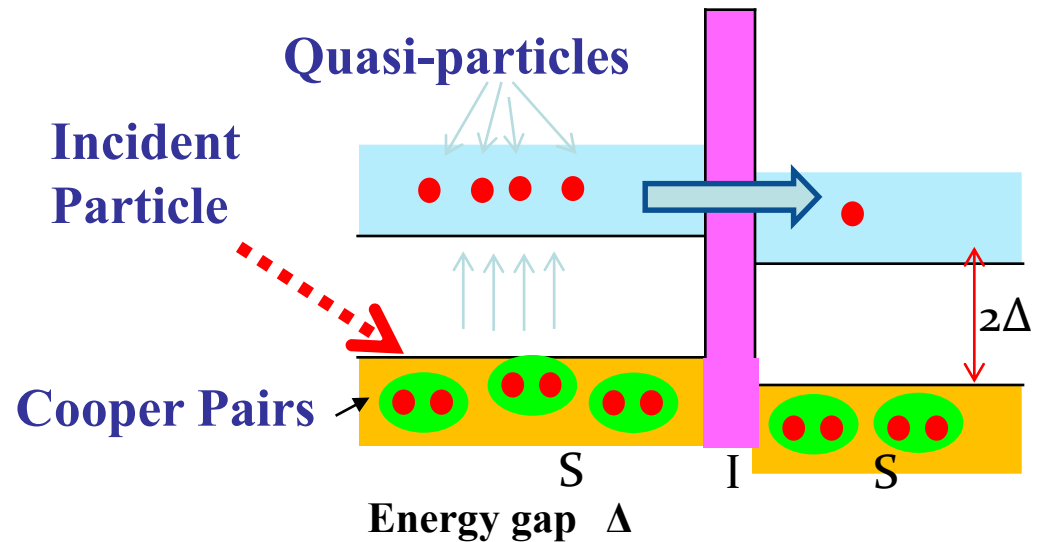
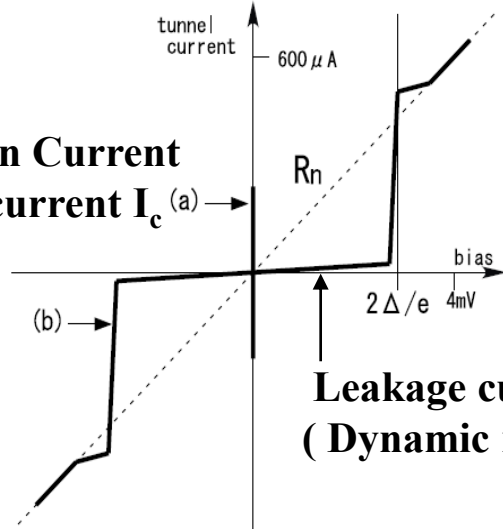
Superconductor / Insulator / Superconductor Josephson Junction



At the superconducting junction, quasi-particles over their energy gap go through tunnel barrier by a tunnel effect. By measuring the tunnel current of quasi-particles excited by an incident particle, we measure the energy of the particle.

current-voltage (I-V) curve for STJ

Josephson Current  
Critical current  $I_c$  (a) →



Material	$T_c(K)$	$\Delta(\text{meV})$
Niobium	9.20	1.550
Aluminum	1.14	0.172
Hafnium	0.13	0.021



# STJ Energy Resolution

## STJ Energy Resolution

$$\sigma_E = \sqrt{1.7\Delta(FE)}$$

Using Hf as a superconductor,

$$\sigma_E / E = 1.7\% \quad \text{at } E = 25\text{meV}$$

$\Delta$ : Band gap energy

F: Fano factor (= 0.2)

E: Incident particle energy

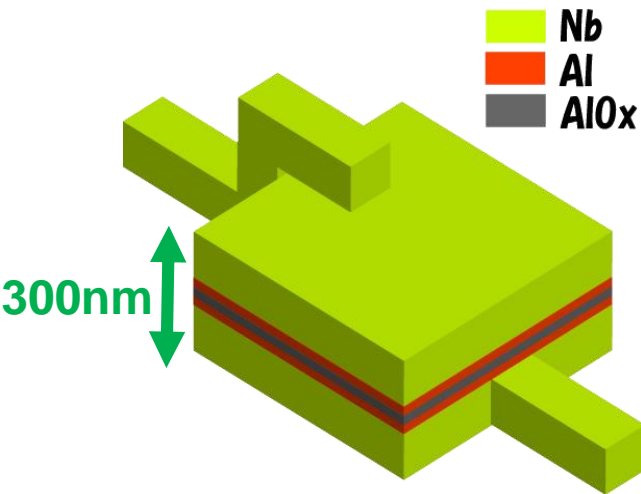
Material	$T_c(K)$	$\Delta(\text{meV})$
Niobium	9.20	1.550
Aluminum	1.14	0.172
Hafnium	0.13	0.021

### **Tc : Critical Temperature**

Operation is done at a temperature around 1/10 of Tc

We reported that Hf-STJ worked as a STJ in 2011.

# Nb/Al-STJ Photon Detector



## Number of Quasi-particles in Nb/Al-STJ

$$N_q = G_{Al} E_0 / 1.7 \Delta$$

$G_{Al}$  : Trapping Gain in Al (~10)

$E_0$  : Photon Energy

$\Delta$  : E-Gap in superconductor

For 25meV single photon

$$N_q = 10 \frac{25 \text{ meV}}{1.7 * 0.57 \text{ meV}} = 250 \text{ e}$$

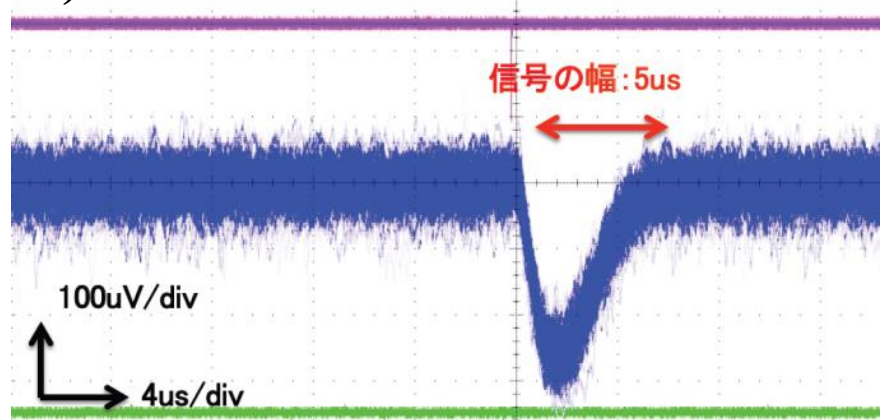
## Back tunneling Effect → Trapping Gain

Quasi-particles near the barrier can mediate Cooper pairs, resulting in true signal gain

- Bi-layer fabricated with superconductors of different gaps  $\Delta_{Nb} > \Delta_{Al}$  to enhance quasi-particle density near the barrier
- Nb(200nm)/Al(70nm)/AlOx/Al(70nm)/Nb(100nm)

$$\Delta_{Nb/Al} = 0.57 \text{ meV}$$

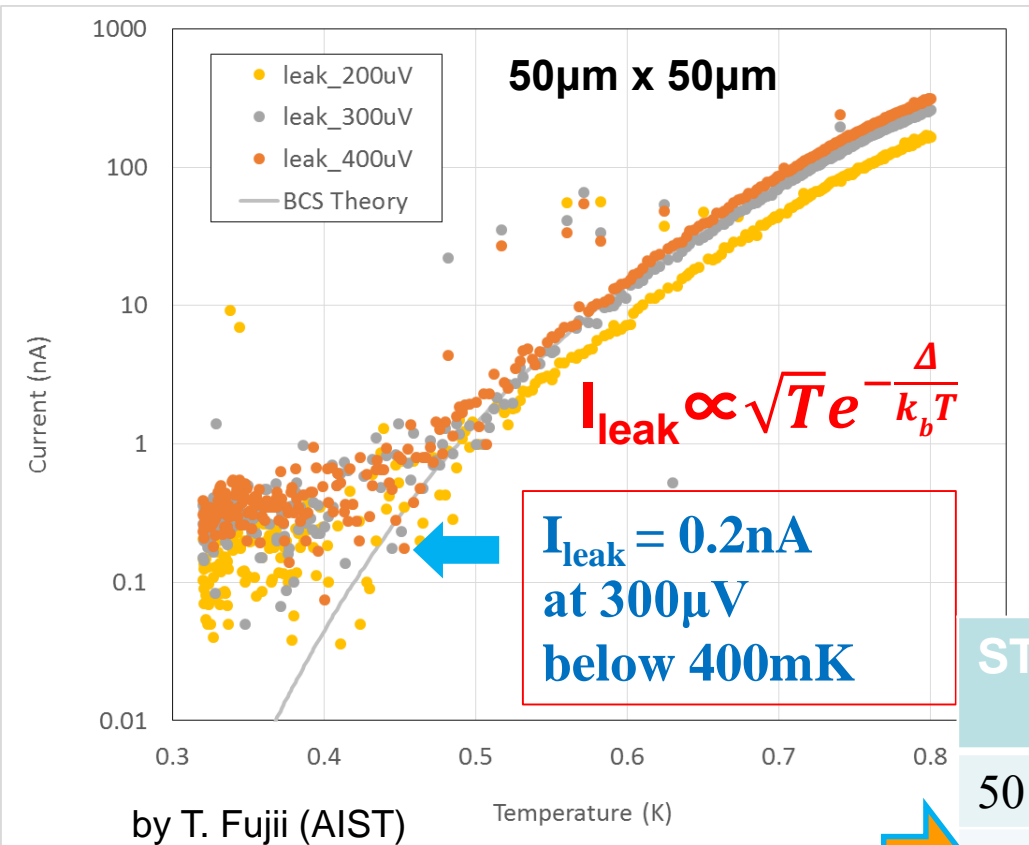
Response of Nb/Al-STJ to visible laser light pulse ( $\lambda=465\text{nm}$ ) at 350mK



# Leakage Current of Nb/Al-STJ

- Leakage current  $I_{\text{leak}}$  is required to be below 0.1nA to detect a single far-infrared photon ( $\lambda = 40 - 80\mu\text{m}$ ).

## Temperature Dependence of Leakage Current



In 2014,  
AIST group joined us and produced  
Nb/Al-STJ with AIST CRAVITY  
processing system.  
Leakage current has satisfied our  
requirement of 0.1nA .

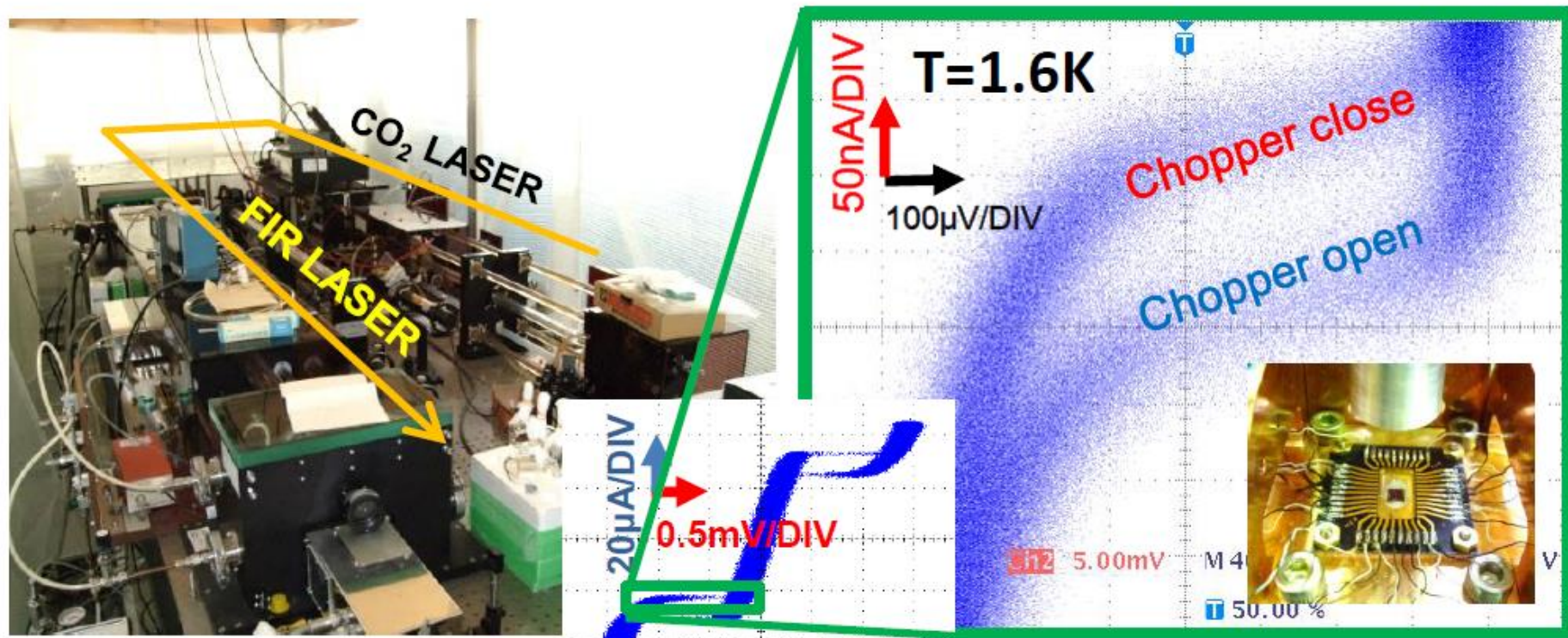


STJ size	# of samples	$I_{\text{leak}}$ at 0.3mV
50 x 50 $\mu\text{m}^2$	18	224 $\pm$ 29 pA
20 x 20 $\mu\text{m}^2$	7	39 $\pm$ 13 pA
10 x 10 $\mu\text{m}^2$	20	14 $\pm$ 7 pA

# Test Results of Nb/Al-STJ with Far-Infrared laser

Far-Infrared Laser at University of Fukui  
(  $\lambda=57.2\mu\text{m}$  )

Nb/Al-STJ Response to Far-Infrared Laser



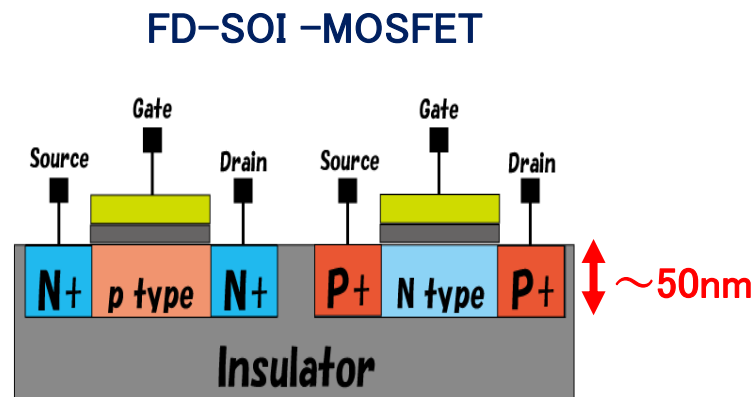
- 20μm-square Nb/Al-STJ made at AIST CRAVITY system
- Laser light was turned on and off with a chopper at a frequency of 200Hz. Measured the change of the I-V curve between the laser on and off to be 50~100nA in current.

## R&D Status of SOI Cryogenic Preamplifier for STJ

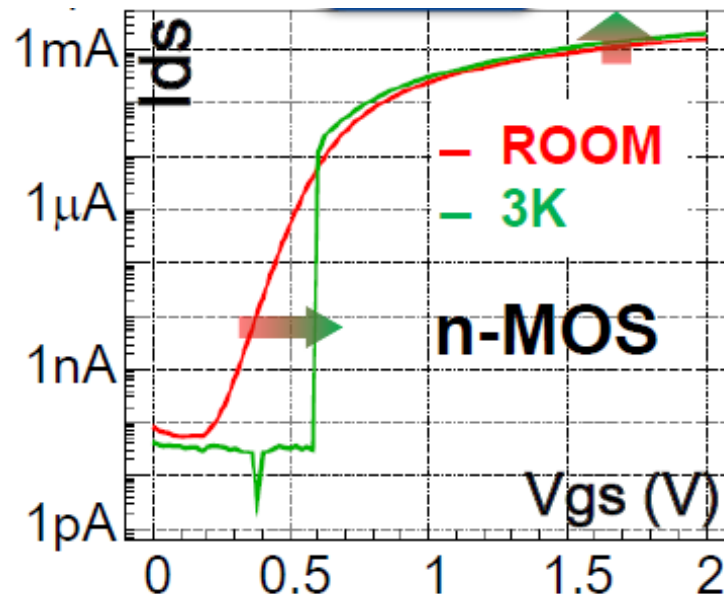
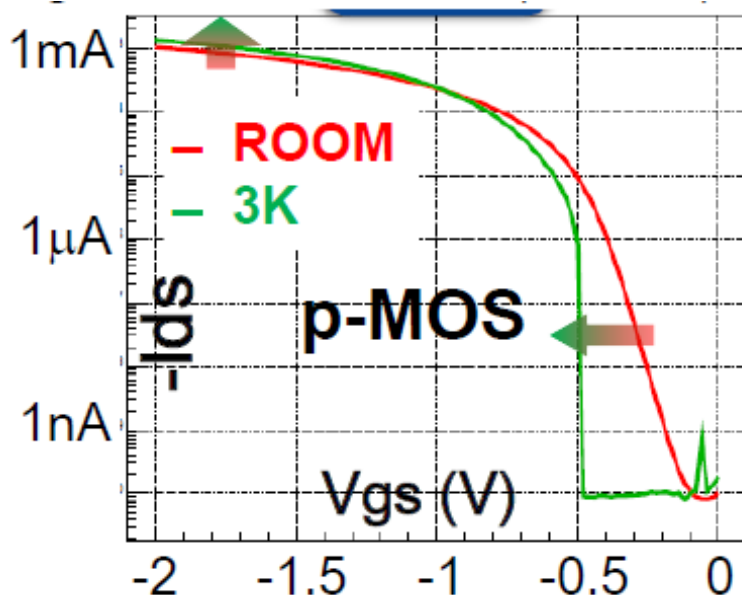


# FD-SOI-MOSFET at Cryogenic Temperature

FD-SOI (Fully Depleted Silicon-On-Insulator) device was proved to operate at 4K by a JAXA/KEK group (AIPC 1185,286-289(2009)). It has the following characteristics:  
low-power consumption, high speed, easy large scale integration and suppression of charge-up by high mobility carrier due to thin depletion layer( $\sim 50\text{nm}$ ).



**$I_{ds}$ - $V_{gs}$  Curve of  $W/L=10\mu\text{m}/0.4\mu\text{m}$  at  $|V_{ds}| = 1.8\text{V}$**



**Both p-MOS and n-MOS show excellent performance at 3K and below.**

# SOI Cryogenic Amplifier

## SOI-STJ4 (the 4<sup>th</sup> prototype)

We updated the SOI cryogenic Amplifier for Nb/Al-STJ.

### Amplification

Replace the resistance by a SOIFET as a current source (M2).

### Feedback

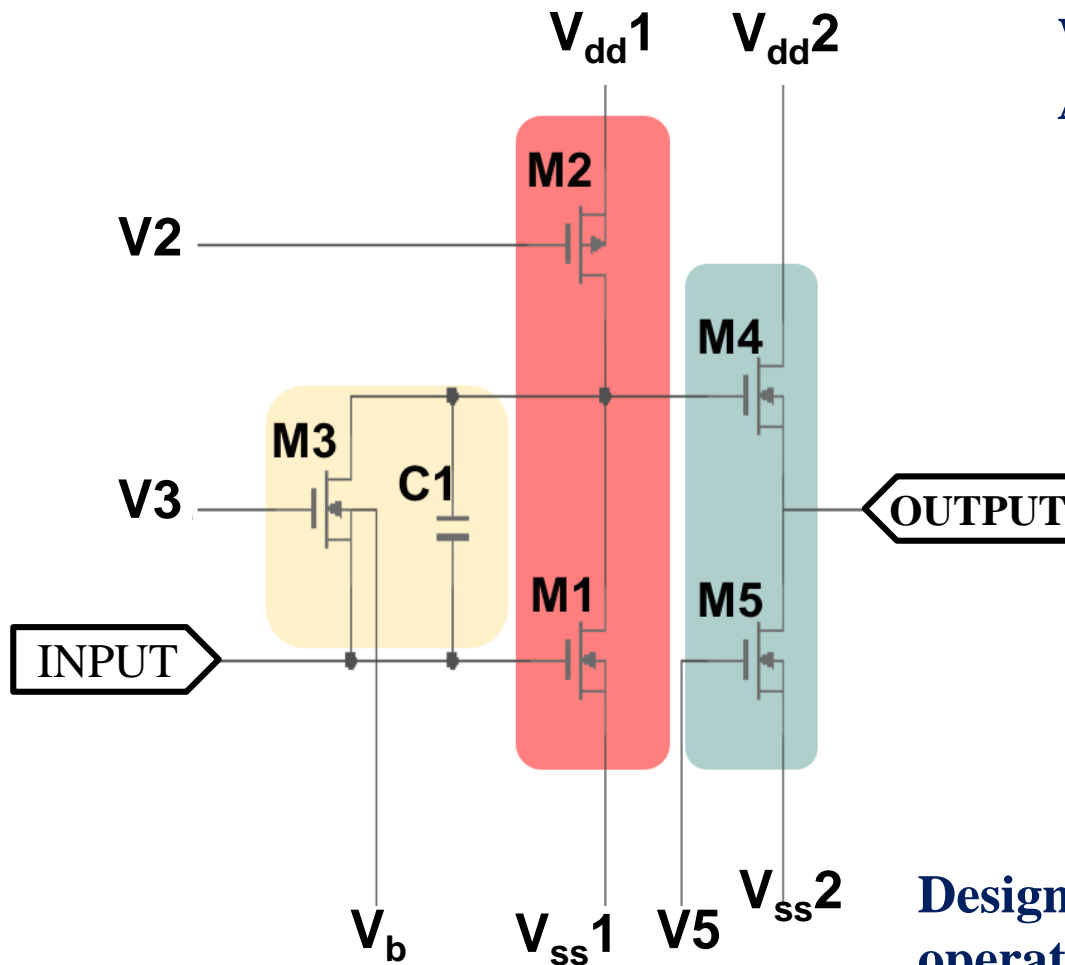
Use the feedback between the drain and the gate of M1 to apply a stable bias voltage (M3).

### Buffer

Add the follower to reduce the output impedance (M4 and M5).

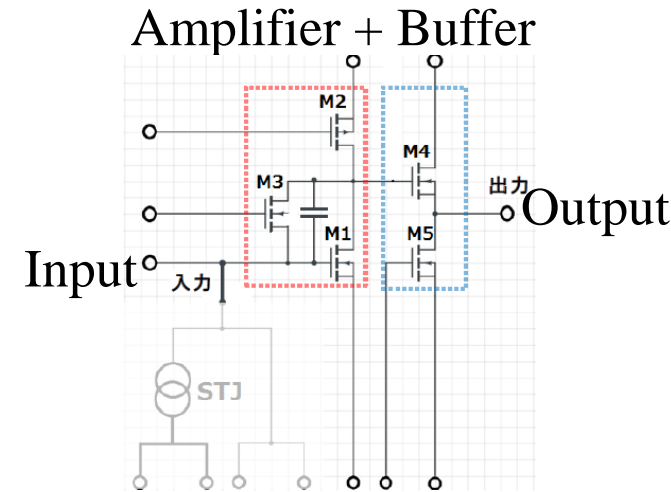
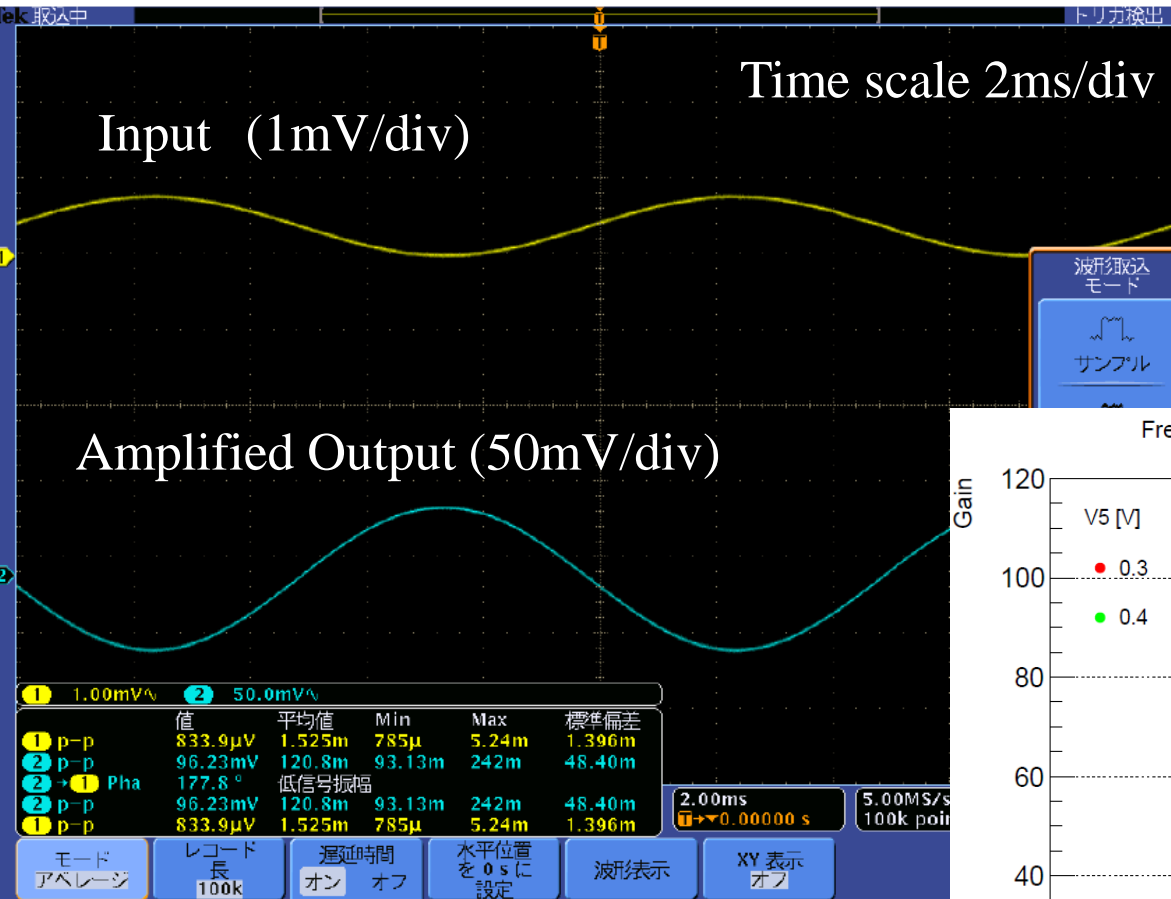
Designed the ratio (W/L) to set the operation power consumption below 120 $\mu$ W.

This SOI amplifier board was made by LAPIS semiconductor company.

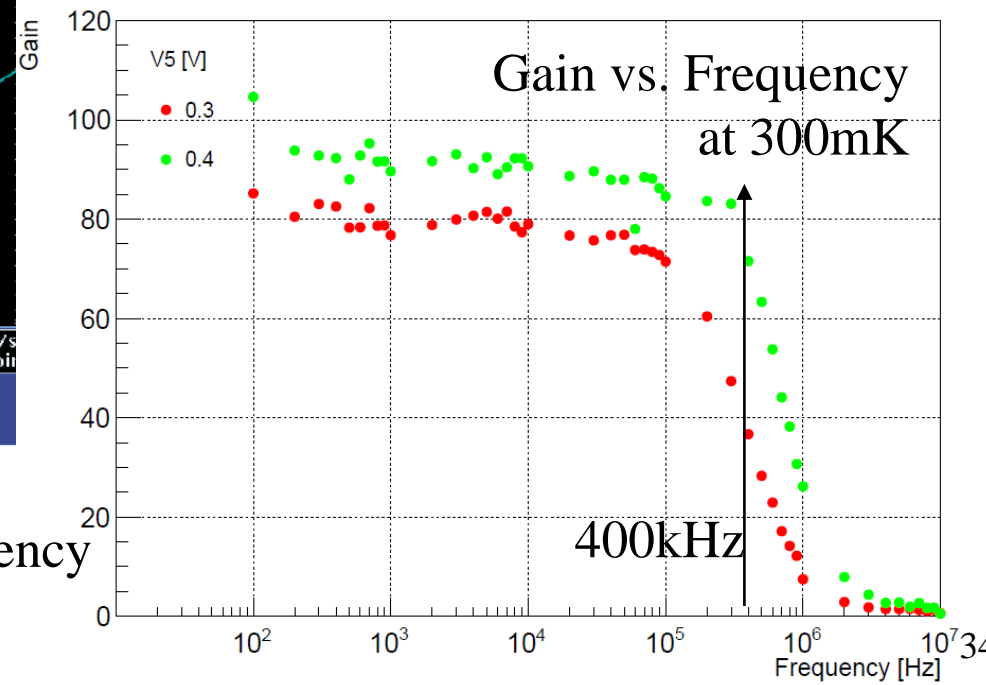


# Test Results of the SOI Cryogenic Amplifier

## Input and Amplified Output

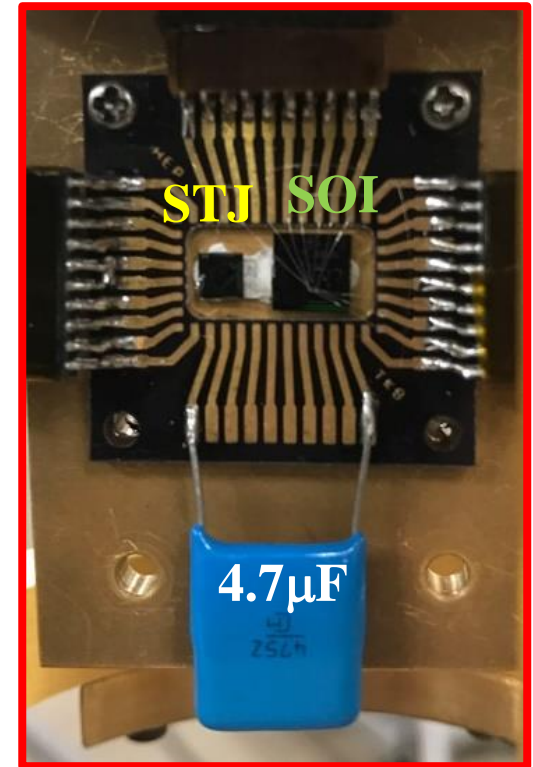
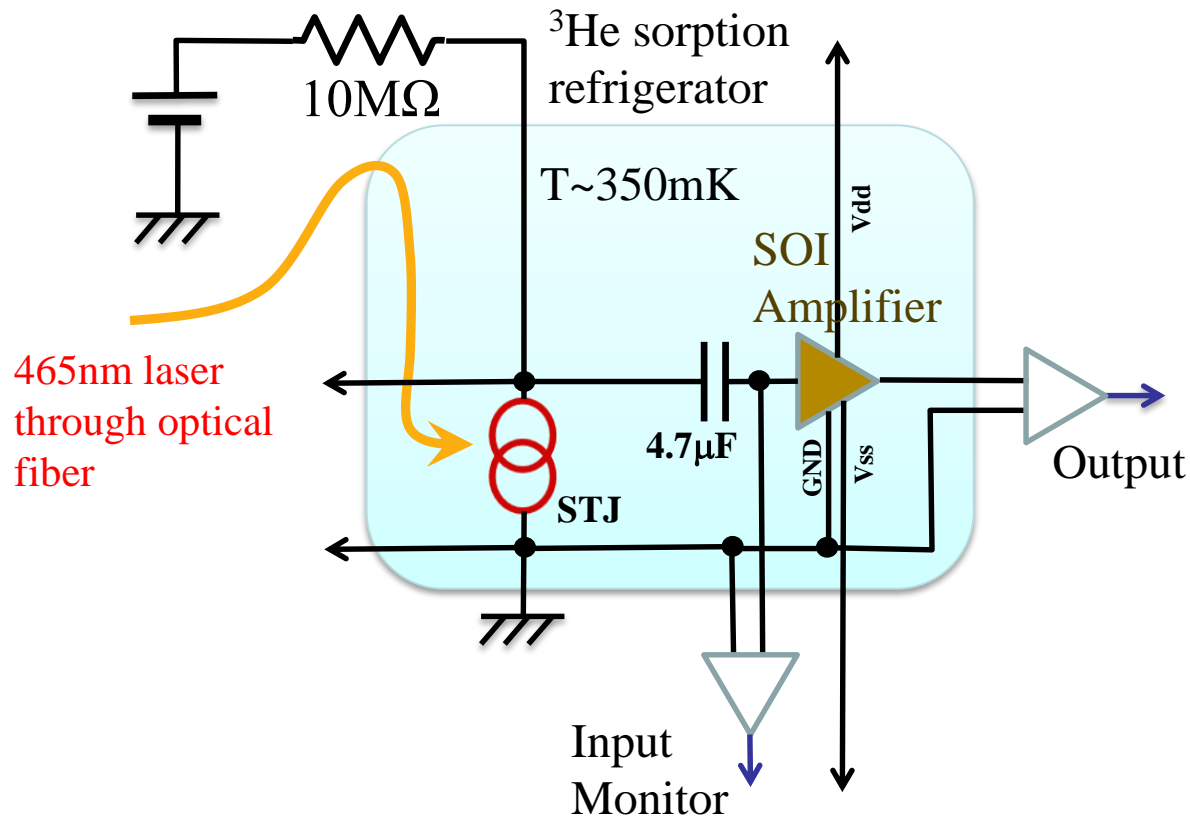


Frequency characteristic of cold amplifier(SOISTJ4) at 300mK



Gain of 80 was achieved for a signal frequency up to 400kHz signals at 300mK.

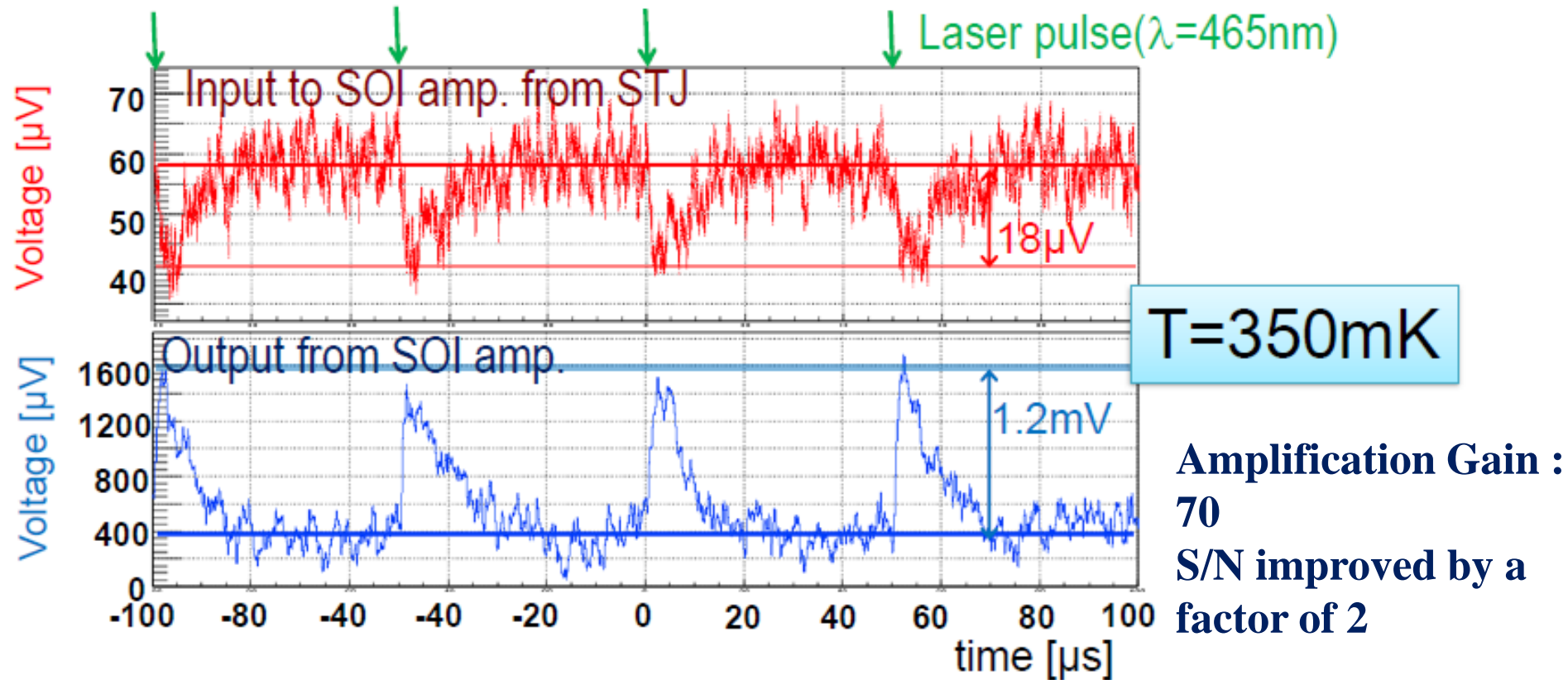
# Setup of STJ Signal Amplification with the SOI Cryogenic Amplifier



- $20\mu\text{m}$ -square Nb/Al-STJ with SOI-STJ4 amplifier through  $4.7\mu\text{F}$  capacitance.
- Input impedance of the SOI amplifier is about  $20\text{k}\Omega$ .
  - **STJ operation at a constant current mode.**
  - STJ bias cable capacitance is around  $1\text{nF}$ :  $Z=160\Omega$  for  $1\mu\text{s}$  signal.

# STJ signal amplified with the SOI cryogenic preamplifier

Nb/Al-STJ laser light response signal was amplified with this SOI cryogenic preamplifier at 350mK.



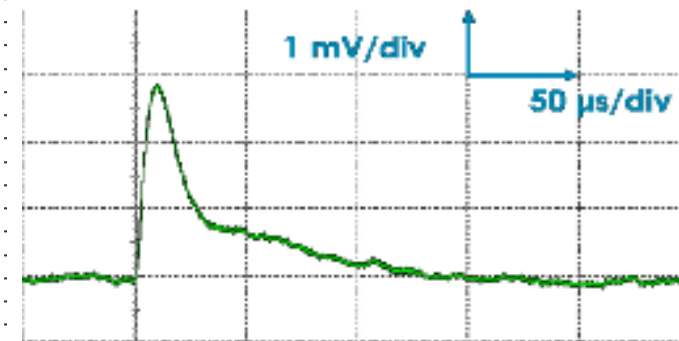
Development of the SOI cryogenic preamplifier is now moving to the stage of the final design for COBAND experiment.



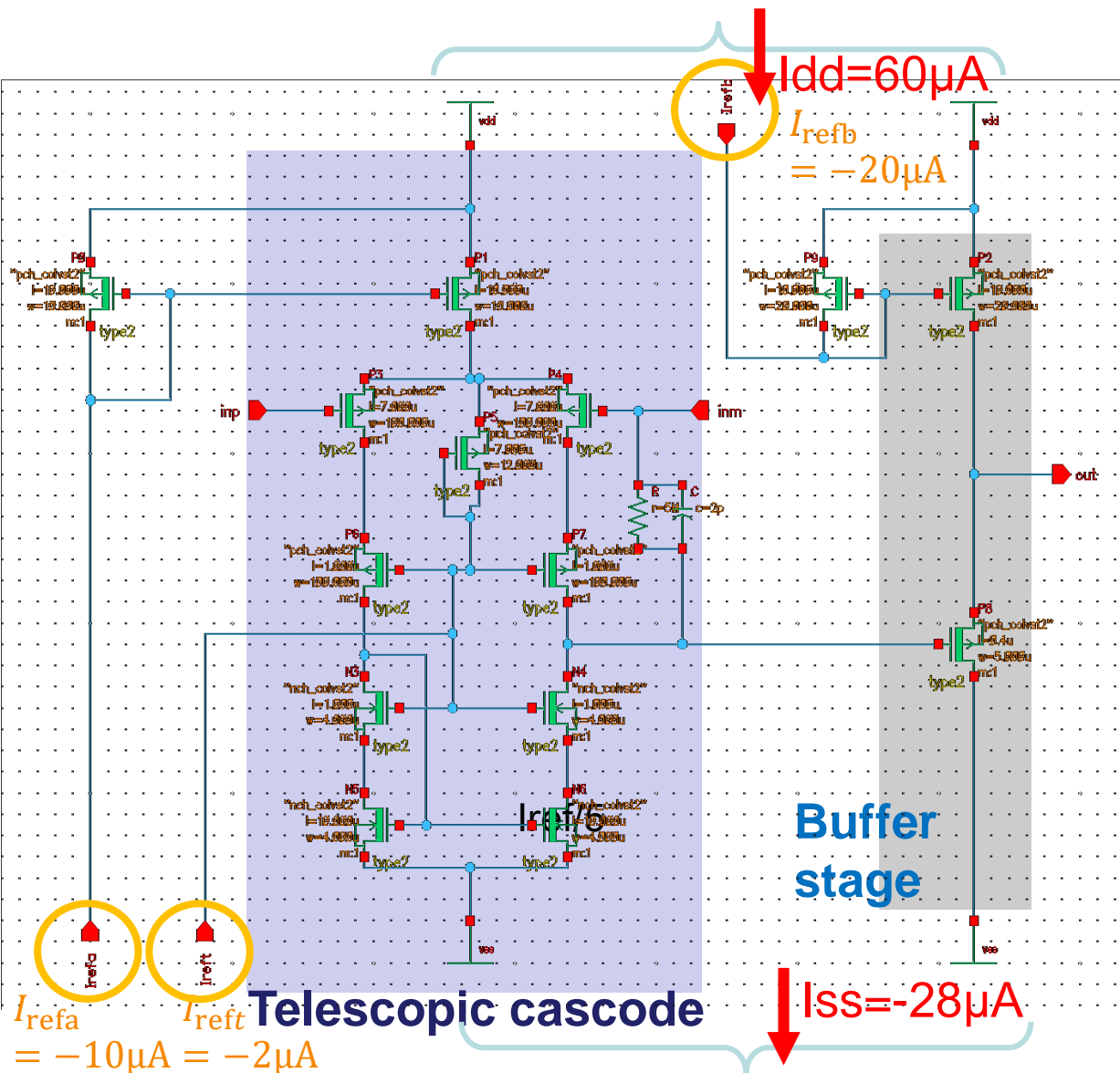
# Charge Amp. Circuit for STJ (SOI-STJ5 design)

- telescopic cascode differential amplifier
- Feedback  $C(2\text{pF}) \times R(5\text{M}\Omega) = 10\mu\text{s}$
- Power consumption  $\sim 150\mu\text{W}$

Test Result of this cryogenic charge amplifier.



STJ signal for visible laser pulse light was amplified by this cryogenic charge amplifier at 350mK.



The schematic illustrates a multi-stage amplifier circuit. It begins with a telescopic cascode stage (yellow background) featuring NMOS and PMOS transistors with various biasing and signal nodes. This is followed by a feedback loop (pink background) containing a resistor network and a capacitor. The final stage is a buffer stage (blue background) consisting of a PMOS transistor and a current source. The circuit is powered by multiple supply rails (Vdd1, Vdd2, Vdd3, Vss1, Vss2, Vss3) and includes input and output nodes for signal processing.

## Feedback capacitance

$$2\text{pF} \rightarrow 60\text{fF}$$

## Power Consumption

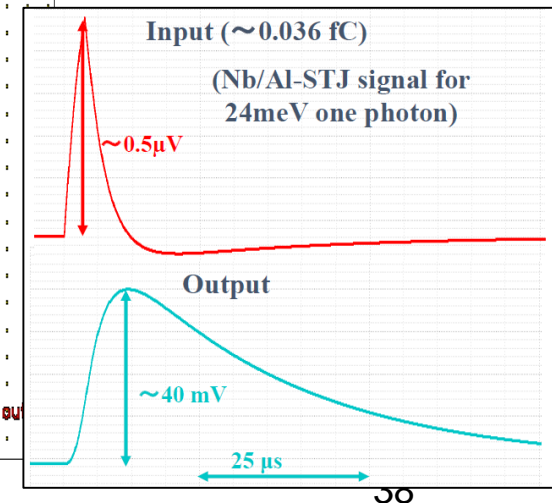
 $\sim 150 \mu\text{W}$ 

## 24meV one photon

(0.03fC) gives  $\sim 40\text{mV}$

## Output

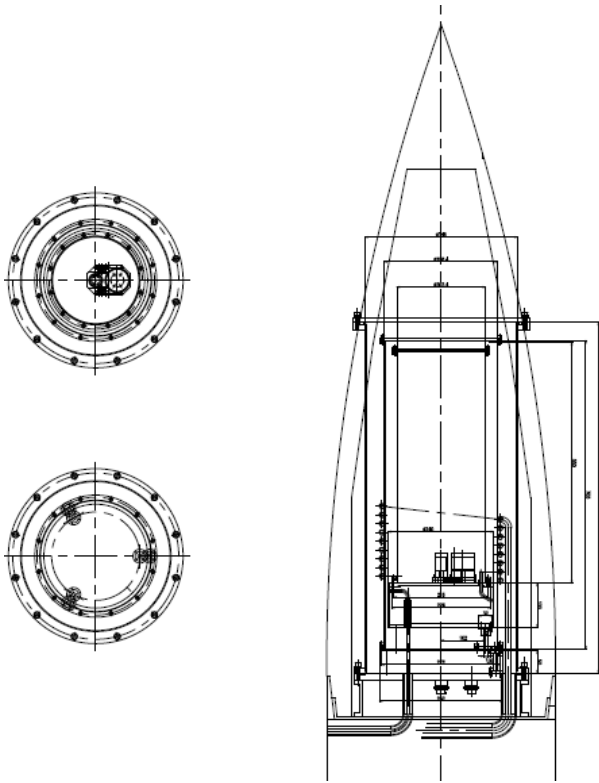
## Simulation Result



**This charge amplifier is now at our University and will be tested soon.**

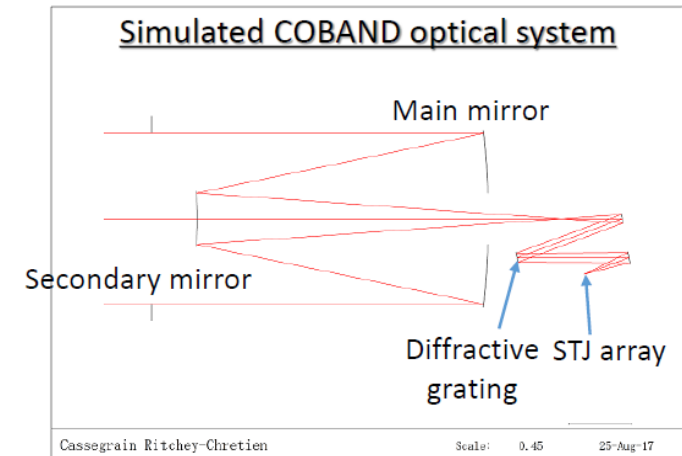
# Prototype of Cryostat for the Rocket Experiment

In April 2017, a prototype of the cryostat for the rocket experiment was made. This is a  $^3\text{He}$  sorption 0.3K refrigerator inside of  $^4\text{He}$  depressurized 1.8K refrigerator. (For now it does not have a  $^3\text{He}$  sorption 0.3K refrigerator inside yet.)



Made by Jeck Tohri company.

We are working on the simulation study on the optics in order to make a final design of COBAND optical system.

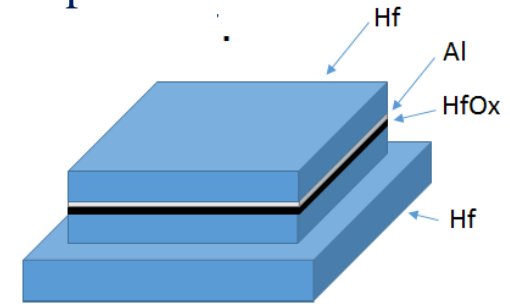


We will test Nb/Al-STJ inside of this cryostat with far-infrared( $50\mu$ ) photon beam at Fukui this winter.

## R&D Status of Hf-STJ

# R&D Status of Hf-STJ - Laser Light Response

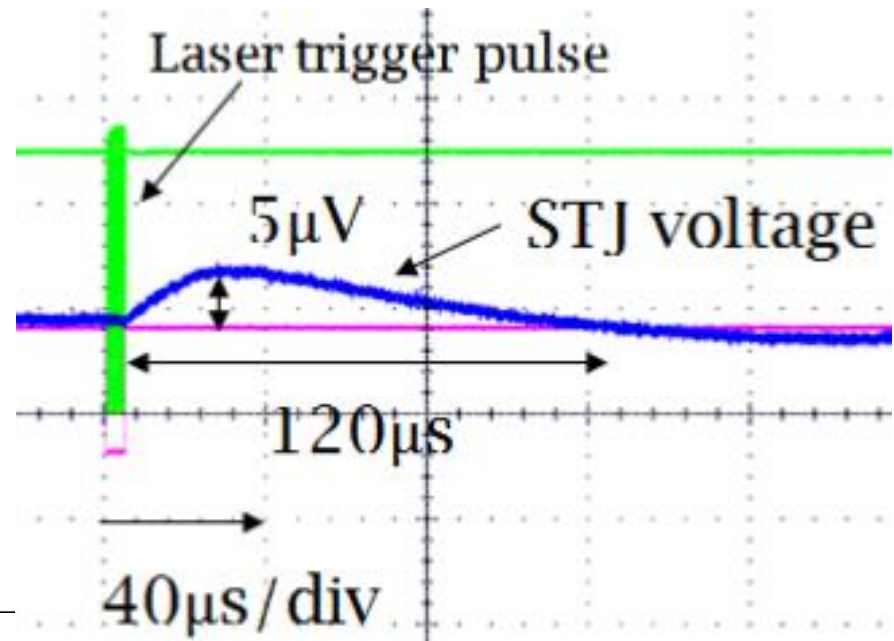
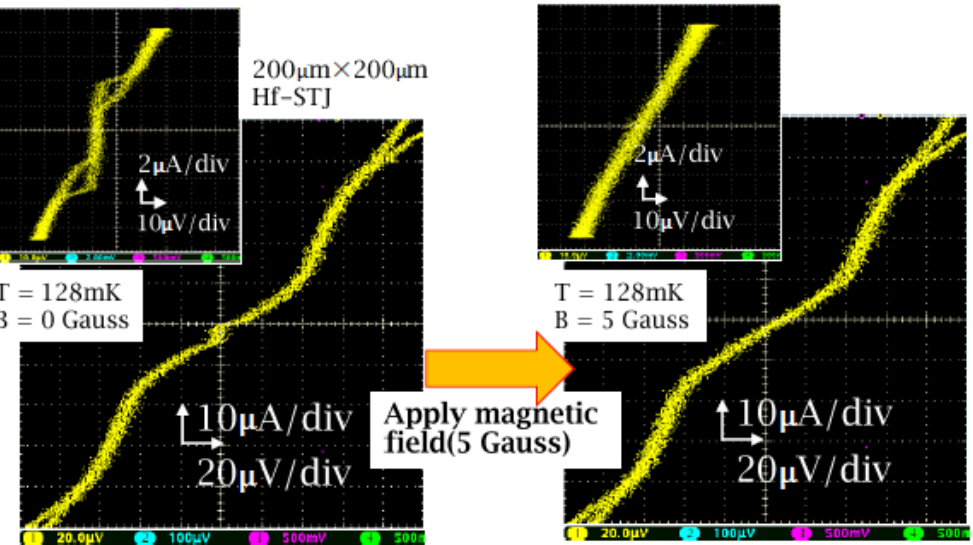
We made a thin aluminum layer (9nm) on the HfO layer (1-2 nm) to improve the insulation of the HfO<sub>x</sub> layer. Hf/Al/HfO<sub>x</sub>/Hf-STJ



$$\Delta = 20 \sim 30 \mu\text{eV}$$

Leakage current =  $5 \mu\text{A}$  @  $128\text{mK}$   
for  $200 \mu\text{m}$ -square sample.

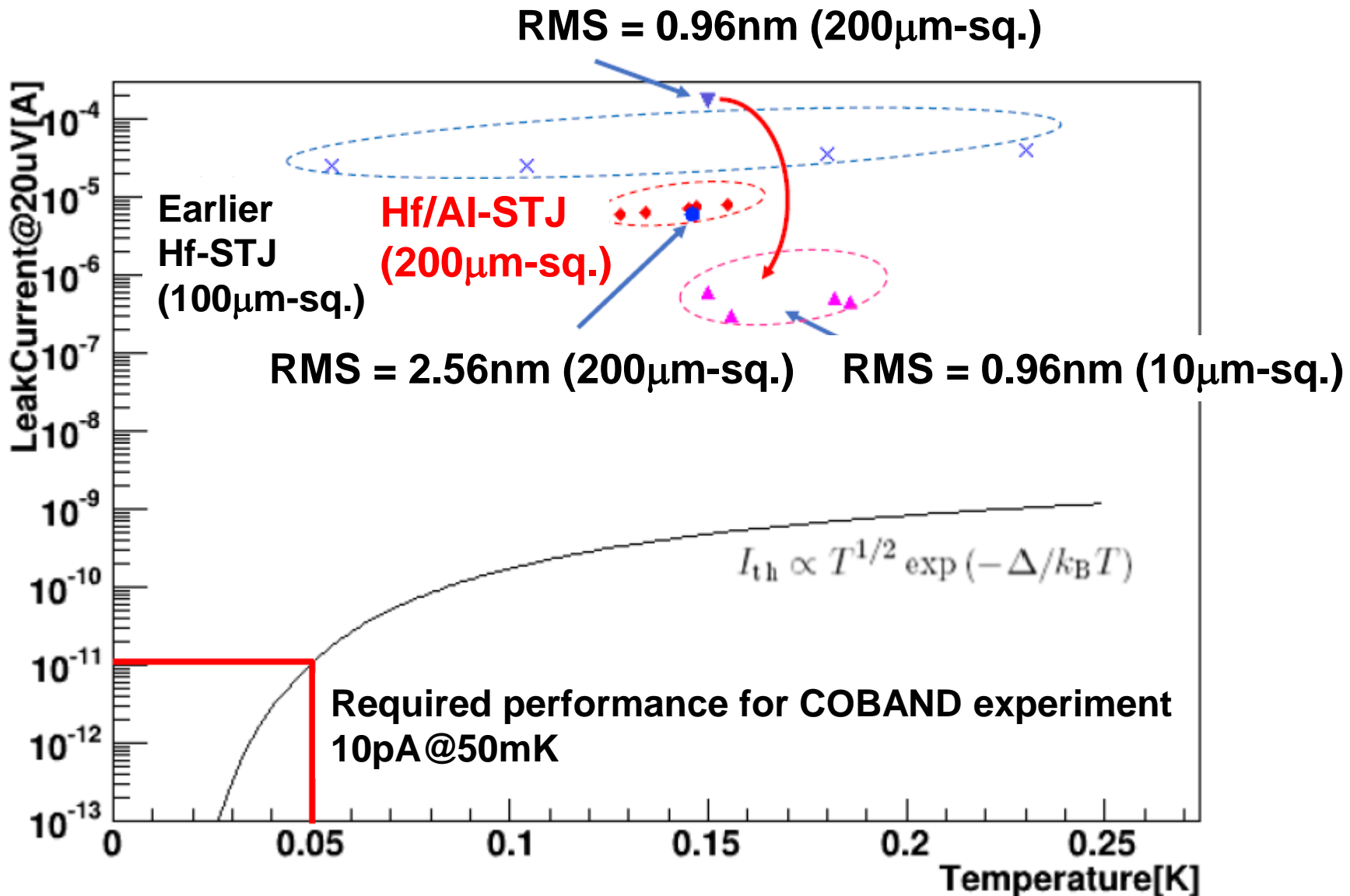
Visible light laser ( $\lambda=465\text{nm}$ ) 10Hz duration



Response speed ( $120 \mu\text{s}$ ) is slower than Nb/Al-STJ response speed ( around a few  $\mu\text{s}$ ).



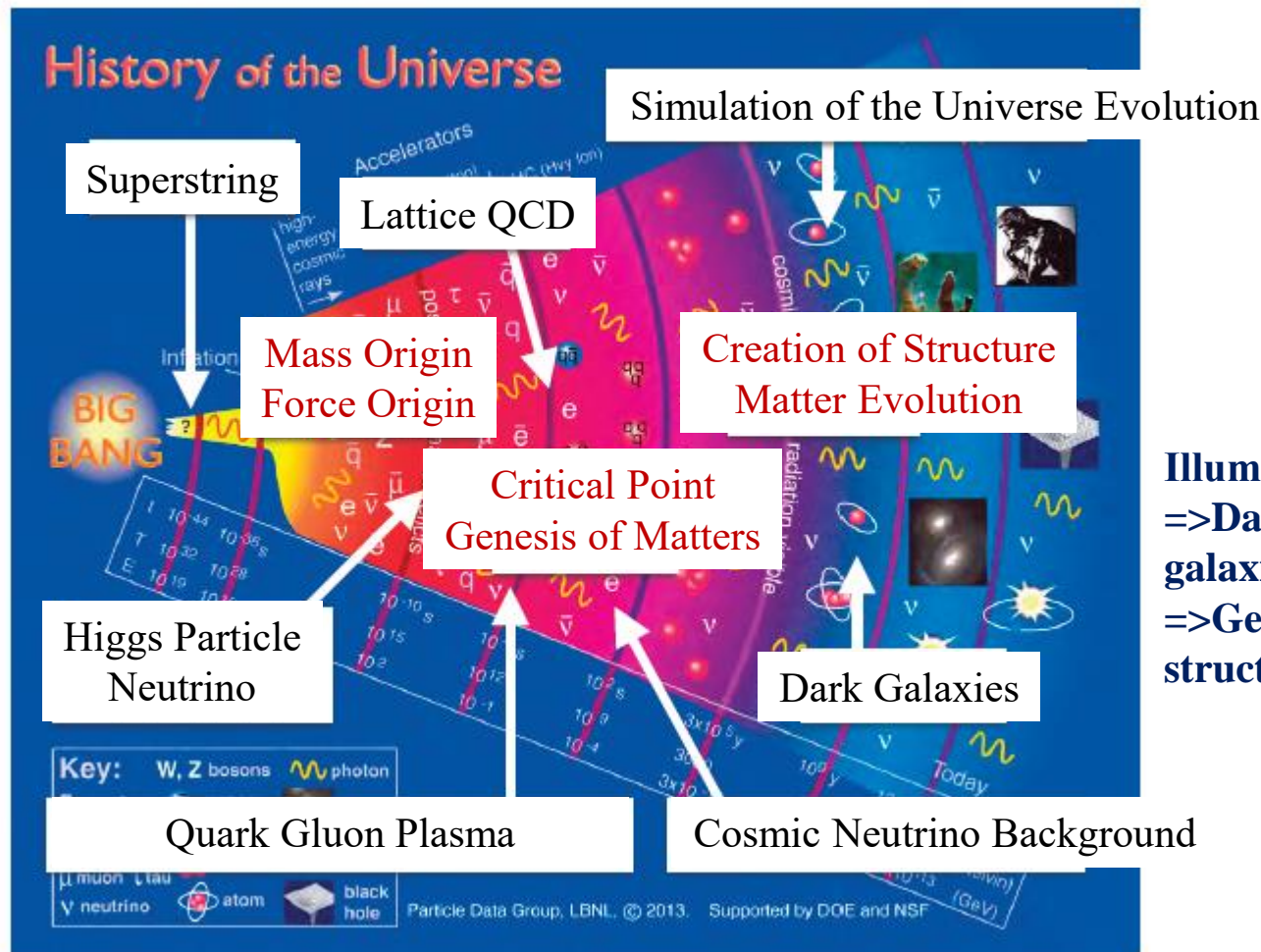
# Improvement of Hf-STJ Leakage Current



# Research Core for the History of the Universe

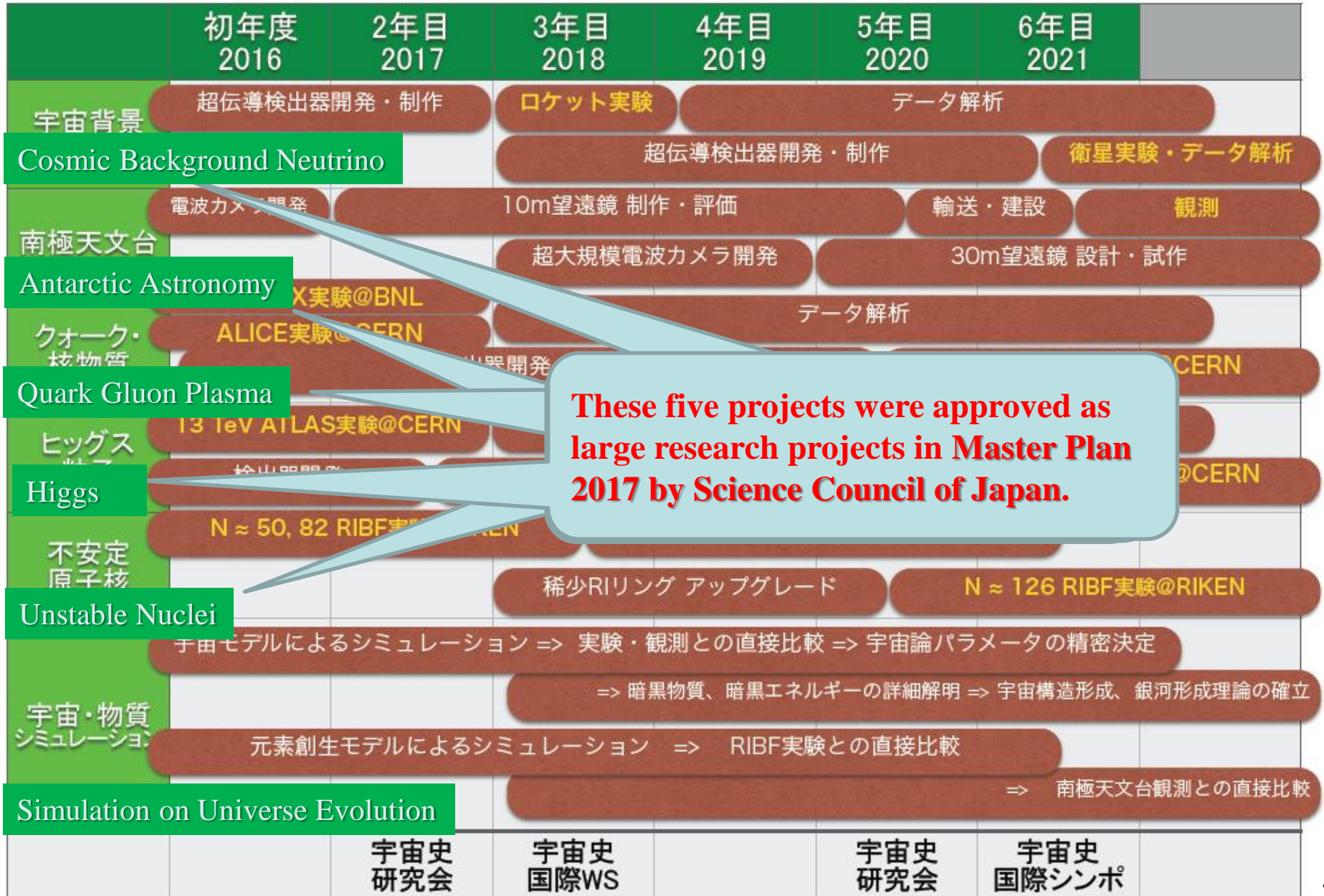
under Center for integrated Research in Fundamental Science and Engineering,  
University of Tsukuba (founded on Sep. 1, 2014)

**Mission:** coordinate the studies in elementary particles, quark nuclear matters and astrophysics to construct an integrated view of the History of the Universe.

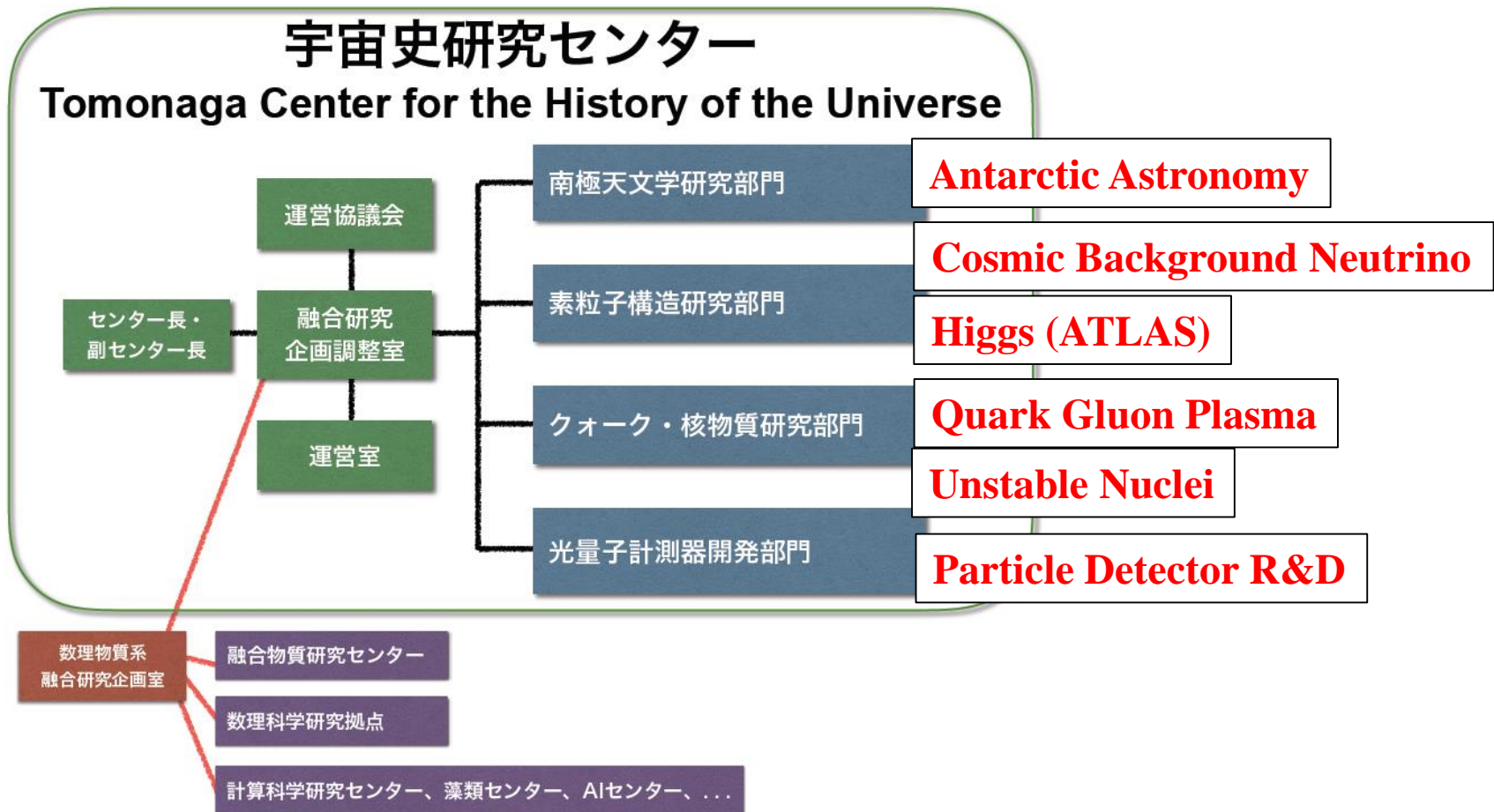
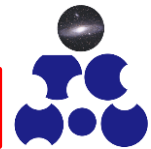


**Illuminate the “Darkness”:**  
=>Dark matter, Dark energy , Dark galaxies  
=>Genesis of matters, creation of structure and their evolution

# Schedule of Projects on the History of the Universe



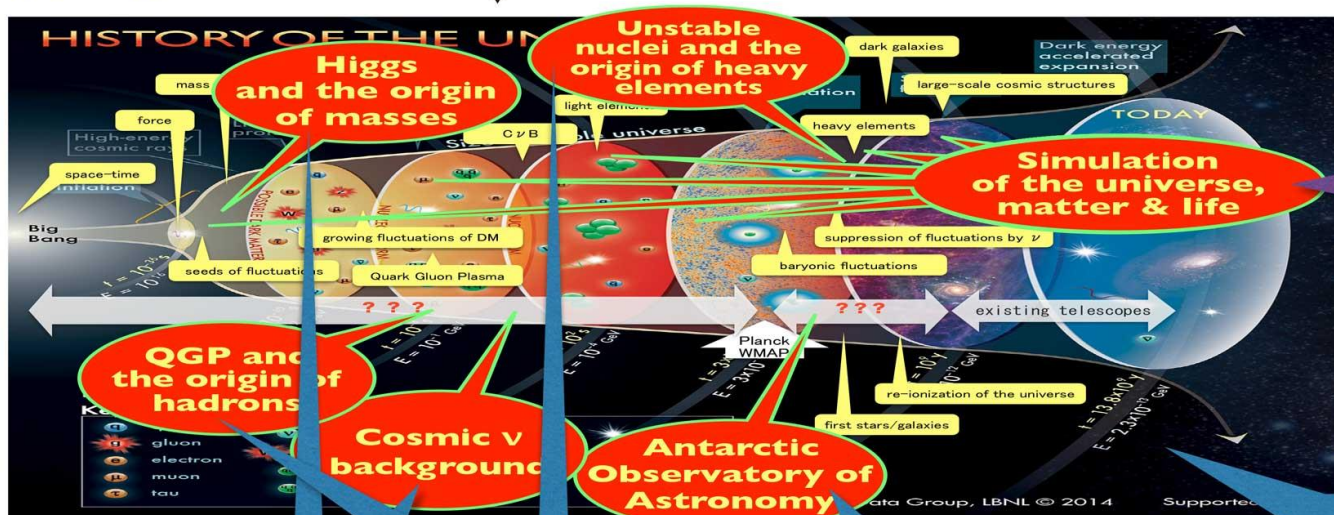






# Tomonaga Center for the History of the Universe

From Oct. 2017



**Center for Computational Sciences**

JICFuS JCAHPC

**Division of Elementary Particles**

- Higgs precision study by ATLAS experiment
- Detection of cosmic neutrino background by COBAND experiment
- Theory of quantum gravity and superstrings

**Division of Quark Nuclear Matters**

- QGP study by ALICE and STAR experiments
- Study of unstable nuclei by RI-beam factory
- QCD simulations towards QGP and nuclear matters

**Division of Antarctic Astrophysics**

- Construction of Antarctic Observatory of Astronomy towards dark galaxies
- Simulation of universe and galaxies

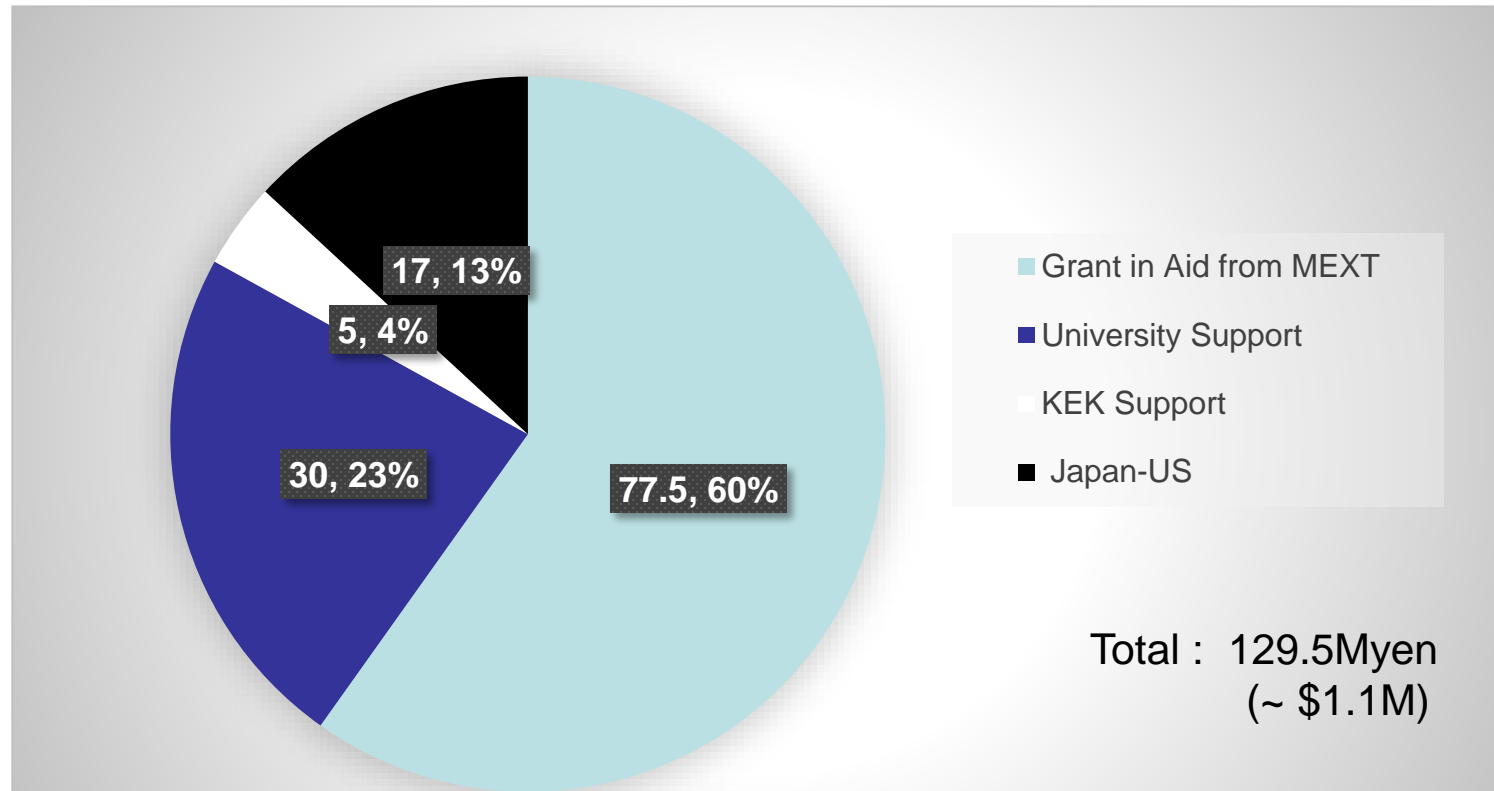
**Division of Photon and Particle Detectors**

- Development of PPD by superconducting and SOI technologies

Tsukuba Research Center for Energy Materials Science  
TIA-ACCELERATE, ...

**Approved for October 2017 to March 2022 (10,500kYen/year).  
Budget for Postdocs, Detector R&D and travel expenses.**

# R&D Budgets for the COBAND in the period of JFY2013-2017



Grant in Aid “Neutrino” from MEXT	R&D of Superconducting Infrared Detector for Neutrino Decay Experiment	JFY2013—2017	77.5 MYen
University Support	$^3\text{He}$ Sorption Refrigerator	JFY2013—2017	30.0 MYen
KEK Support	Cooperative Project of University of Tsukuba and KEK	JFY2013-2016	5.0 MYen
Japan-US Project	Neutrino Decay Search Experiment Cooperative	JFY2012-2014	17.0 MYen



# Summary

- R&D of STJ detectors and the design of the COBAND rocket experiment are underway.
  - Nb/Al-STJ satisfied our requirement for leakage current less than 100pA.
  - Cryogenic amplifier with the SOI technology worked at 300mK.

**We have succeeded in amplifying the STJ signal with the SOI cryogenic amplifier.**

  - **Hf-STJ signal for visible laser light was observed.**
- Many applications of the STJ detector as
  - a single photon detector in the far-infrared range,
  - a very low energy particle detector,
  - X-ray energy measurement with very higher energy resolution and so on.

**People who will join COBAND collaboration are very welcome.**

**Please join us.**

BACKUP

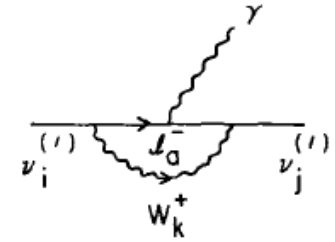
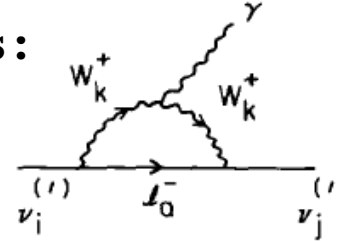
# Neutrino Lifetime by Left-Right Symmetric Model

In the Left - Right Symmetric Model  $SU(2)_L \otimes SU(2)_R \otimes U(1)$   
(PRL 38, 1252(1977), PRD 17, 1395(1978) NP B206, 359(1982)),

there are two Weak Boson mass eigenstates :

$$W_1 = W_L \cos \zeta - W_R \sin \zeta,$$

$$W_2 = W_L \sin \zeta + W_R \cos \zeta.$$



$W_L$  and  $W_R$  are fields with pure V-A and V+A couplings, respectively, and  $\zeta$  is a mixing angle.

$$\tau^{-1} = \frac{\alpha G_F^2}{128\pi^4} \left( \frac{m_3^2 - m_2^2}{m_3} \right)^3 \times |U_{32}|^2 |U_{33}|^2 \left[ \frac{9}{64} (m_3^2 + m_2^2) \frac{m_\tau^4}{M_{W1}^4} \left( 1 + \frac{M_{W1}^2}{M_{W2}^2} \right)^2 + 4m_\tau^2 \left( 1 - \frac{M_{W1}^2}{M_{W2}^2} \right)^2 \sin^2 2\zeta \right],$$

where  $\alpha$  is a fine structure constant,  $G_F$  is a Fermi coupling constant,  $m_\tau$ ,  $M_{W1}$  and  $M_{W2}$  are masses of  $\tau$ ,  $W_1$  and  $W_2$ , respectively.<sup>21,22)</sup>  $U_{ij}$  is the (i, j)-th element of the Maki-Nakagawa-Sakata mixing matrix<sup>23)</sup> and we took  $|U_{32}| = 1/\sqrt{2}$  and  $|U_{33}| = 1/\sqrt{2}$ .

$$\tau^{-1} = \frac{\alpha G_F^2}{128\pi^4} \left( \frac{\Delta m_{32}^2}{m_3} \right)^3 m_\tau^2 \sin^2 2\zeta$$

Using a lower mass limit  $M(W_R) > 715 \text{ GeV}$ , a mixing angle limit  $\zeta < 0.02$ , and  $m_3 = 50 \text{ meV}$ ,

$$\tau = 1.5 \times 10^{17} \text{ year}$$

Measured neutrino lifetime limit  $\tau < 3 \times 10^{12} \text{ year}$  from CIB results measured by COBE and AKARI

# Other papers citing our JPSJ paper

PHYSICAL REVIEW D 88, 013019 (2013)

## Radiative decays of cosmic background neutrinos in extensions of the MSSM with a vectorlike lepton generation

Amin Aboubrahim,<sup>2,\*</sup> Tarek Ibrahim,<sup>1,2,†,§,||</sup> and Pran Nath<sup>3,4,‡,¶</sup>

<sup>1</sup>*Department of Physics, Faculty of Science, University of Alexandria, Alexandria 21511, Egypt*

<sup>2</sup>*Department of Physics, Faculty of Sciences, Beirut Arab University, Beirut 11-5020, Lebanon*

<sup>3</sup>*Department of Physics, Northeastern University, Boston, Massachusetts 02115-5000, USA*

<sup>4</sup>*KITP, University of California, Santa Barbara, California 93106-4030, USA*

(Received 11 June 2013; published 30 July 2013)

An analysis of radiative decays of the neutrinos  $\nu_j \rightarrow \nu_l \gamma$  is discussed in minimal supersymmetric standard model extensions with a vector like lepton generation. Specifically we compute neutrino decays arising from the exchange of charginos and charged sleptons where the photon is emitted by the charged particle in the loop. It is shown that while the lifetime of the neutrino decay in the Standard Model is  $\sim 10^{43}$  yrs for a neutrino mass of 50 meV, the current lower limit from experiment from the analysis of the Cosmic Infrared Background is  $\sim 10^{12}$  yrs and thus beyond the reach of experiment in the foreseeable future. However, in the extensions with a vectorlike lepton generation the lifetime for the decays can be as low as  $\sim 10^{12} - 10^{14}$  yrs and thus within reach of future improved experiments. The effect of  $CP$  phases on the neutrino lifetime is also analyzed. It is shown that while both the magnetic and the electric transition dipole moments contribute to the neutrino lifetime, often the electric dipole moment dominates even for moderate size  $CP$  phases.

MSSM extension with a vectorlike lepton generation

$$\rightarrow \tau_\nu \sim 10^{12} \sim 10^{14} \text{ years}$$

# MSSM extension model with a vectorlike lepton generation

PHYSICAL REVIEW D 88, 013019 (2013)

## Radiative decays of cosmic background neutrinos in extensions of the MSSM with a vectorlike lepton generation

Amin Aboubrahim,<sup>2,\*</sup> Tarek Ibrahim,<sup>1,2,†,§</sup> and Pran Nath<sup>3,4,‡,¶</sup>

PHYSICAL REVIEW D 89, 055009 (2014)

Large neutrino magnetic dipole moments in MSSM extensions

Amin Aboubrahim,<sup>2,\*</sup> Tarek Ibrahim,<sup>1,2,†</sup> Ahmad Itani,<sup>2,‡</sup> and Pran Nath<sup>3,§</sup>

$$SU(3)_C \times SU(2)_L \times U(1)_Y \quad Q = T_3 + Y$$

$$\psi_{iL} = \begin{pmatrix} \nu_{iL} \\ l_{iL} \end{pmatrix} \sim \left(1, 2, -\frac{1}{2}\right), \quad l_{iL} \sim (1, 1, 1),$$

$$\nu_{iL}^c \sim (1, 1, 0), \quad i = 1, 2, 3 \quad (3)$$

$$\chi^c = \begin{pmatrix} E_L^c \\ N_L^c \end{pmatrix} \sim \left(1, 2, \frac{1}{2}\right), \quad E_L \sim (1, 1, -1), \quad \text{vectorlike lepton generation}$$

$$N_L \sim (1, 1, 0). \quad V+A^{(4)} \text{ interaction}$$

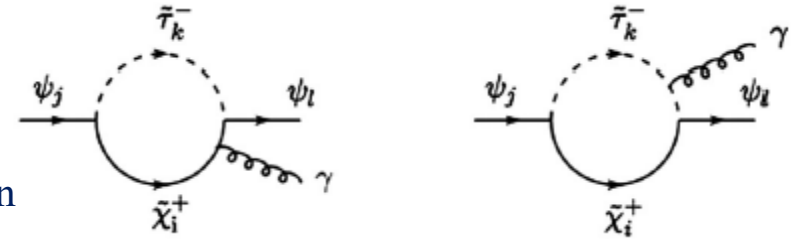
$$\begin{pmatrix} \nu_{\tau R} \\ N_R \\ \nu_{\mu R} \\ \nu_{e R} \end{pmatrix} = D_R^\nu \begin{pmatrix} \psi_{1R} \\ \psi_{2R} \\ \psi_{3R} \\ \psi_{4R} \end{pmatrix}, \quad \begin{pmatrix} \nu_{\tau L} \\ N_L \\ \nu_{\mu L} \\ \nu_{e L} \end{pmatrix} = D_L^\nu \begin{pmatrix} \psi_{1L} \\ \psi_{2L} \\ \psi_{3L} \\ \psi_{4L} \end{pmatrix}. \quad (17)$$

In Eq. (16)  $\psi_1, \psi_2, \psi_3, \psi_4$  are the mass eigenstates for the neutrinos, where in the limit of no mixing we identify  $\psi_1$  as the tau neutrino,  $\psi_2$  as the heavier mass eigenstate,  $\psi_3$  as the muon neutrino and  $\psi_4$  as the electron neutrino. To make contact with the normal neutrino hierarchy we relabel the states so that

$$\nu_1 = \psi_4, \quad \nu_2 = \psi_3, \quad \nu_3 = \psi_1, \quad \nu_4 = \psi_2, \quad (18)$$

which we assume has the mass hierarchical pattern

$$m_{\nu_1} < m_{\nu_2} < m_{\nu_3} < m_{\nu_4}. \quad (19)$$



Neutrino mass eigenvalues (GeV)		$m_{\nu_3} = 5.2 \times 10^{-11}$ $m_{\nu_2} = 9.2 \times 10^{-12}$ $m_{\nu_1} = 9.7 \times 10^{-13}$
(i) $m_{\chi^\pm} = 256$ GeV $m_{\tilde{\tau}} = 162$ GeV	$\mu_2$ $\mu_1$ $\nu_3$ lifetime	$1.2 \times 10^{-10}$ $2.5 \times 10^{-13}$ $3.9 \times 10^{14}$ yr
(ii) $m_{\chi^\pm} = 267$ GeV $m_{\tilde{\tau}} = 202$ GeV	$\mu_2$ $\mu_1$ $\nu_3$ lifetime	$4.6 \times 10^{-10}$ $1.3 \times 10^{-12}$ $2.5 \times 10^{14}$ yr
(iii) $m_{\chi^\pm} = 268$ GeV $m_{\tilde{\tau}} = 158$ GeV	$\mu_2$ $\mu_1$ $\nu_3$ lifetime	$2.2 \times 10^{-10}$ $1.1 \times 10^{-13}$ $1.8 \times 10^{14}$ yr
(iv) $m_{\chi^\pm} = 272$ GeV $m_{\tilde{\tau}} = 195$ GeV	$\mu_2$ $\mu_1$ $\nu_3$ lifetime	$-7.6 \times 10^{-10}$ $-1.3 \times 10^{-13}$ $8.8 \times 10^{13}$ yr

$$-\mathcal{L}_{CC} = \frac{g}{2\sqrt{2}} W_\rho^\dagger \{ \bar{\nu}_\tau \gamma^\rho (1 - \gamma_5) \tau + \bar{\nu}_\rho \gamma^\mu (1 - \gamma_5) \mu + \bar{\nu}_e \gamma^\rho (1 - \gamma_5) e + \bar{N} \gamma^\rho (1 + \gamma_5) E \} + H.c. \quad (24)$$



## Publications

1. "Development of Superconducting Tunnel Junction Photon Detectors with Cryogenic Preamplifier for COBAND experiment", S. H. Kim *et al.*, to be published in Proceeding of TIPP2017
2. "Development of Superconducting Tunnel Junction Detector using Hafnium for COBAND experiment"(Poster presentation), K.Takemasa *et al.*, , to be published in Proceeding of TIPP2017
3. "Development of Superconducting Tunnel Junction detectors as a far-infrared photon-by-photon spectrometer for neutrino decay search", Y. Takeuchi *et al.*, Instrumentation and Measurement Technology Conference (I2MTC), 2015 IEEE International, 551 - 555 (2015), DOI:10.1109/I2MTC.2015.7151327
4. "Development FD-SOI MOSFET amplifiers for integrated read-out circuit of superconducting-tunnel-junction single-photon-detectors" K. Kiuchi *et al.* , Proceedings of International Workshop on SOI Pixel Detector, FERMILAB-CONF-15-355-E-TD (2015) arXiv:1507.07424
5. "Development of Superconducting Tunnel Junction Detectors as a far-infrared single photon detector for neutrino decay search", Y. Takeuchi *et al.*, PoS(TIPP2014)155
6. "Development of Superconducting Tunnel Junction Photon Detector on SOI Preamplifier Board to Search for Radiative decays of Cosmic Background Neutrino", K. Kasahara *et al.*, PoS(TIPP2014)074
7. "Search for Cosmic Background Neutrino Decay", S. H. Kim *et al.*, JPS Conf. Proc. 1, 013127 (2014)
8. "Search for Radiative Decays of Cosmic Background Neutrino using Cosmic Infrared Background Energy Spectrum", S.H.Kim, K.Takemasa, Y.Takeuchi, and S.Matsuura, JPSJ 81 (2012) 024101, arXiv:1112.4568
9. "Development of superconducting tunnel junction photon detector using Hafnium", S.H.Kim , H.S. Jeong, K. Kiuchi, S. Kanai, T. Onjo, K. Takemasa, Y. Takeuchi, H. Ikeda, S. Matsuura, H. Sato, M. Hazumi, and S.B. Kim, Physics Procedia 37 (2012) 667-674 (Proceedings of TIPP2011)

# Cosmic Background Neutrino

*Fermi and Bose Distribution Function*

$$F(E) = \frac{1}{e^{(E-\mu)/kT} \pm 1}$$

where + for fermions and - for bosons, and E is energy and  $\mu$  is a chemical potential.

For  $\mu \ll T$  and  $m \ll T$ ,

$$\text{Energy density } \rho = g \int \frac{d^3p}{(2\pi)^3} E F(E) = g \left(\frac{7}{8}\right)^F \frac{\pi^2}{30} T^4$$

$$\text{Number density } n = g \int \frac{d^3p}{(2\pi)^3} F(E) = g \left(\frac{3}{4}\right)^F \frac{\zeta(3)}{\pi^2} T^3$$

$$\text{Entropy } s = \frac{4\rho}{3T} = g \left(\frac{7}{8}\right)^F \frac{2\pi^2}{45} T^3$$

*Temperature:*

Below 3MeV,  $\nu$  is decoupled from other particles because the weak interaction cross section becomes too small.

Below 1MeV,  $e^+e^- \rightarrow \gamma\gamma$  is possible, but  $\gamma\gamma \rightarrow e^+e^-$  is impossible. so photons are reheated by this process. The entropies before and after this time are equal to each other:

$$\text{Entropy } s \propto g \left(\frac{7}{8}\right)^F T^3 \quad g=2(\text{ for } \gamma), 2(\text{ for } e^- \text{ or } e^+), 1(\text{for } \nu \text{ or anti-}\nu)$$

where  $g$  is the spin degree of freedom, and  $F = 1$  ( for fermions) and 0 (for bosons).

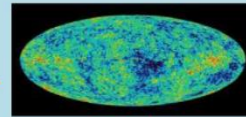
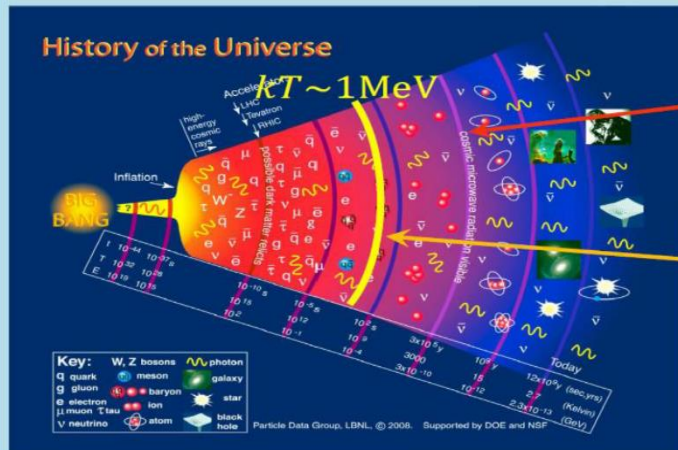
The present entropies of photons and neutrinos,  $s_{\gamma 0}$  and  $s_{\nu 0}$  are given by

$$s_{\gamma 0} = a^3 (s_{\gamma} + s_{e^- + e^+}) , \quad s_{\nu 0} = a^3 s_{\nu} \quad \text{where } a \text{ is a scale factor.}$$

$$\rightarrow \frac{s_{\nu 0}}{s_{\gamma 0}} = \frac{s_{\nu}}{s_{\gamma} + s_{e^- + e^+}} = \frac{2 \times \frac{7}{8}}{2 + 4 \times \frac{7}{8}} = \frac{7}{22} \quad \therefore s_{\nu 0} = \frac{7}{22} s_{\gamma 0}$$

$$2 \times \frac{7}{8} T_{\nu}^3 = \frac{7}{22} \times 2 T_{\gamma}^3 \rightarrow T_{\nu} = \left(\frac{4}{11}\right)^{\frac{1}{3}} T_{\gamma} \quad \text{As } T_{\gamma} = 2.73K, \quad \therefore T_{\nu} = 1.95K$$

# Cosmic Background Neutrino



**CMB**

$$n_\gamma = 411/\text{cm}^3$$

$$T_\gamma = 2.73 \text{ K}$$

**CvB**

$$n_\nu = n_{\bar{\nu}} = \frac{3}{4} \left( \frac{T_\nu}{T_\gamma} \right)^3 \frac{n_\gamma}{2}$$

$$= 56/\text{cm}^3$$

$$T_\nu = \left( \frac{4}{11} \right)^{\frac{1}{3}} T_\gamma = 1.95 \text{ K}$$

$$E_\gamma = \frac{m_3^2 - m_2^2}{2m_3}$$



Temperature:

$$T_\nu = 1.95 \text{ K}$$

Number density:

$$\text{As } \mu/T \ll 1, \quad n \propto g \left( \frac{3}{4} \right)^F T^3$$

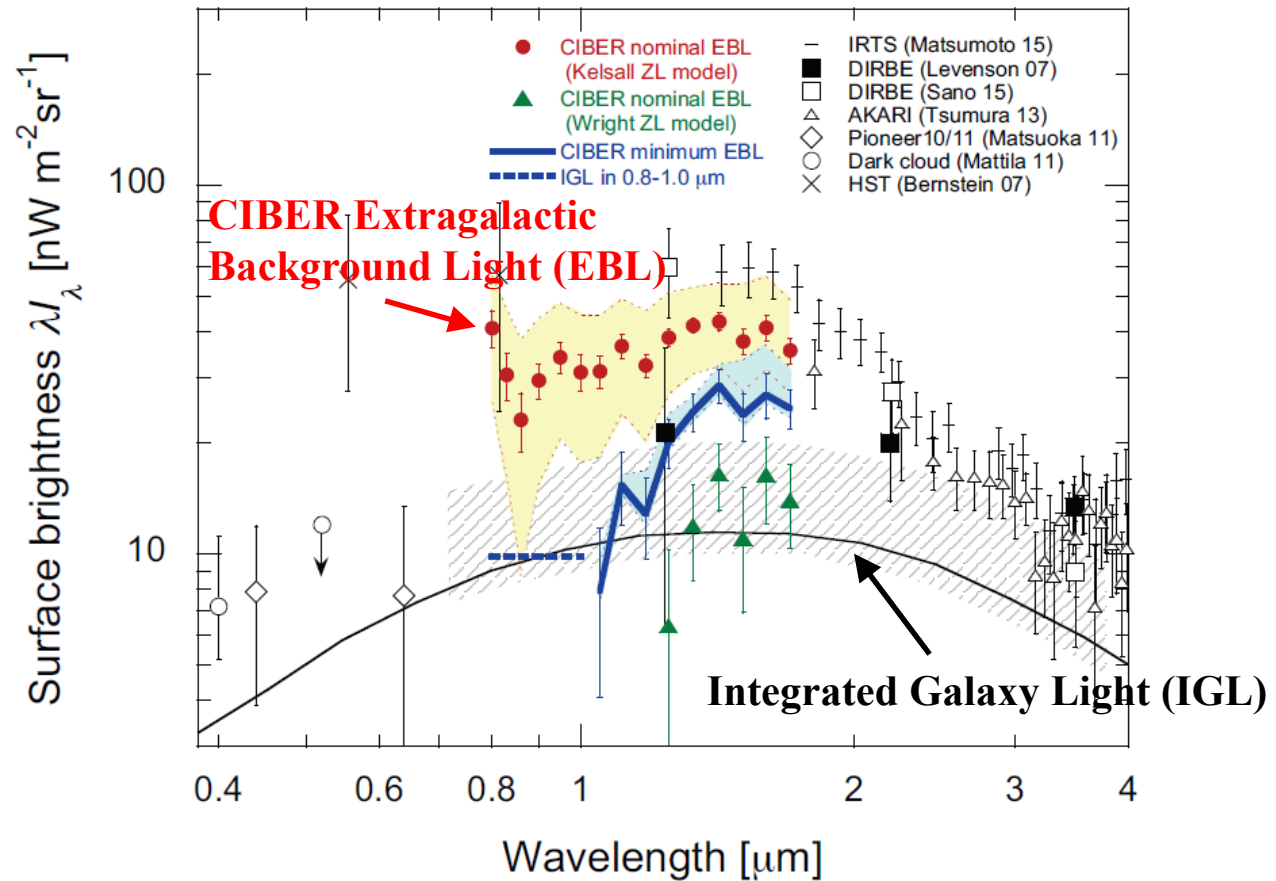
where  $g$  is the spin degree of freedom, and  $F = 1$  ( for fermions) and 0 (for bosons).

$$\rightarrow n_\nu = \frac{3}{4} \left( \frac{T_\nu}{T_\gamma} \right)^3 \frac{n_\gamma}{2}$$

$$\therefore n_{\nu_\alpha} \approx n_{\bar{\nu}_\alpha} \approx 56 \text{ cm}^{-3} \quad (\alpha = e, \mu, \tau)$$

# Near-Infrared Photon Energy Spectrum from Outer Space

April 2017 (Press Release): CIBER rocket experiment measured the energy spectrum of the cosmic infrared background (CIB) in the near-infrared region. They found an excess of the CIB spectrum over the prediction by integrated galaxy light (IGL).



This excess might be due to first-generation stars Ly- $\alpha$ .

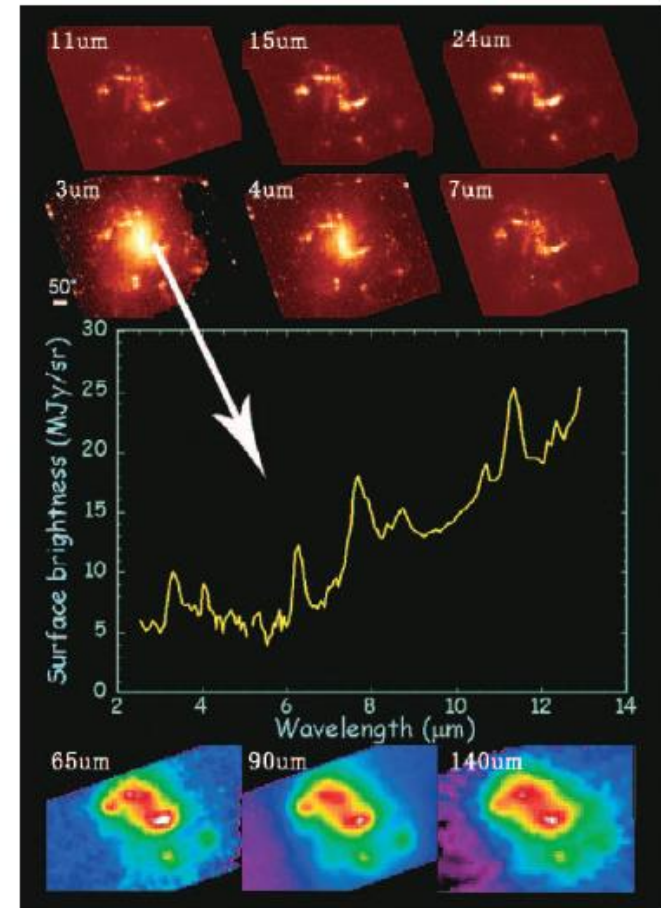
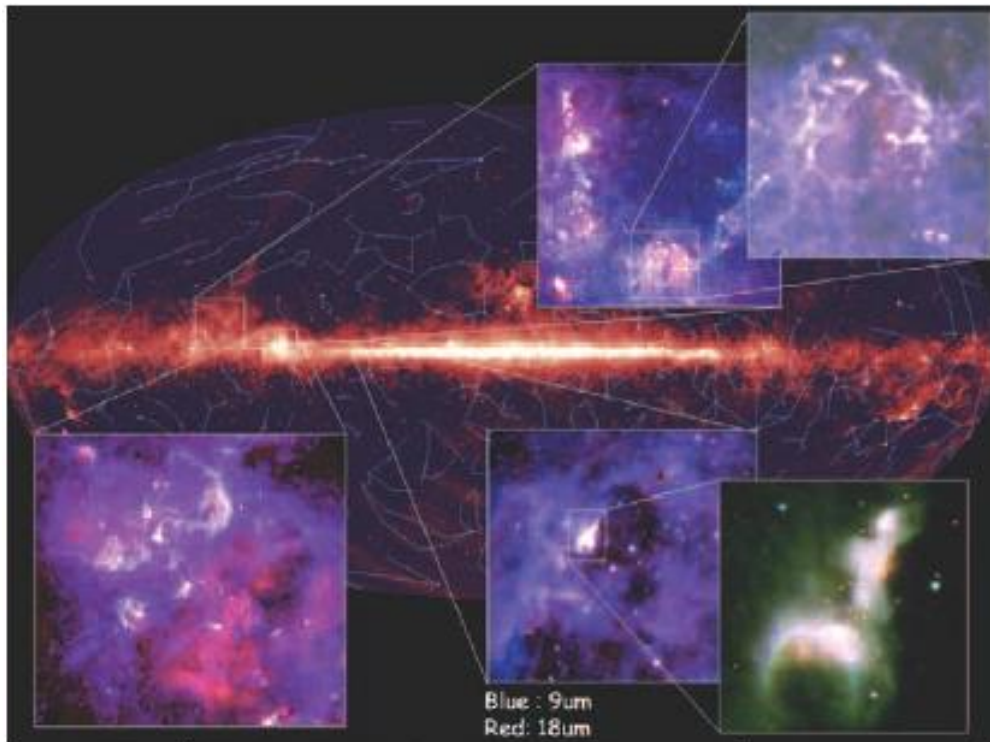
CIBER collaboration plans to do CIBER-2 rocket experiment which has 10 times larger sensitivity.

# Astrophysics with NIR and FIR by AKARI



- NIR and FIR from interstellar gas → Formation and evolution of galaxies, stars and dusts
- NIR: Search for first-generation stars using Ly- $\alpha$  line.
- FIR: Search for first-generation heavy element dusts.

Map of NIR from Milky Way Galaxy measured by AKARI



Molecular cloud IR in the arm of NGC1313 measured by AKARI



# Existing FIR photo-detectors

Detectors	$\lambda(\mu\text{m})$	Operation Temp.	NEP ( $\text{W}/\text{Hz}^{1/2}$ )	
Monolithic Ge:Ga	50-110	2.2K	$\sim 10^{-17}$	Akari-FIS
Stressed Ge:Ga	60-210	0.3K	$\sim 0.9 \times 10^{-17}$	Herschel-PACS

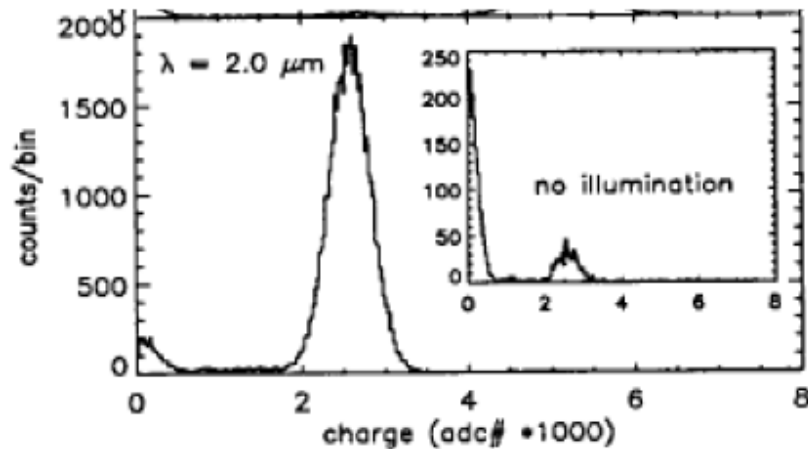
Need **more than 2 orders improvement** from existing photoconductor-based detectors

We are trying to achieve  $\text{NEP} \sim O(10^{-20}) \text{ W}/\sqrt{\text{Hz}}$  **by using**

- **Superconducting tunneling junction detector**
- **FIR single-photon counting technique**

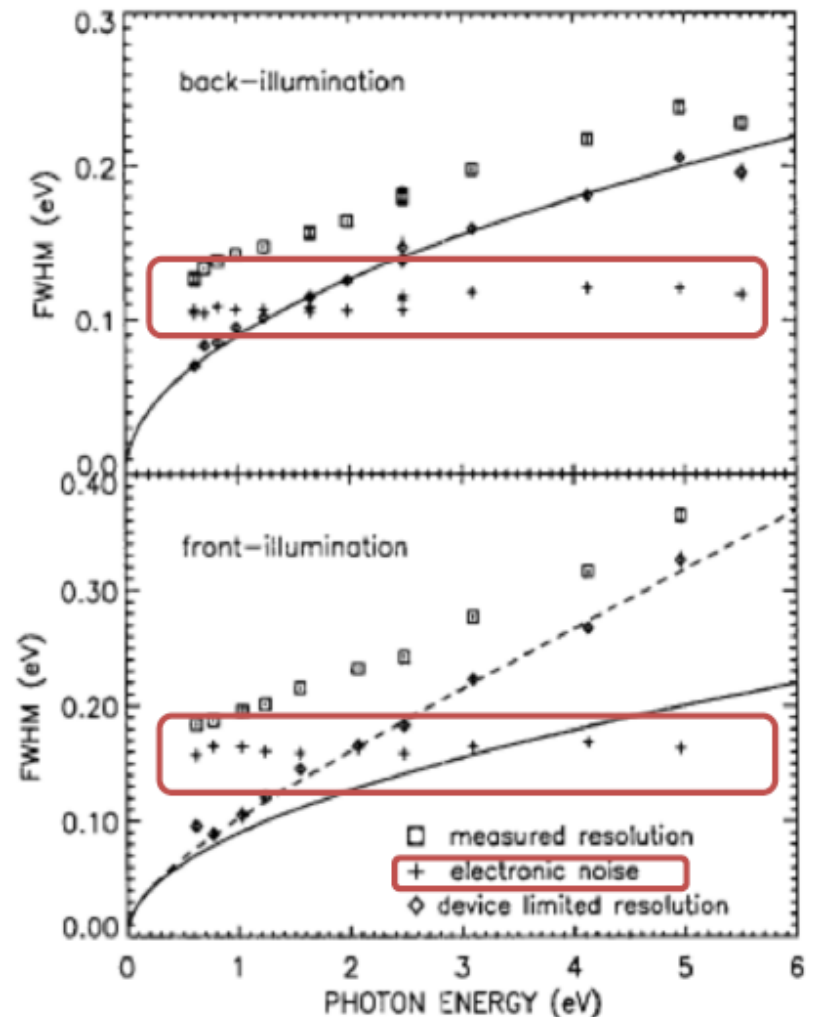


# STJ Energy Resolution for Near-Infrared Photon



P. Verhoeve et. al 1997

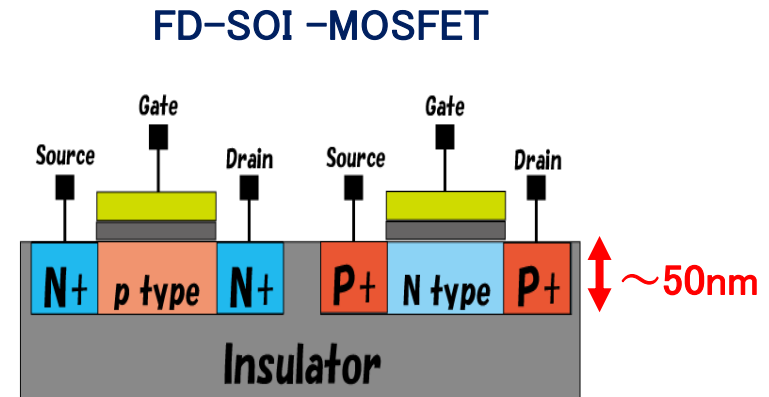
- 30 $\mu\text{m}$  sq. Ta/Al-STJ
- $\Delta E \sim 130 \text{ meV}$  @  $E = 620 \text{ meV} (\lambda = 2 \mu\text{m})$
- Charge sensitive amplifier at room temp.
- Electronic noise  $\sim 100 \text{ meV}$



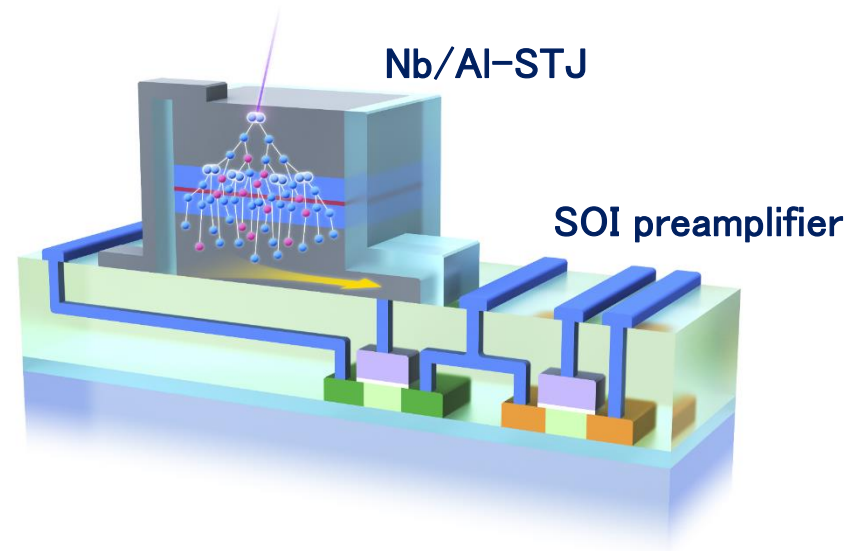
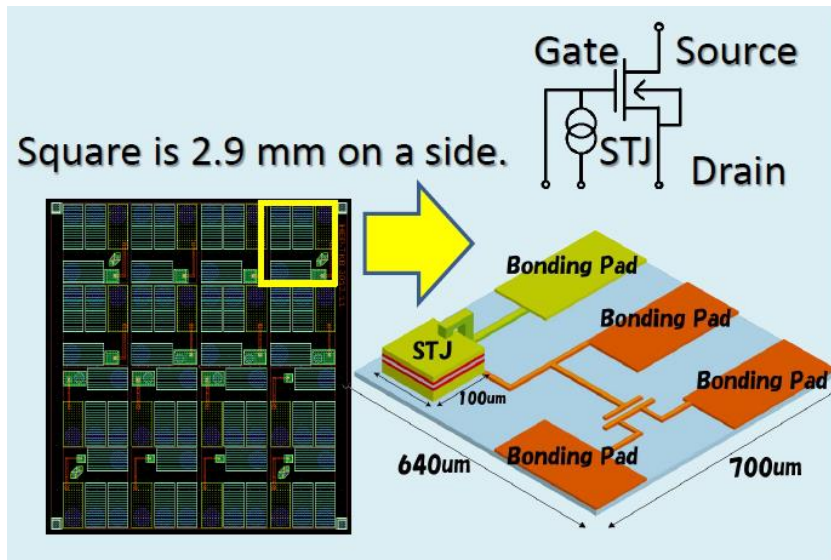
In sub-eV  $\sim$  several-eV region, STJ gives the best energy resolution among superconductor based detectors, but limited by readout electronic noise.

# R&D of SOI-STJ Detector

FD-SOI (Fully Depleted Silicon-On-Insulator) device was proved to operate at 4K by a JAXA/KEK group (AIPC 1185,286-289(200 FD-SOI 9)). It has the following characteristics: low-power consumption, high speed, easy large scale integration and suppression of charge-up by high mobility carrier due to thin depletion layer( $\sim 50\text{nm}$ ).

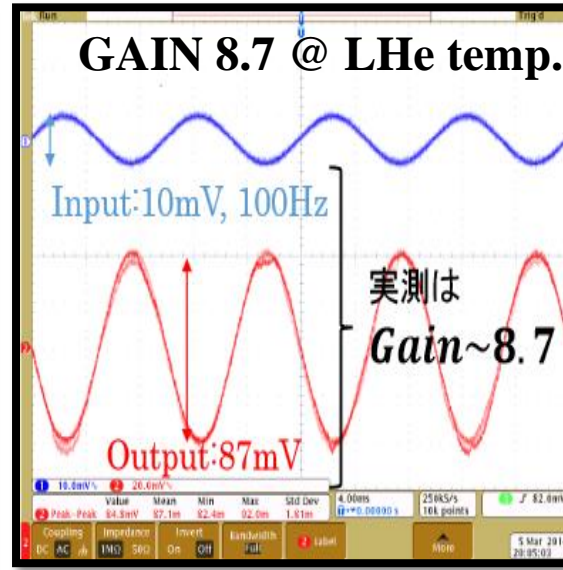
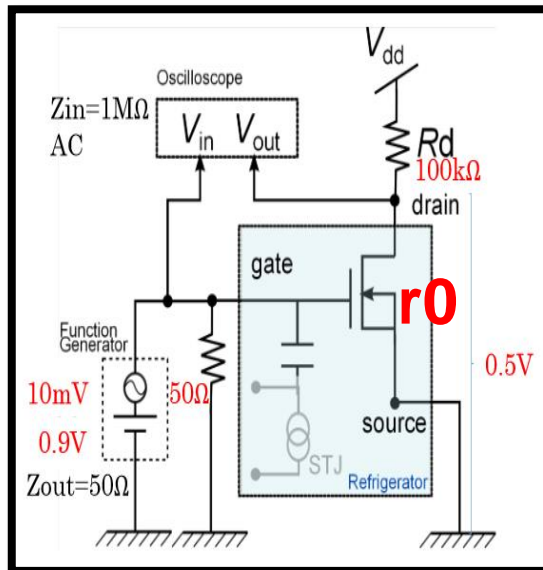
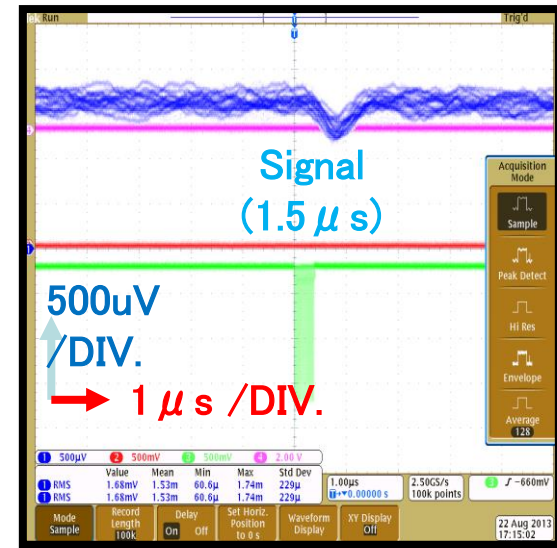


To improve the **signal-to-noise ratio** and to make **multi-pixel device** easily, we made a SOI-STJ detector where we processed Nb/Al-STJ on a SOI transistor board.



# Performance of STJ and SOIFET in SOI-STJ detector

- We observed the signal of Nb/Al-STJ processed on the SOI board to 465nm laser pulse at 700mK.

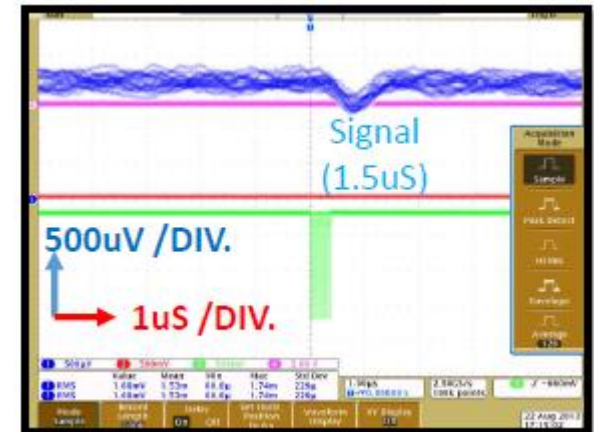
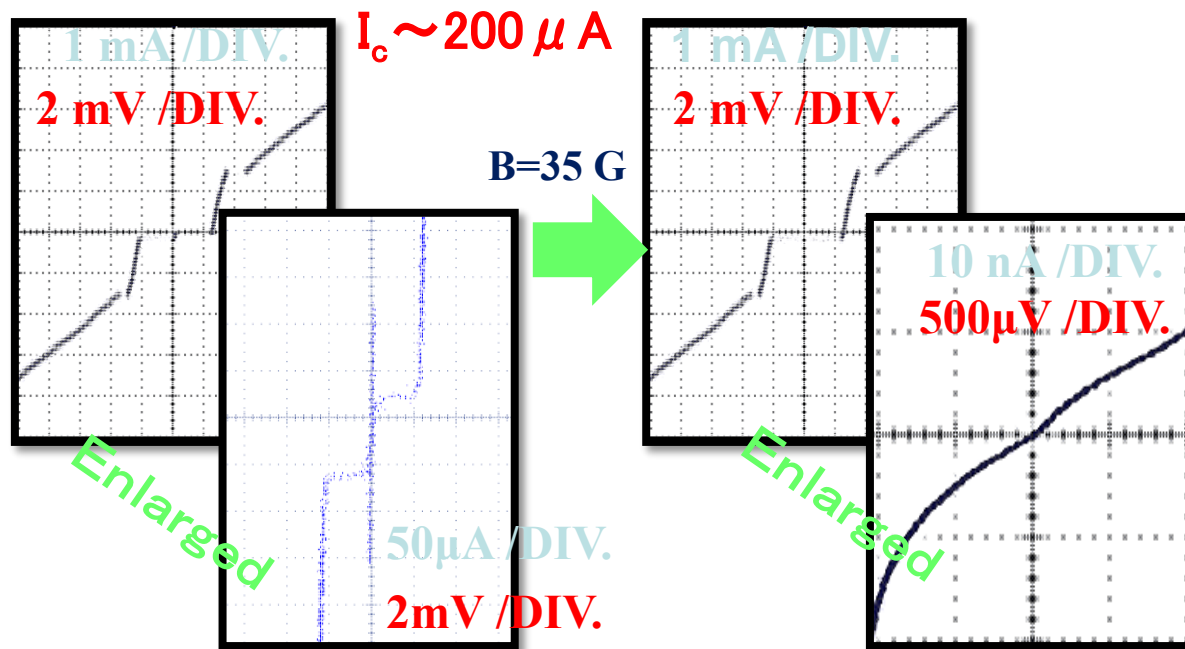
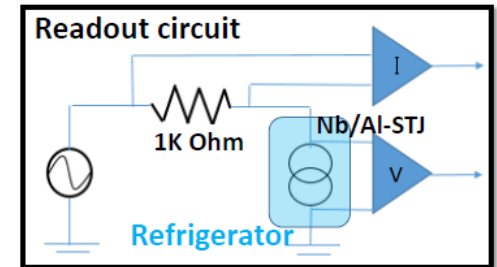


- We confirmed that the SOI-FET work as a preamplifier with a gain of 8.7 at 4K up to 100kHz.

# Performance of Nb/Al-STJ in SOI-STJ Detector

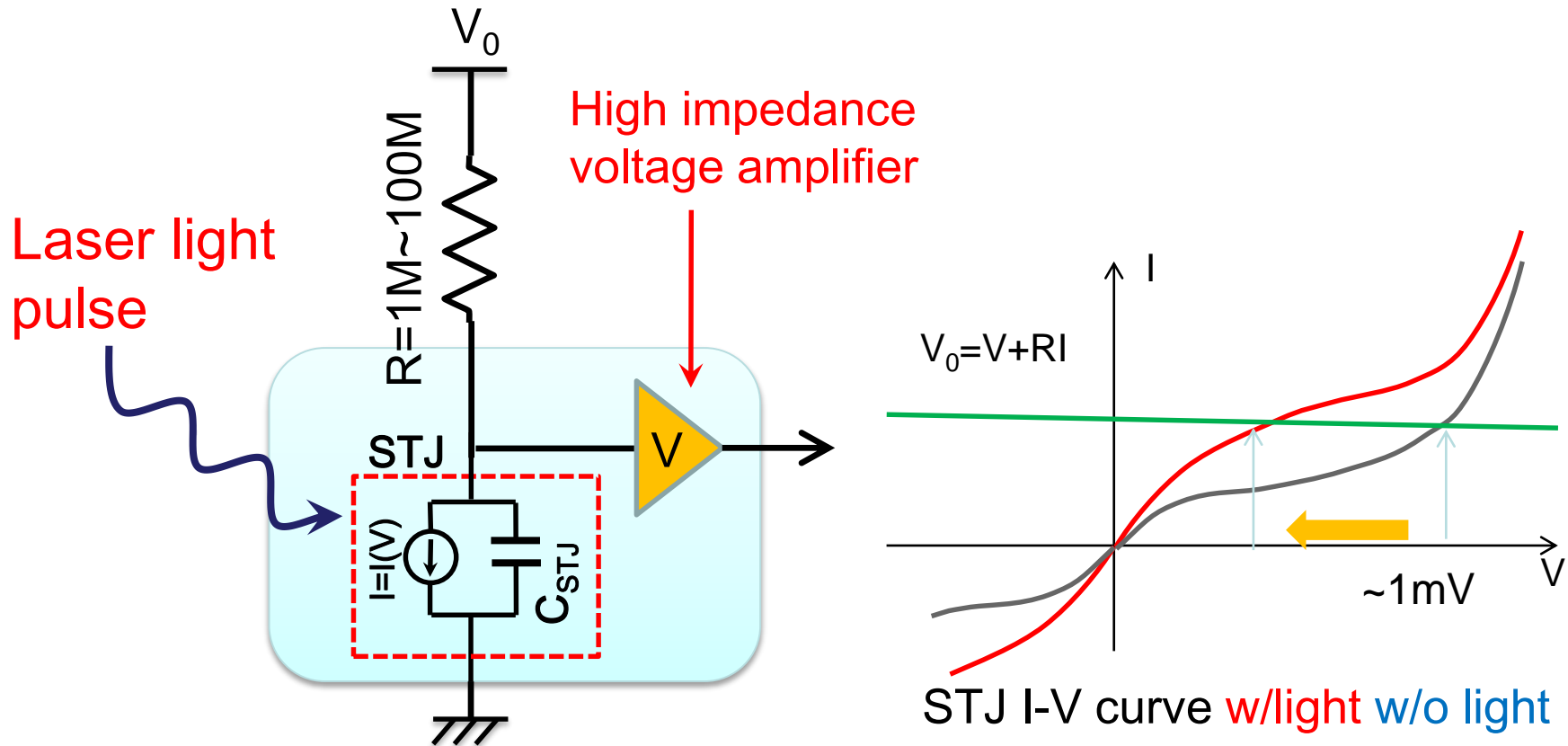
We measured the I-V curve of the Nb/Al-STJ ( $50 \times 50 \mu\text{m}^2$  junction) processed on the SOI wafer at **700mK** with a dilution refrigerator.

## □ I-V curve of Josephson Junction



- Quality Factor  $(R_{\text{dynamic}}/R_{\text{normal}})$ 
  - On Si wafer :  $5 \times 10^5$
  - On SOI wafer :  $3 \times 10^5$

# STJ response at a constant current mode



- STJ has a large resistance in series.  
→ constant current mode
- STJ response to light is observed as a voltage drop.

# **COBAND Collaboration**

## **( Cosmic Background Neutrino Decay Search )**

A part of the consortium of the History of the Universe

**Seoul National Univ.**  
**S. B. Kim**  
STJ detector

**Fukui Univ.**  
**T. Yoshida**  
FIR photon  
beam source

**Kindai  
Univ.**  
**Y. Kato**  
Data transfer

**FNAL**  
**E. Ramberg**  
Electronics

**KEK**  
**Y. Arai, M. Hazumi**  
Electronics

**AIST**  
**M. Ohkubo, M. Ukibe**  
SOI-STJ detector

**Univ. of Tsukuba**  
**S. H. Kim, Y. Takeuchi**  
STJ detector,  
Electronics,  
Cryostat, Optical system

**JAXA/ISAS**  
**T. Wada, H. Ikeda**  
Rocket, Electronics

**Kwansei Gakuin Univ.**  
**S. Matsuura**  
Cryostat, Optical system

**RIKEN**  
**S. Mima**  
STJ detector

**Okayama  
Univ.**  
**H. Ishino**  
STJ detector

**Shizuoka  
Univ.**  
**S. Kawahito**  
Electronics

**2015 Shizuoka Univ.**  
**Kwansei Gakuin  
Univ.**

**2014 AIST**

**2011 FNAL, Okayama Univ., Fukui Univ., Kindai Univ.**

**2007 Univ. of Tsukuba, JAXA/ISAS, RIKEN, KEK, Seoul National Univ.**

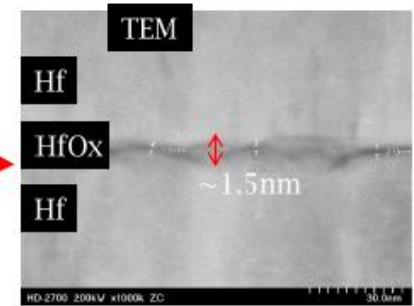
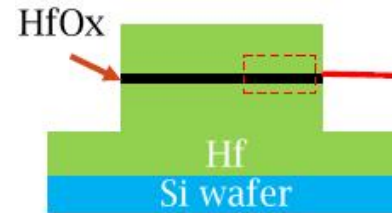
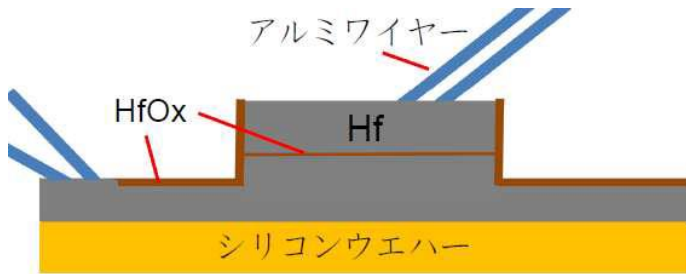


# R&D Status of Hf-STJ

Goal: Measure energy of a single far-infrared photon for neutrino decay search experiment within 2% energy resolution.

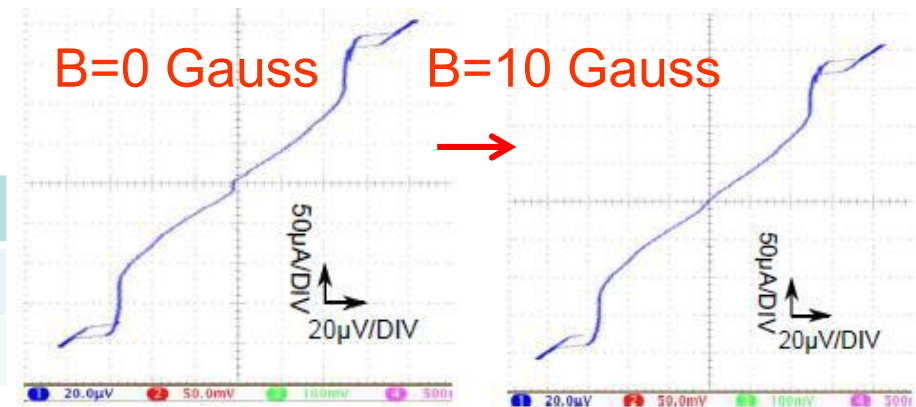
Micro-calorimeter: Hf-STJ can generate enough quasi-particles from cooper pair breakings to achieve 2% energy resolution for photons with  $E_\gamma = 25\text{meV}$ .

## Direct wire bonding on Hf layer



## I-V curve of Hf-STJ ( $100 \times 100 \mu\text{m}^2$ )

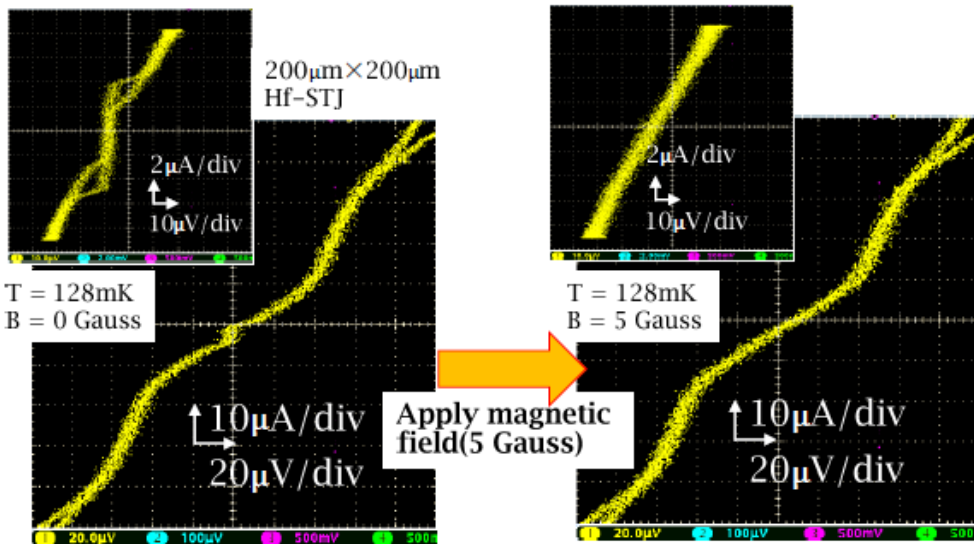
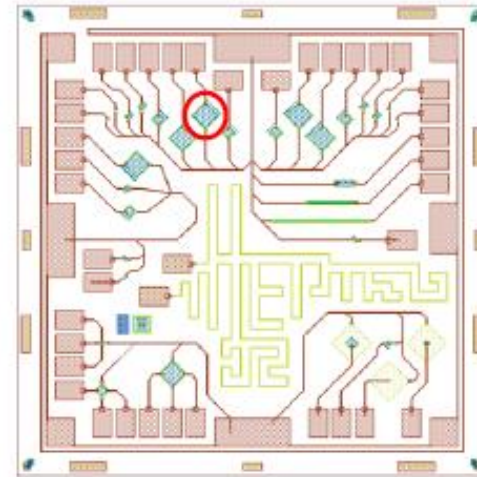
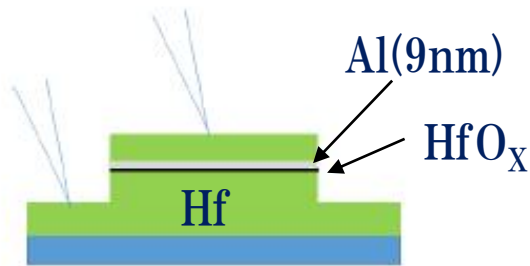
- $T \sim 40\text{mK}$ ,  $I_c = 10 \mu\text{A}$ ,  $R_d = 0.6 \Omega$



STJ size	# of samples	$R_d$
$200 \times 200 \mu\text{m}^2$	3	$0.22 \pm 0.01 \Omega$
$100 \times 100 \mu\text{m}^2$	3	$0.60 \pm 0.10 \Omega$

# Results of Hf/Al-STJ R&D

In 2016, we made a thin aluminum layer (9nm) on the HfO layer (1-2 nm) to improve the insulation of the HfO<sub>x</sub> layer. Hf/Al/HfO<sub>x</sub>/Hf-STJ



$$I_{\text{leak}} = 5 \mu\text{A at } 10 \mu\text{V.}$$

$$R_d = 2 \Omega$$

Dynamic Resistance was improved by a factor of 15 in Hf/Al/HfO<sub>x</sub>/Hf-STJ.

# Improvement of Hf-STJ Leakage Current

- Hf-STJ was proved to work as a photon detector.
- To satisfy our leakage current requirement for COBAND experiment, we need decrease the leakage current of Hf-STJ.

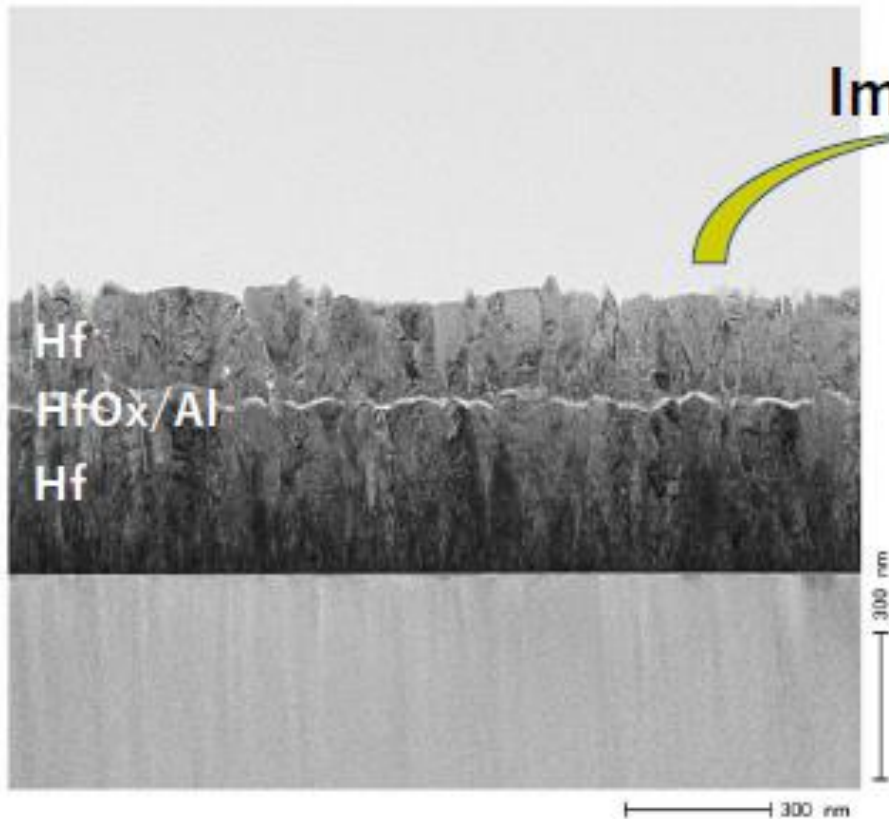
We are working on this improvement by the following methods:

- Optimize the Hf sputtering conditions to make Hf surface smoother. Because the imperfect insulator layer causes high leakage current.
- Downsize Hf-STJ. Because the leakage current is proportional to the junction size.

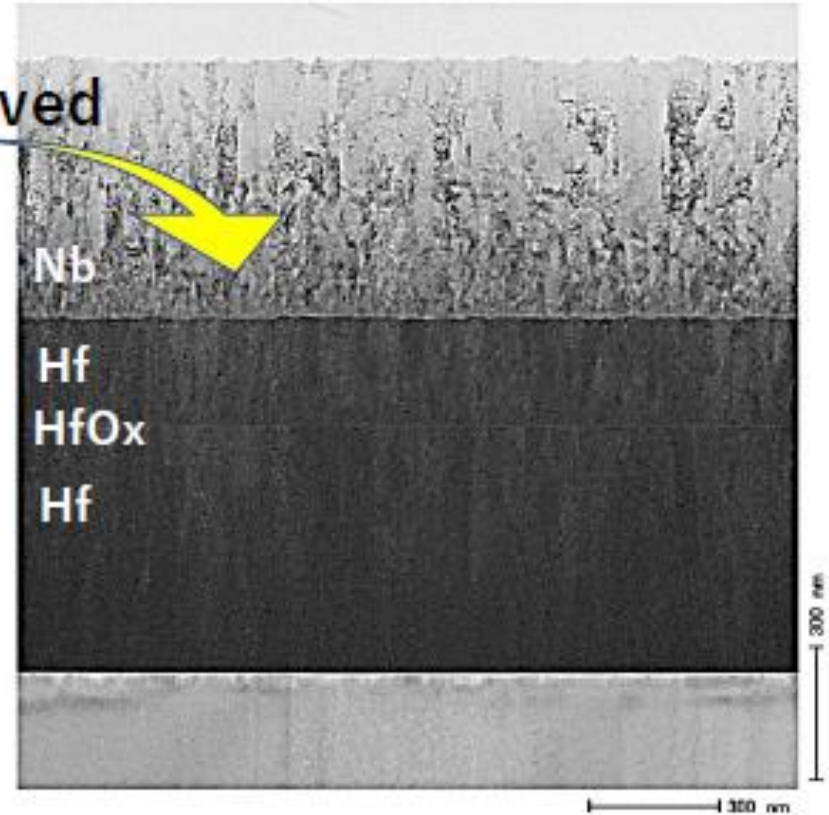
# Improvement of Hf Surface Smoothness

We improved the Hf surface smoothness by optimizing the Hf sputtering parameters.

Old sputtering condition  
Ar 2.0Pa, 80W



New sputtering condition  
Ar 0.5Pa, 50W

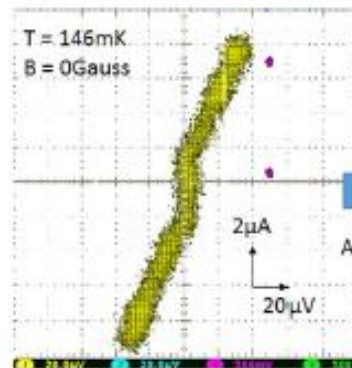
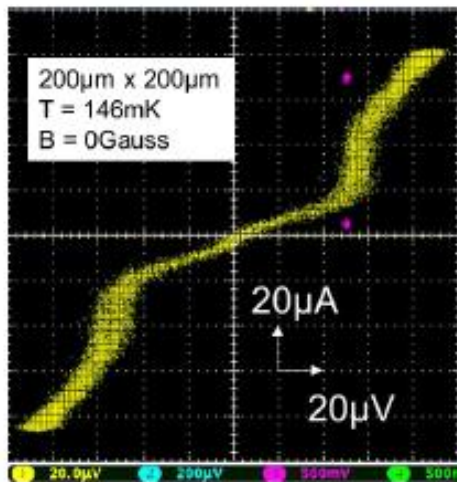


Improved

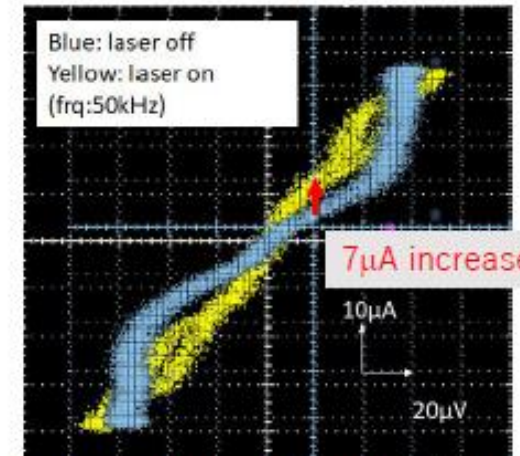
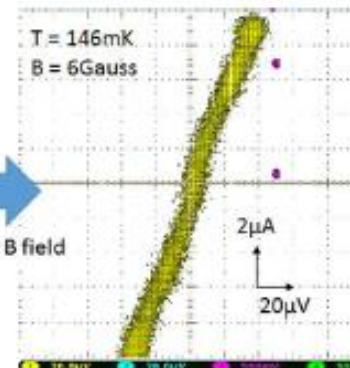
# Hf-STJ with Improved Smoothness

200 $\mu$  square Hf-STJ with improved smoothness. Wire bonding readout line.

- Josephson current is 2 $\mu$ A.
- Effective energy gap  $\Delta = 25\mu$ eV.
- Leakage current at 20 $\mu$ V is 7 $\mu$ A.



Apply B field





# Small-size Hf-STJ with improved smoothness

10 $\mu$ -square Hf-STJ. Signal line was made by sputtered Nb line not by wire bonding.

- Josephson current is 0.7 $\mu$ A.
- Effective energy gap  $\Delta = 130\mu\text{eV}$ .
- Leakage current at 20 $\mu\text{V}$  is 0.3 $\mu\text{A}$  (1/24 of 200 $\mu$ -square Hf-STJ).

