

超伝導体検出器

– SOI技術との融合による遠赤外－光子検出 STJ 開発 –

2015年11月30日 /光量子計測器開発推進室発足会議

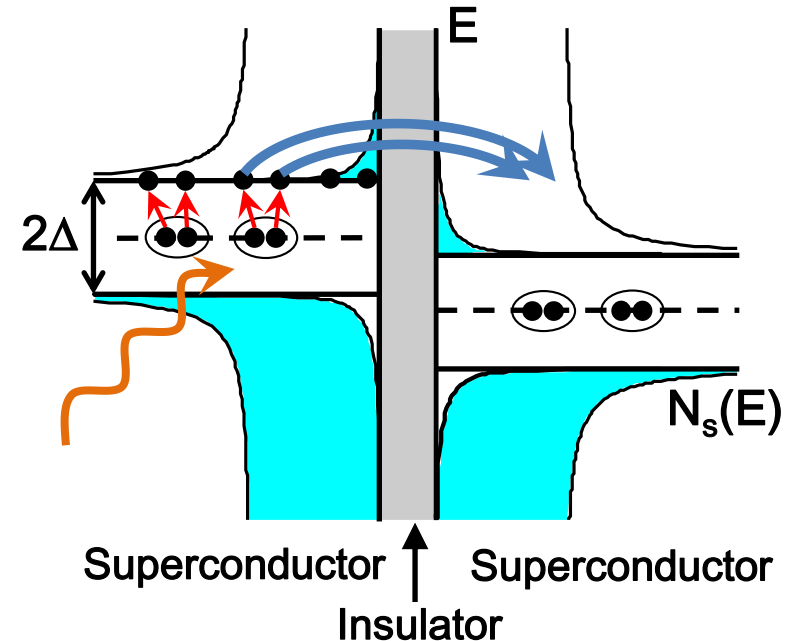
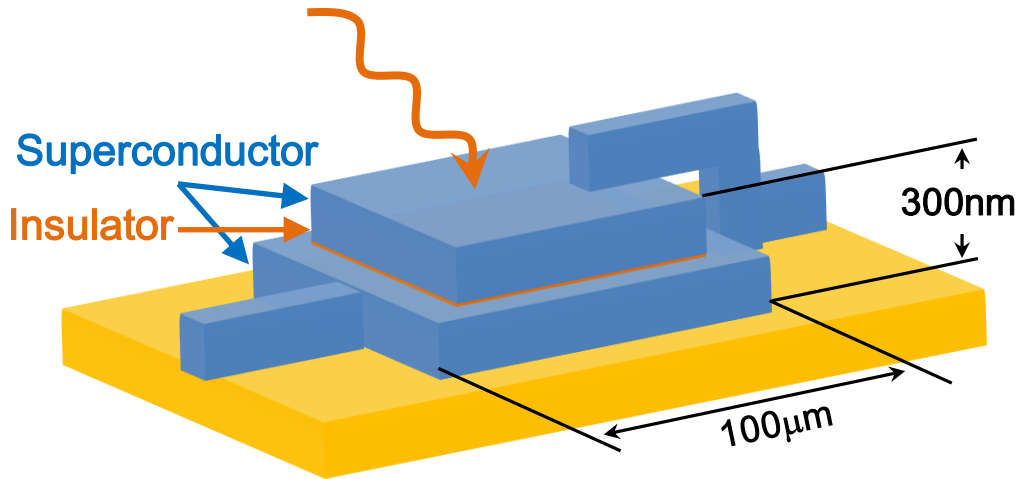
武内勇司 (筑波大 数理物質融合科学センター)

on behalf of Neutrino Decay Collaboration

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Superconducting Tunnel Junction (STJ)

- Superconductor / **Insulator** / Superconductor
Josephson junction device



接合面を挟んで電位差($|V| < 2\Delta$)を印加.

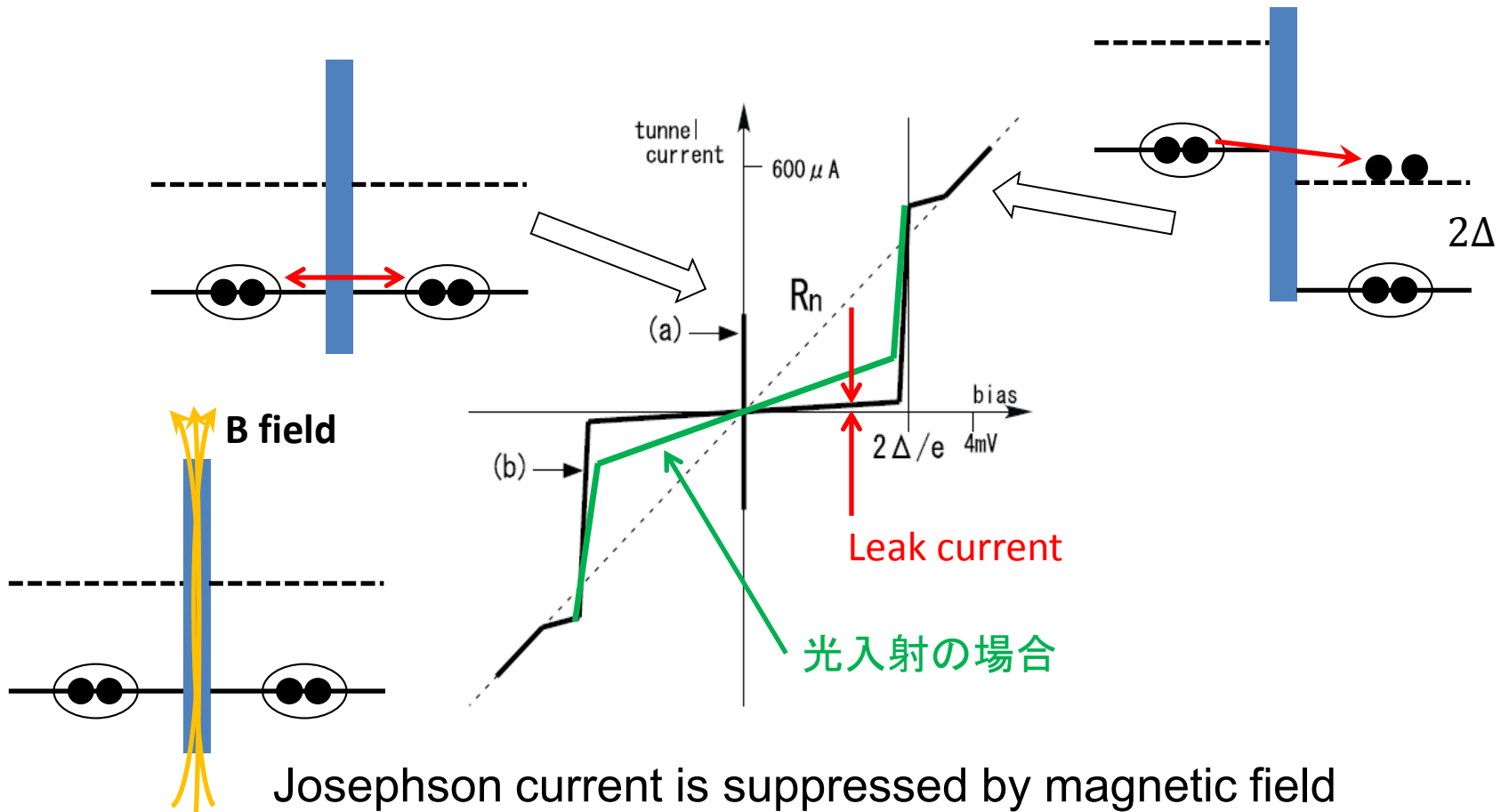
超伝導体に吸収された光子のエネルギーにより複数のクーパー対が解離(励起)し、生成された準粒子によって、エネルギーに比例したトンネル電流が発生.

Δ : Superconducting gap energy

- 超伝導ギャップ(Δ)は遠赤外光子のエネルギーよりもずっと小さい → 原理的には、遠赤外域一光子を検出可能
- $\sim \mu\text{s}$ 程度の比較的高速なパルス応答(Nbの場合) → 光子計数することでS/Nの著しい向上

STJ I-V 特性

- Sketch of a current-voltage (I-V) curve for STJ
- ➔ The Cooper pair tunneling current (DC Josephson current) is seen at $V = 0$, and the quasi-particle tunneling current is seen for $|V| > 2\Delta$



宇宙背景ニュートリノ崩壊探索への応用

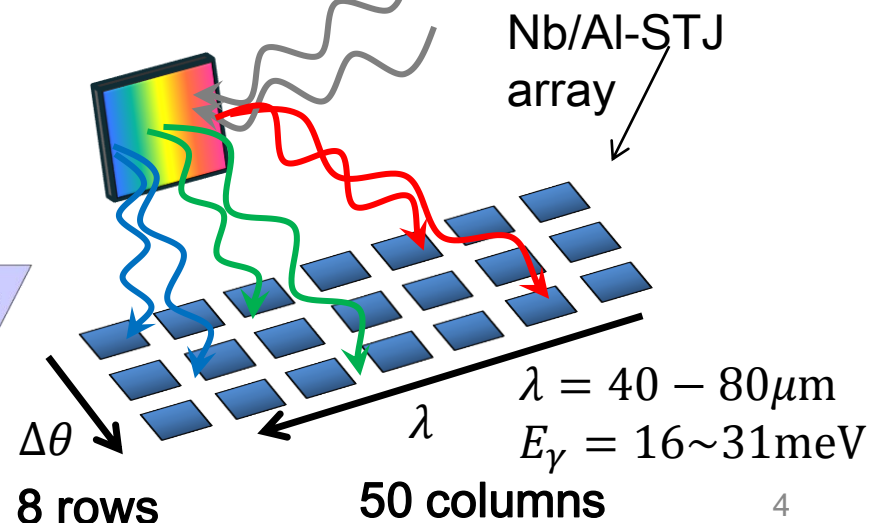
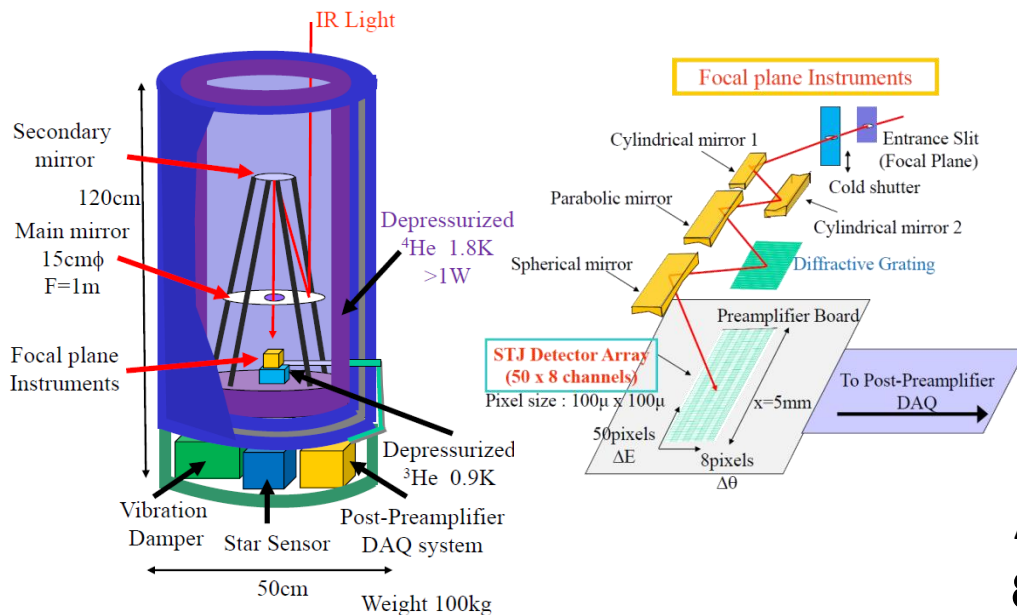
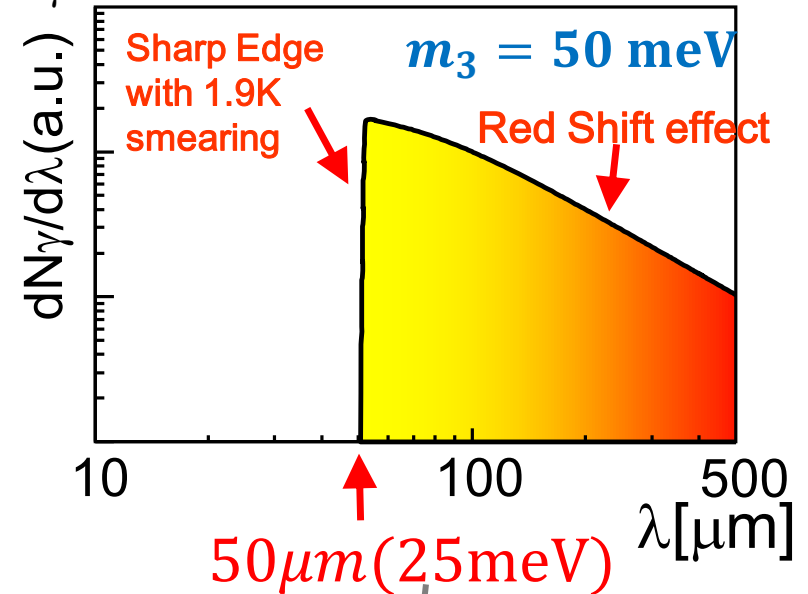
$$\nu_3 \rightarrow \nu_{1,2} + \gamma$$

in the ν_3 rest frame

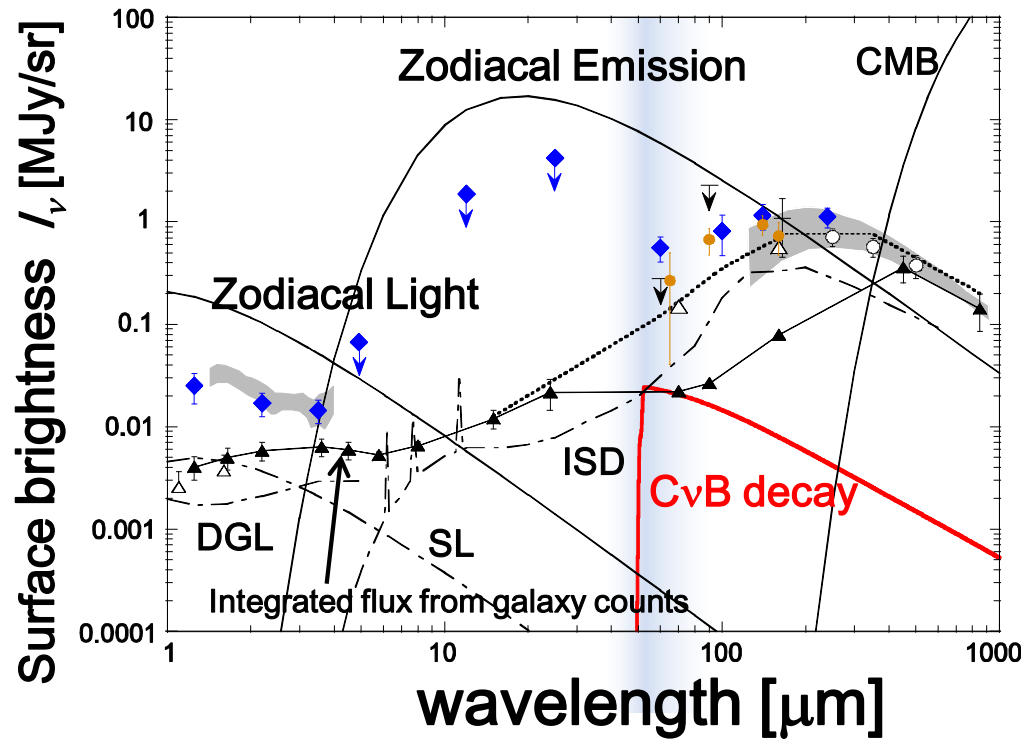
$$E_\gamma = \frac{m_3^2 - m_{1,2}^2}{2m_3}$$

Two body decay

λ_γ distribution in $\nu_3 \rightarrow \nu_2 + \gamma$



宇宙背景ニュートリノ崩壊探索への応用



Telescope parameters

- Main mirror
 - $D=15\text{cm}$, $F=1\text{m}$
- detector
 - 波長 $0.8\mu\text{m}$ あたり
 - $100\mu\text{m} \times 100\mu\text{m} \times 8 \text{ pixels}$
 - 視野角 : $8 \times 10^{-8} \text{ sr}$

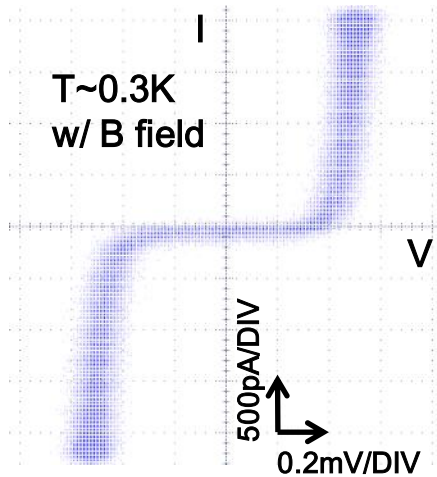
$$\tau = 1 \times 10^{14} \text{ yrs}$$

- Neutrino decay ($m_3 = 50 \text{ meV}$, $\tau_\nu = 1 \times 10^{14} \text{ yrs}$): $I_\nu = 25 \text{ kJy/sr}$
 - $3.3 \times 10^{-20} \text{ W} / 8 \text{ pixels @ } \lambda = 50 \mu\text{m}$
- 200sec の測定でこれを検出
 - ➔ 検出器の性能として最低でも $\text{NEP} < 6.6 \times 10^{-19} \text{ W} / \sqrt{\text{Hz}}$ が必要
 - 実際には, 本物のフォトンによるバックグラウンド(Zodiacal emission)があるので, 更に1 order 位低い必要がある ($\text{NEP} < 1 \times 10^{-19} \text{ W} / \sqrt{\text{Hz}}$)

産総研 CRAVITY 製 Nb/Al-STJ

M. Ukibe et al., Jpn. J. Appl. Phys. 51, 010115 (2012)

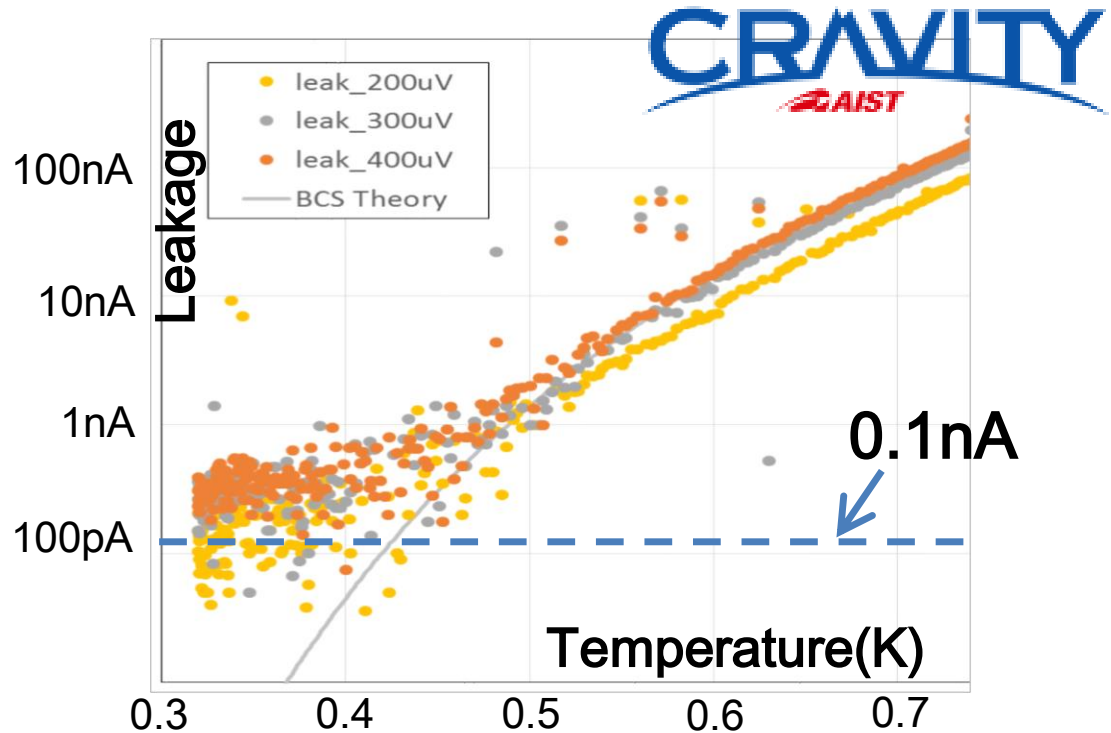
M. Ohkubo et al., IEEE Trans. Appl. Super., 24, 2400208 (2014)



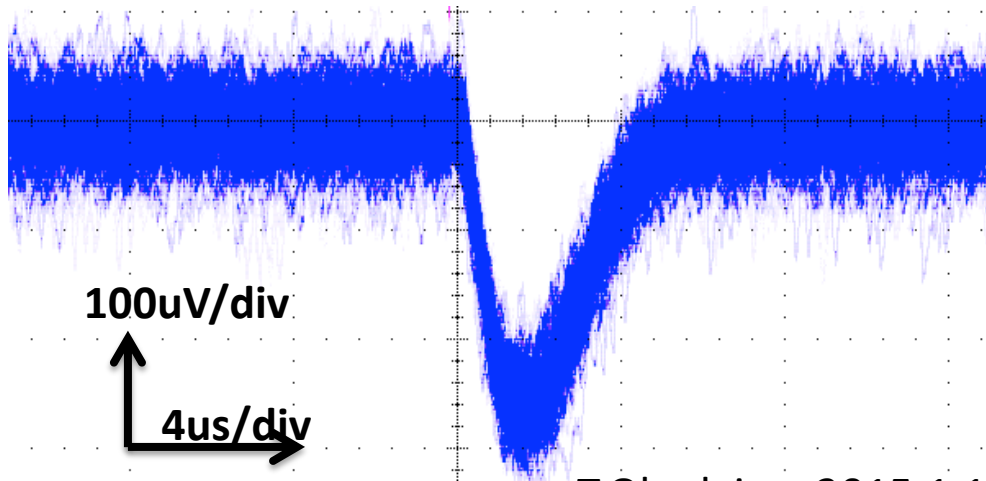
50 μ m \times 50 μ m Nb/Al-STJ
fabricated in CRAVITY at AIST

- $I_{\text{leak}} \sim 0.2 \text{ nA}$ を達成
- 更に小さな junction size のものでテスト中
- リーク電流のショットノイズ由来のNEPは，リーク電流 $i_L = 50 \text{ pA}$ ，超伝導ギャップエネルギー $\Delta = 0.6 \text{ meV}$ ，トラッピングゲイン $G = 10$ とすると

$$\text{NEP} = \frac{1.7\Delta}{G} \sqrt{\frac{2i_L}{e}} \sim 4 \times 10^{-19} \text{ W}/\sqrt{\text{Hz}}$$



STJパルス光応答特性

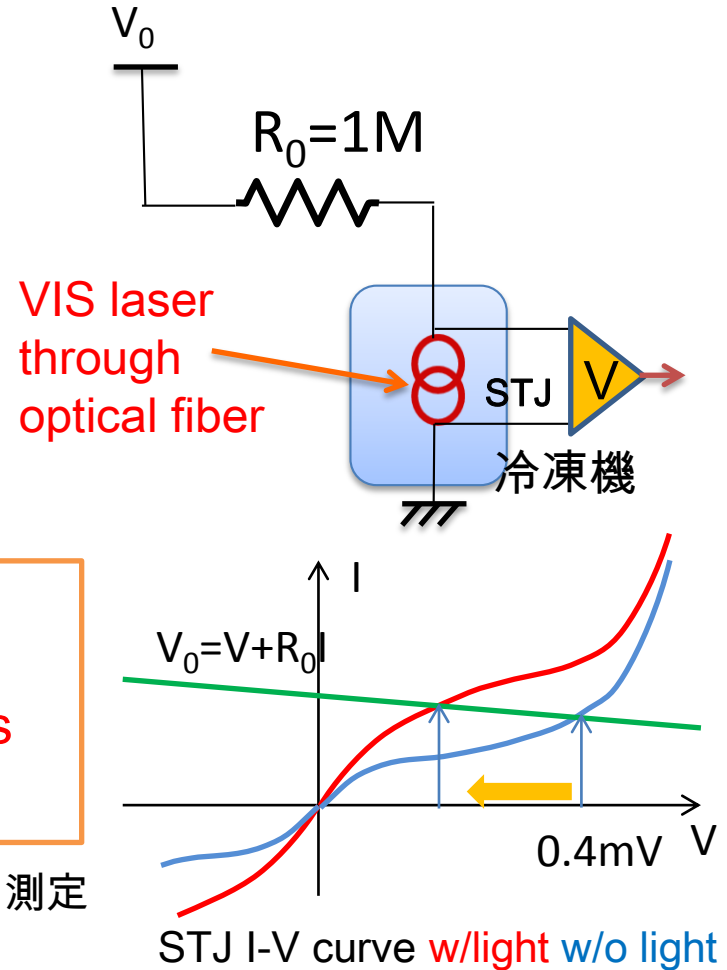


T.Okudaira, 2015.1.16

可視光(465nm)レーザーパルス応答性(AIST製 Nb/Al-STJ 100um角)

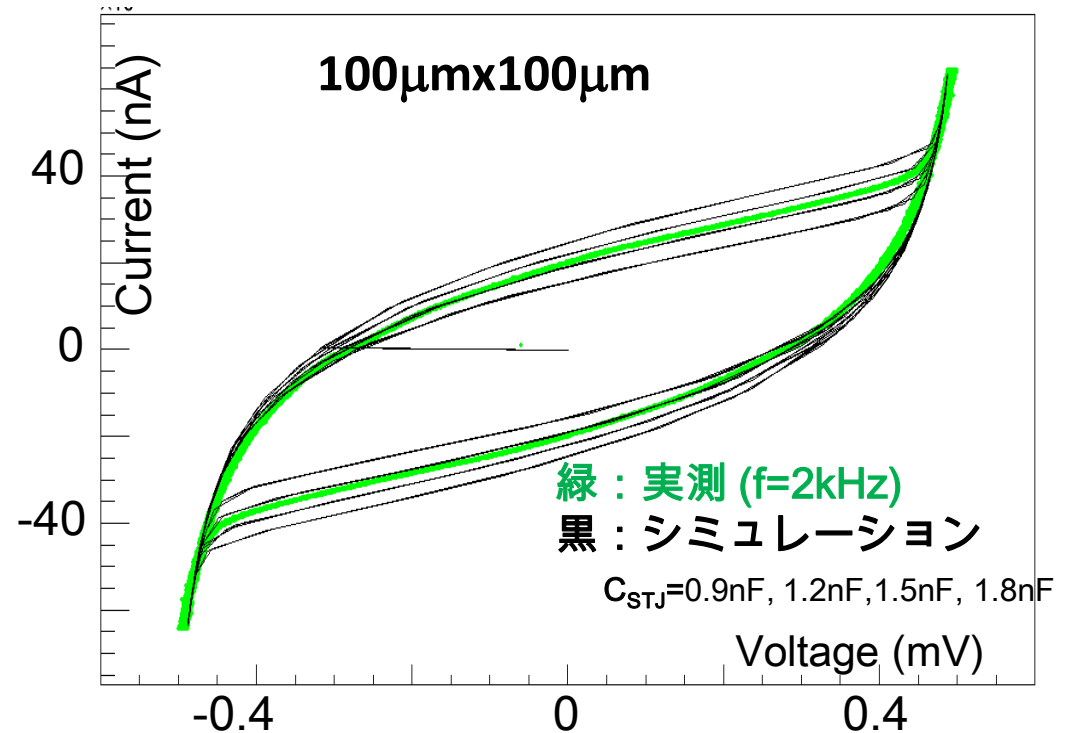
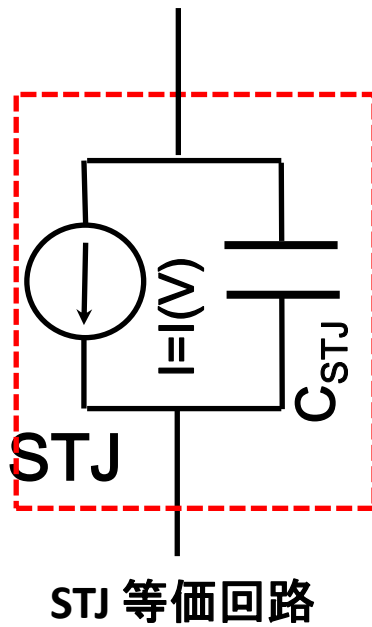
STJ 応答信号時定数: 立下り $\sim 1\mu\text{s}$, 立上り $\sim 2\mu\text{s}$
(もしくは, これより早い)

定電流モードの回路(右上)で測定



- Nb/Al-STJは, $\sim 1\mu\text{s}$ という比較的早い応答速度
➔ 光子計数を行えば, 実効的なNEPは劇的に改善可能
 - 但し読出し系の帯域は $>1\text{MHz}$ を確保する必要あり

STJ キャパシタンス測定

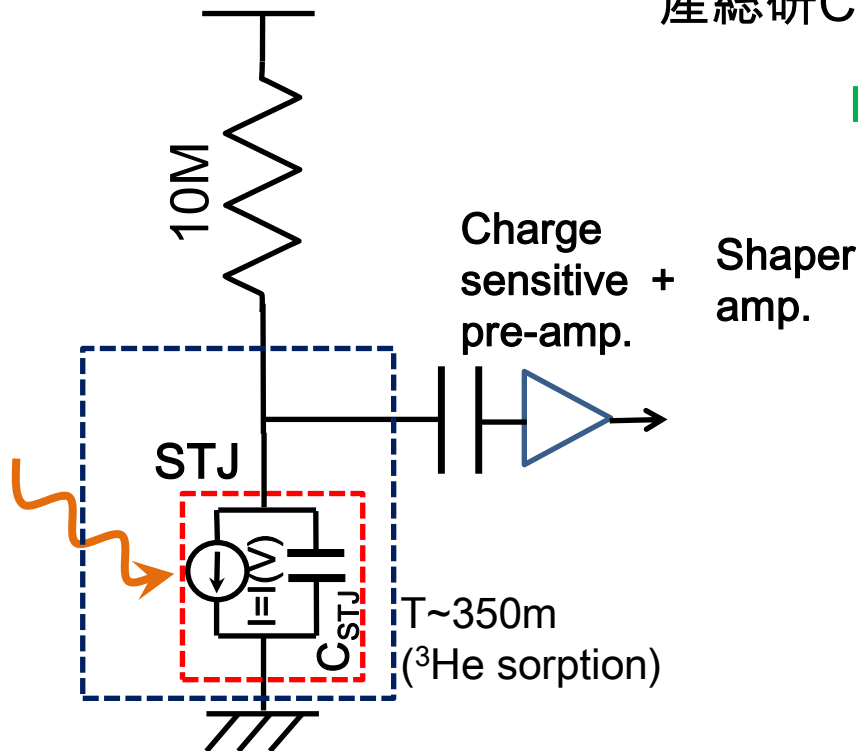


STJ は , junction size に比例したキャパシタンス

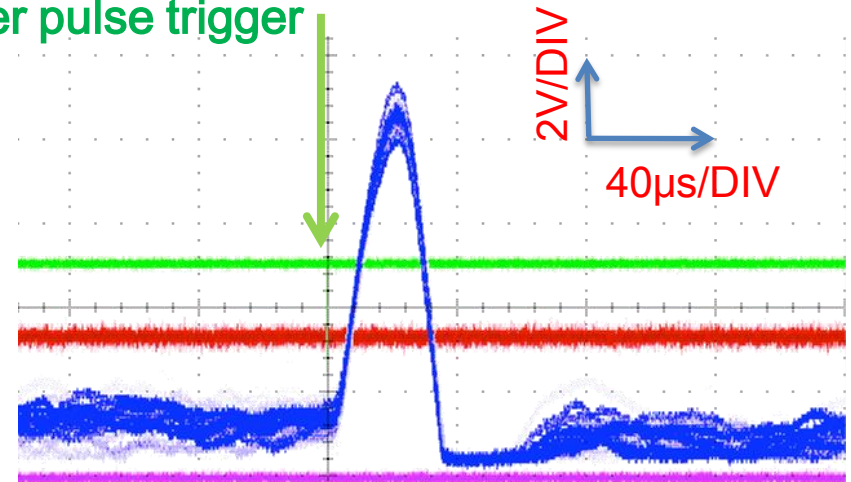
- STJのI-V測定からSTJのキャパシタンスを測定
- SIS接合面の面積に比例する成分 : $\sim 34\text{fF}/\mu\text{m}^2$
- ➔ 20 μm 角のSTJでも 14pF
- ➔ 低入力インピーダンスの電荷積分型アンプでの読出しが必要

100x100 μm^2 Nb/Al-STJ response to 465nm multi-photons

産総研CRAVITY製100x100 μm^2 Nb/Al-STJ



Laser pulse trigger



Pulse height dispersion is consistent with 10-photon detection in STJ

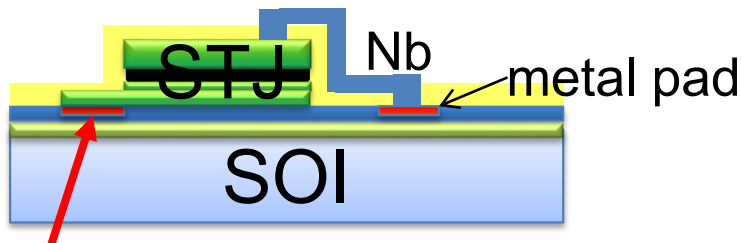
Nb/Al-STJ の低入カインピーダンス電荷積分アンプ読出

- 可視光パルス(波長465nm)に対する応答
- 室温に置かれた電荷積分型のアンプでの読出
 - 観測した出力電荷量は, およそ10光子の検出に対応
- 一光子検出には, 読み出し系のS/N改善が必要

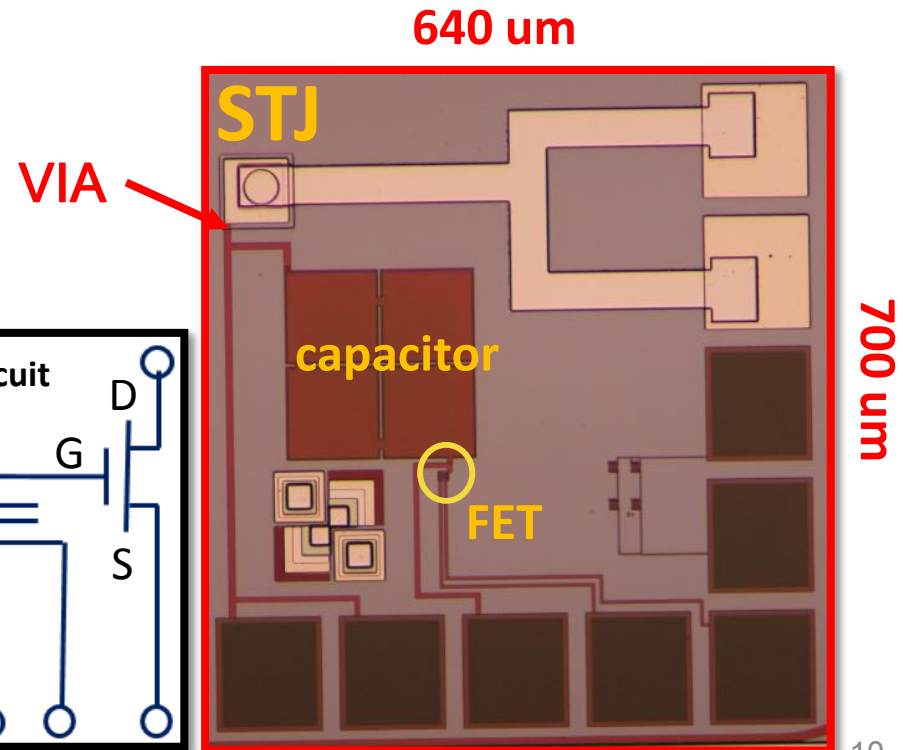
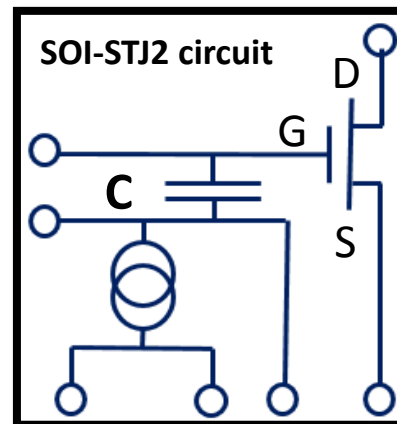
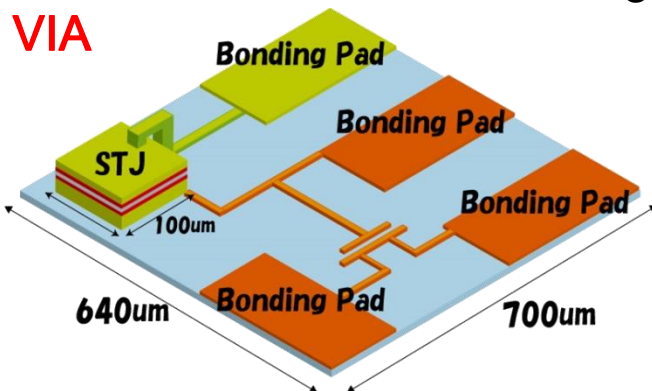
→ 極低温電荷積分型アンプの開発

Development of SOI-STJ

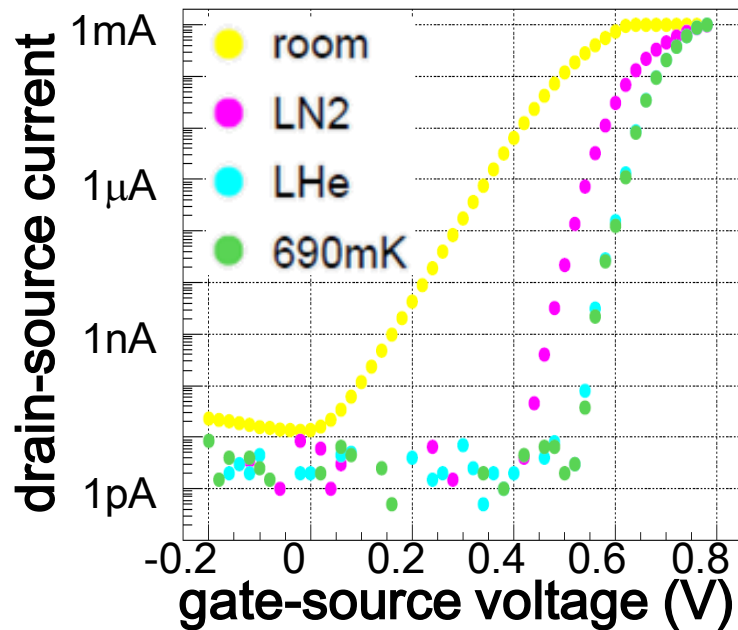
- SOI: Silicon-on-insulator
 - CMOS in FD-SOI is reported to work at 4K by T. Wada (JAXA), et al.
J Low Temp Phys 167, 602 (2012)
- SOI と STJ の融合 (SOI回路一体型 STJ の基礎研究)
 - STJ layers are fabricated **directly on** a SOI pre-amplifier board and cooled down together with the STJ
- Started test with Nb/Al-STJ on SOI with p-MOS and n-MOS FET



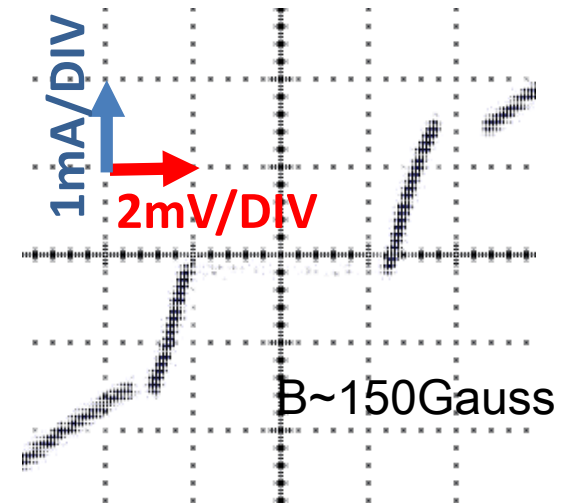
STJ lower layer has electrical contact with SOI circuit through **VIA**



SOI上にSTJを形成後の特性



nMOS-FET in FD-SOI wafer on which a STJ is fabricated at KEK

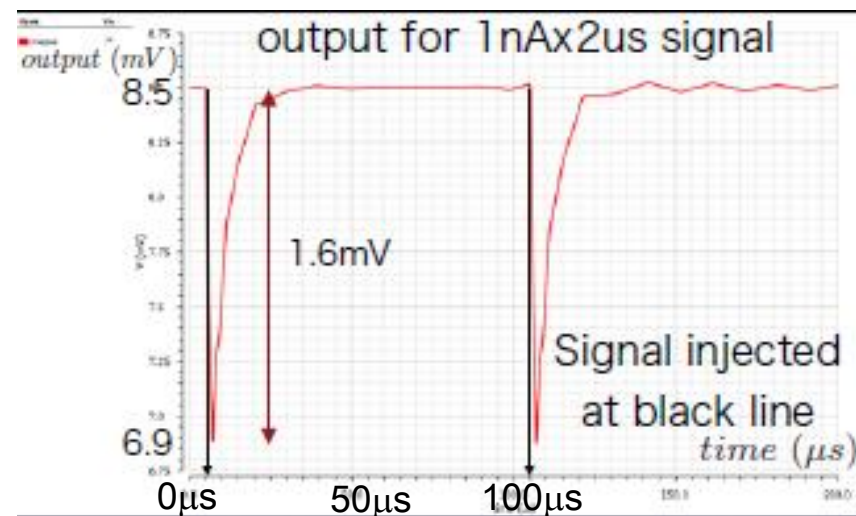
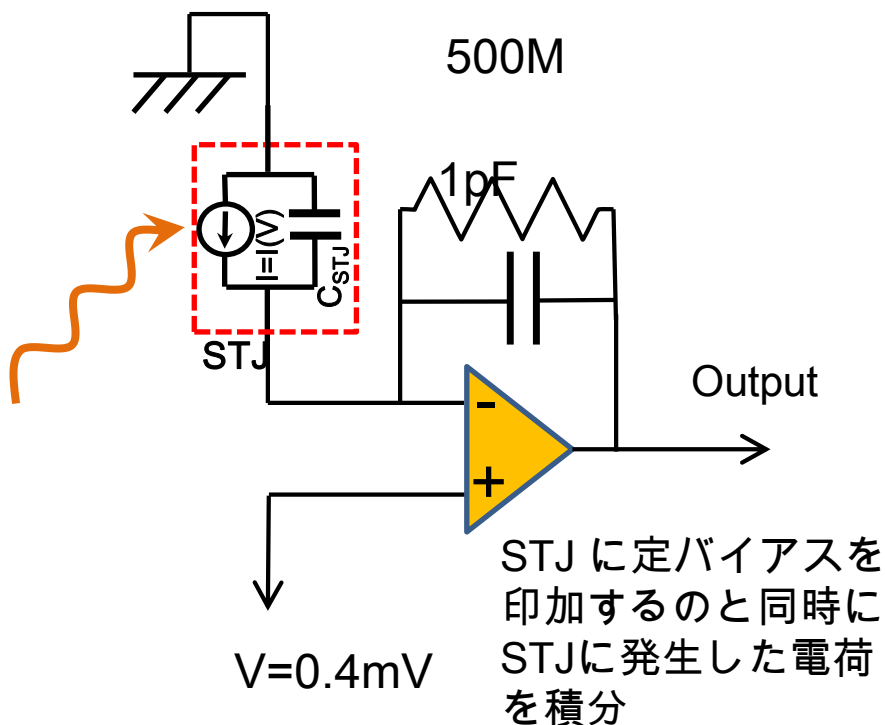


I-V curve of a STJ fabricated at KEK on a FD-SOI wafer

- Both nMOS and pMOS-FET in FD-SOI wafer on which a STJ is fabricated work fine at temperature down below 1K
 - 極低温では，スレッシュホールド電圧のシフト，サブスレッシュホールド領域のドレイン電流抑制，飽和領域でのドレイン電流の上昇など，特性が変動
- Nb/Al-STJ fabricated at KEK on FD-SOI works fine
- We are also developing SOI-STJ where STJ is fabricated at CRAVITY

SOI Pre-amplifier development

- 遠赤外一光子検出に向け，前段階として近赤外一光子検出に最適化した電荷積分型アンプをVDEC* が提供するSPICE simulation で設計中
- 極低温での SOI MOSFET の振る舞いをシミュレーションに組み込むため KEK や JAXA と共同研究で SPICE用MOSFETパラメータを構築中
 - 様々なL(チャンネル長：L=0.4~5um)やW(チャンネル幅：W=1~10um)を持ったFETの3Kにおける特性の測定



室温でのFET パラメータを仮定したSPICE simulation
入力電荷:2fC
1.3eV(波長1 μ m)の一光子入射相当

* VLSI Design and Education Center(VDEC), the U. Tokyo in collaboration with Synopsys, Inc., Cadence Design Systems, Inc., and Mentor Graphics, Inc.

まとめ

- 遠赤外(50 μ m)の一光子検出が可能な検出器を STJ + SOI の技術を用いて開発中
 - 光子計数により，実効的にNEP で $\sim 10^{-20} W / \sqrt{Hz}$ を目指す
- Nb/Al-STJ は，産総研CRAVITY で世界最高水準の低リーク電流のものが得られている(<200pA @ 50 μ m角)
 - 20 μ m角，10 μ m角のものもテスト中
- SOIに技術を用いた極低温アンプによる読出し回路を開発中
 - 様々なW/LをもつSOI MOSFETの極低温でのI-V測定→SPICE シミュレーションにもちいるFETパラメータ抽出
 - 光子計数の利点を最大限に生かす高速アンプ(帯域>1MHz)
 - SOI アンプー体型STJの可能性

Backup

Neutrino

- Neutrino has 3 mass generations (ν_1, ν_2, ν_3)
- Neutrino flavor states (ν_e, ν_μ, ν_τ) are not mass eigenstates

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

➔ Neutrino flavor oscillates during the flight, and squared mass differences ($\Delta m_{12}^2, |\Delta m_{23}^2|$) have been measured, **but their absolute masses are not measured yet!**

❑ **Heavier neutrinos (ν_2, ν_3) are not stable**

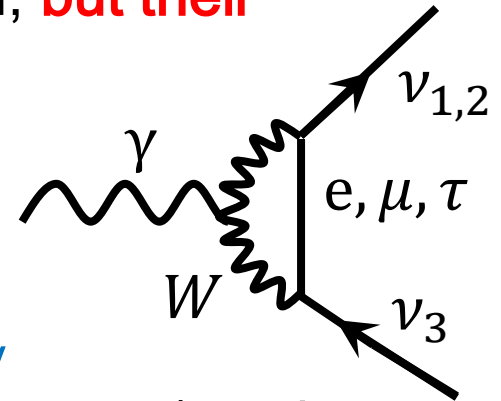
– Neutrino can decay through the loop diagrams

– $\nu_3 \rightarrow \nu_{1,2} + \gamma$

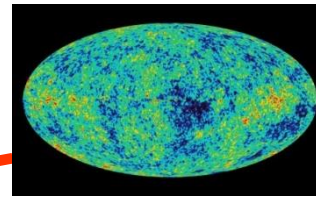
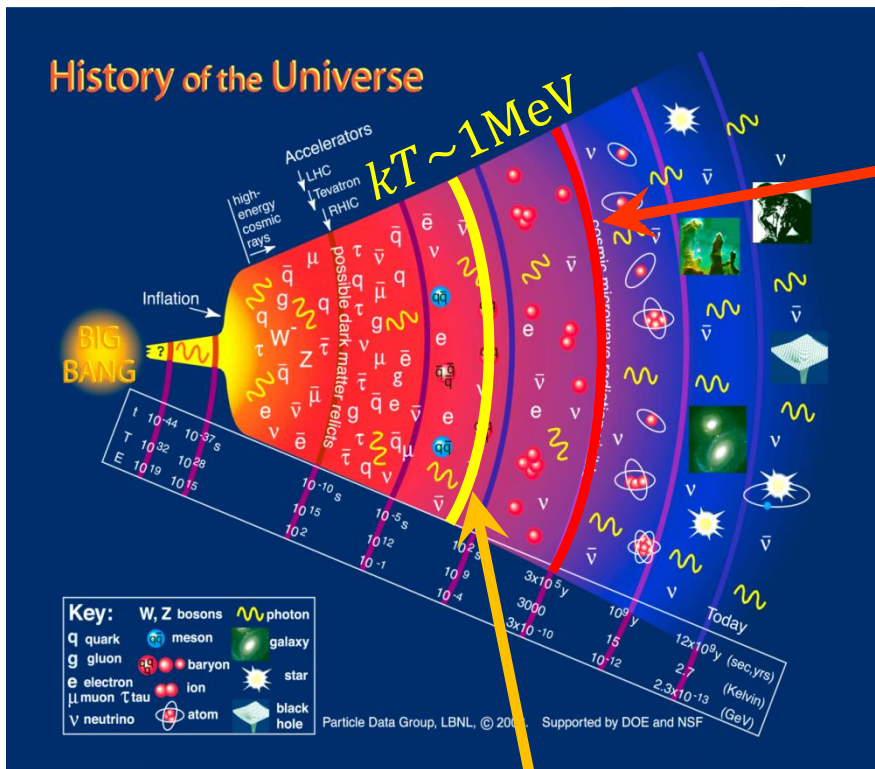
– Neutrino mass can be determined from the decay

✓ However, neutrino lifetime is expected to be very long (much longer than the age of universe)

➔ **We adopt Cosmic neutrino background (CvB) as the neutrino source for neutrino decay search**



Cosmic neutrino background (CνB)



CMB

$$n_\gamma = 411/\text{cm}^3$$

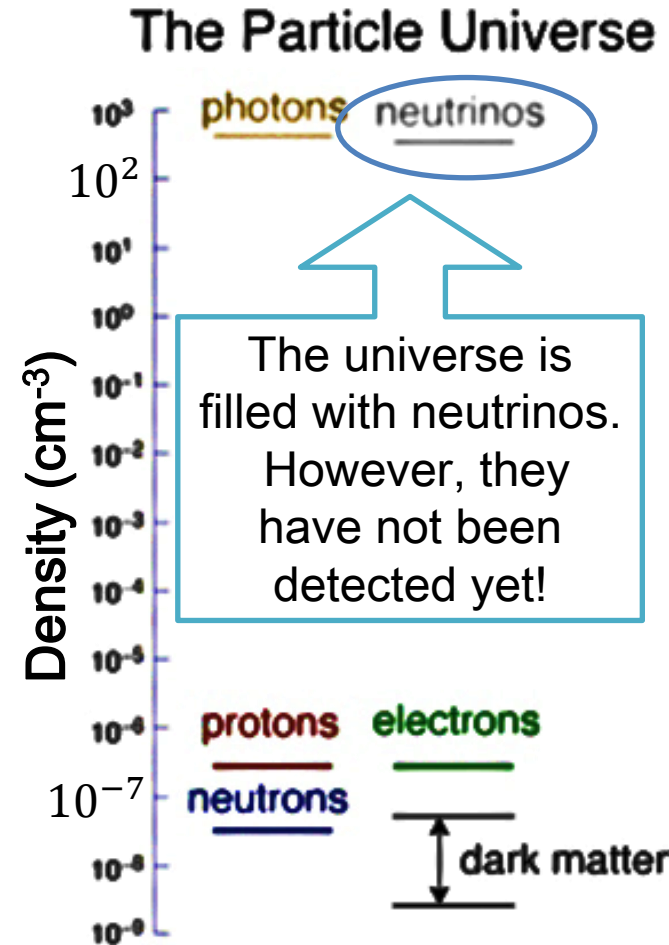
$$T_\gamma = 2.73 \text{ K}$$

CνB (=neutrino decoupling)
~1s after the big bang

$$T_\nu = \left(\frac{4}{11}\right)^{\frac{1}{3}} T_\gamma = 1.95 \text{ K}$$

$$\langle p_\nu \rangle = 0.5 \text{ meV}/c$$

$$n_\nu + n_{\bar{\nu}} = \frac{3}{4} \left(\frac{T_\nu}{T_\gamma}\right)^3 n_\gamma = 110/\text{cm}^3$$

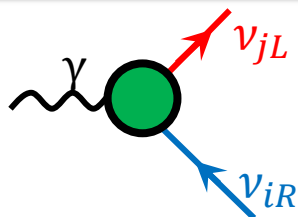


Motivation of ν -decay search in C ν B

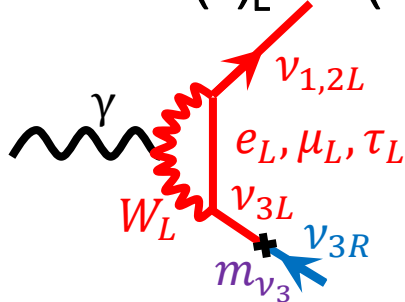
- Search for $\nu_3 \rightarrow \nu_{1,2} + \gamma$ in cosmic neutrino background (C ν B)
 - Search for anomalous magnetic moment of neutrino
 - Direct detection of C ν B
 - Determination of neutrino mass: $m_3 = (m_3^2 - m_{1,2}^2)/2E_\gamma$
- Aiming at a sensitivity to ν lifetime for $\tau(\nu_3) = O(10^{17} \text{ yrs})$
 - Standard Model expectation: $\tau = O(10^{43} \text{ yrs})$
 - Experimental lower limit: $\tau > O(10^{12} \text{ yrs})$
 - L-R symmetric model (for Dirac neutrino) predicts down to $\tau = O(10^{17} \text{ yrs})$ for W_L - W_R mixing angle $\zeta < 0.02$

Magnetic moment term
(need L-R coupling)

$$\bar{\nu}_{jL} i \sigma_{\mu\nu} q^\nu \nu_{iR}$$



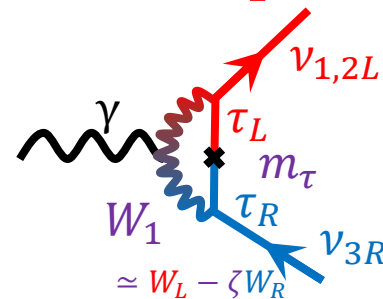
SM: $SU(2)_L \times U(1)_Y$



$$\Gamma \sim (10^{43} \text{ yr})^{-1}$$

Suppressed by m_ν and GIM

LRS: $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$



$$\Gamma \sim (10^{17} \text{ yr})^{-1}$$

Only suppressed by L-R mixing (ζ)

PRL 38,(1977)1252, PRD 17(1978)1395

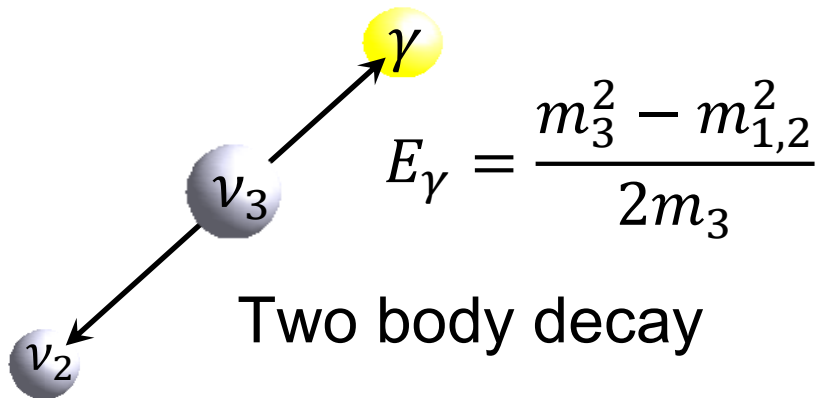
$$\begin{pmatrix} W_1 \\ W_2 \end{pmatrix} = \begin{pmatrix} \cos\zeta & -\sin\zeta \\ \sin\zeta & \cos\zeta \end{pmatrix} \begin{pmatrix} W_L \\ W_R \end{pmatrix}$$

**10^{26}
enhancement to
SM**

Photon Energy (Wavelength) in Neutrino Decay

$$\nu_3 \rightarrow \nu_{1,2} + \gamma$$

in the ν_3 rest frame



Two body decay

$m_3 = 50 \text{ meV}$

$E_\gamma = 24.8 \text{ meV}$
($\lambda = 50 \mu\text{m}$)

$E_\gamma = 24 \text{ meV}$
($\lambda = 51 \mu\text{m}$)

$m_2 = 8.7 \text{ meV}$

$m_1 = 1 \text{ meV}$

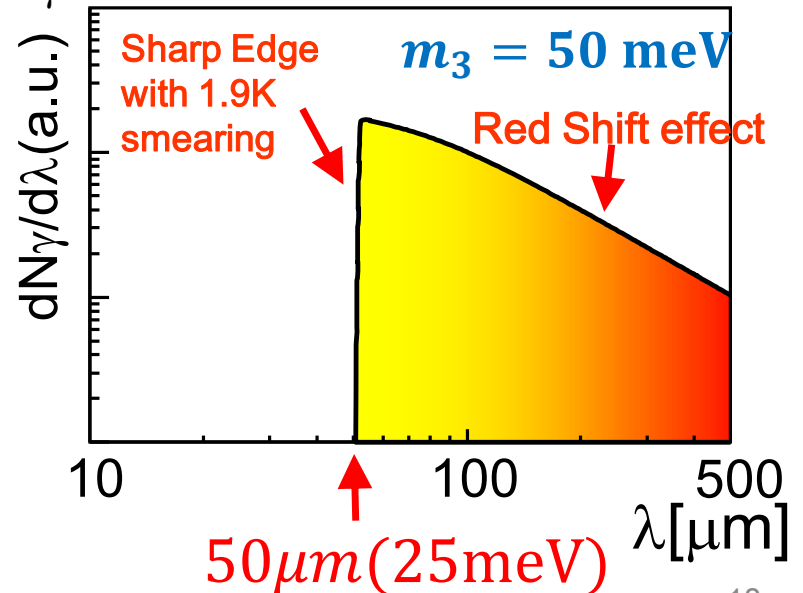
$E_\gamma = 4.4 \text{ meV}$
($282 \mu\text{m}$)

- From neutrino oscillation
 - $|\Delta m_{23}^2| = |m_3^2 - m_2^2| \sim 2.4 \times 10^{-3} \text{ eV}^2$
 - $\Delta m_{12}^2 \sim 7.65 \times 10^{-5} \text{ eV}^2$
- From Planck+WP+highL+BAO
 - $\sum m_i < 0.23 \text{ eV}$

→ $50 \text{ meV} < m_3 < 87 \text{ meV}$

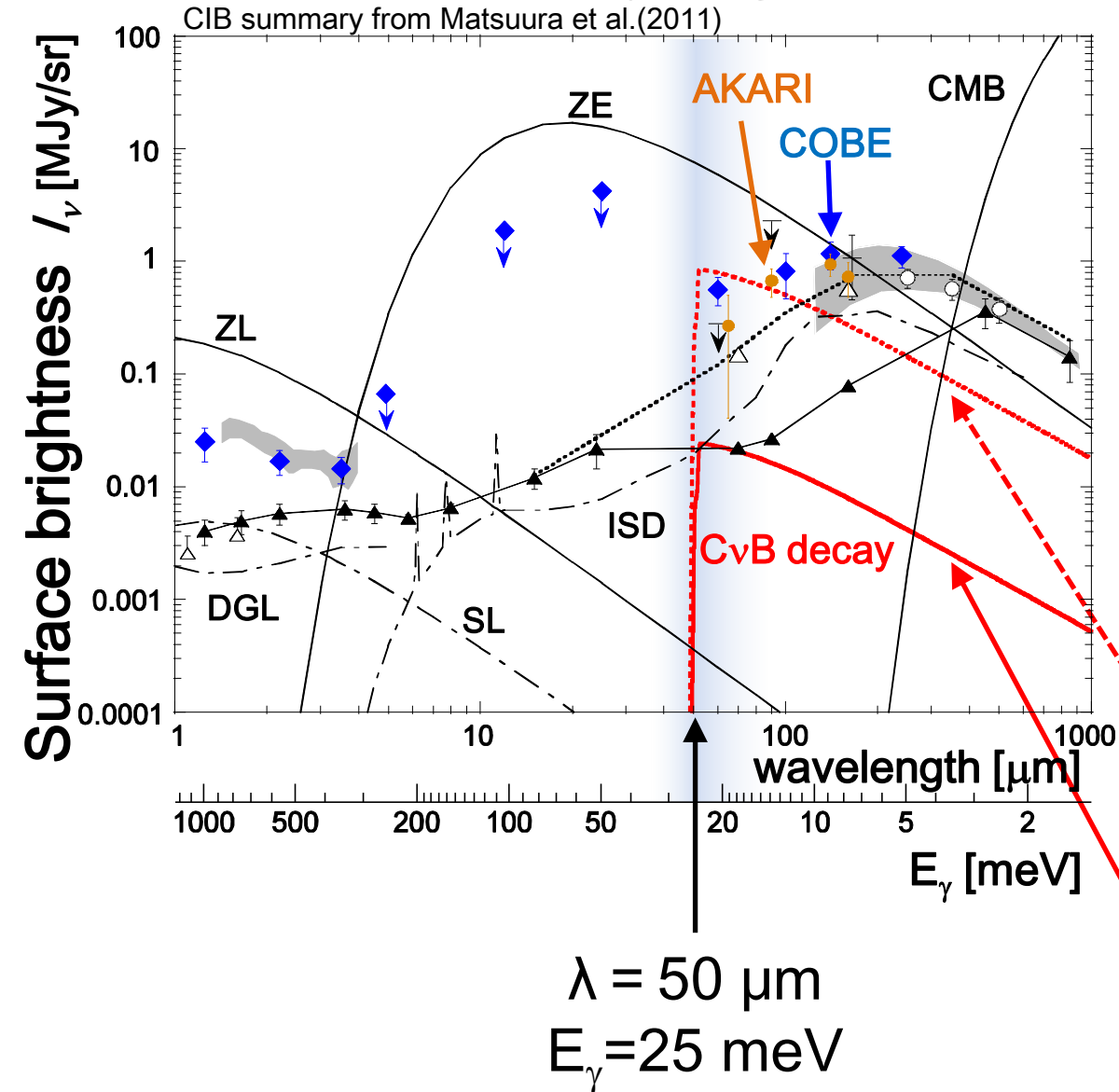
$E_\gamma^{\text{rest}} = 14 \sim 24 \text{ meV}$ ($\lambda_\gamma = 51 \sim 89 \mu\text{m}$)

λ_γ distribution in $\nu_3 \rightarrow \nu_2 + \gamma$



CνB decay signal and Backgrounds

CIB summary from Matsuura et al.(2011)



at $\lambda = 50 \mu\text{m}$

Zodiacal Emission(ZE)

$$I_\nu \sim 8 \text{ MJy/sr}$$

CIB

$$\lambda I_\lambda \sim 0.1\text{-}0.5 \text{ MJy/sr}$$

CνB decay

Expected E_γ spectrum

$$m_3 = 50 \text{ meV}$$

$$\tau = 3 \times 10^{12} \text{ yrs}$$

$$I_\nu \sim 0.8 \text{ MJy/sr}$$

Excluded by S.H.Kim et. al 2012

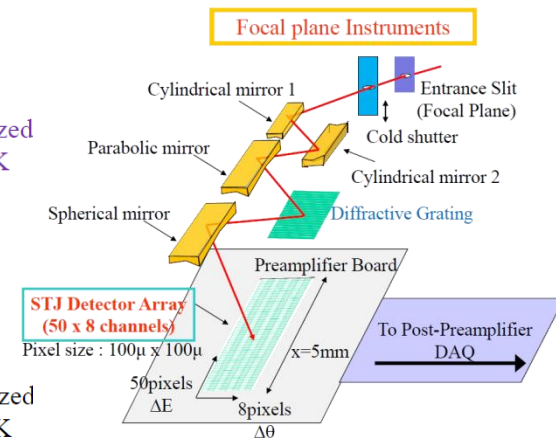
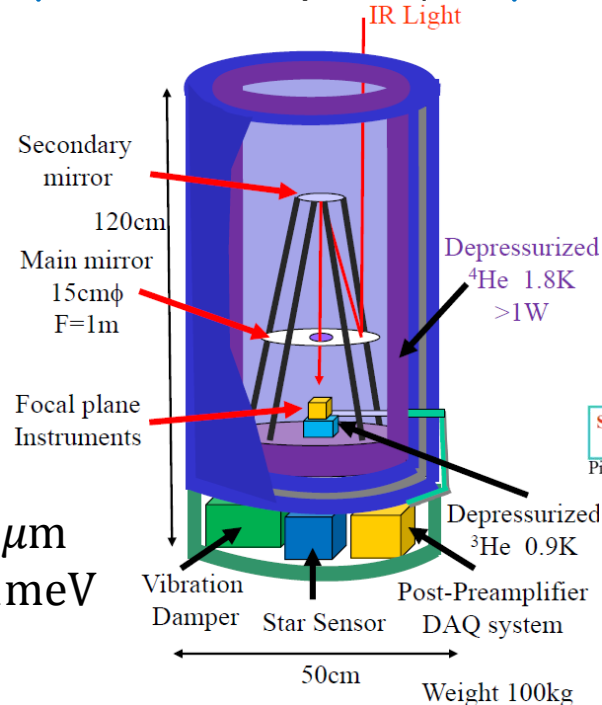
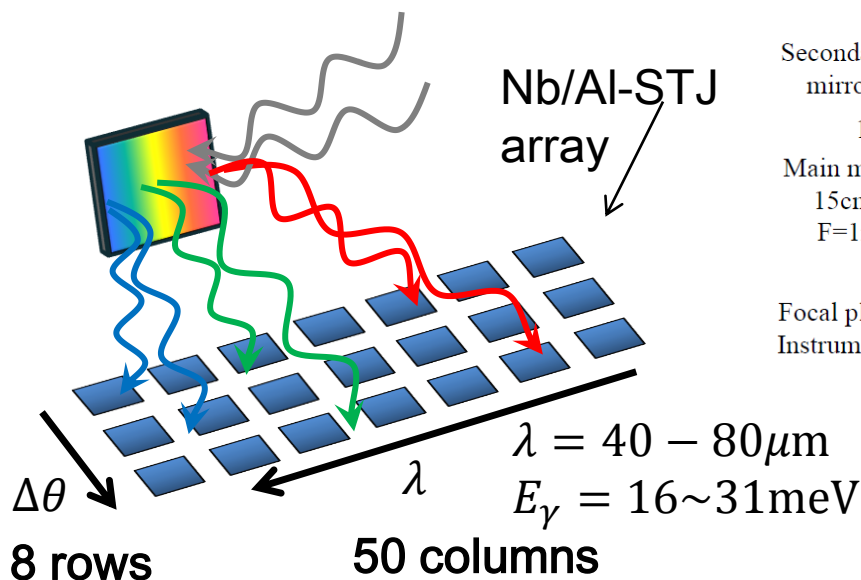
$$\tau = 1 \times 10^{14} \text{ yrs}$$

$$I_\nu \sim 25 \text{ kJy/sr}$$

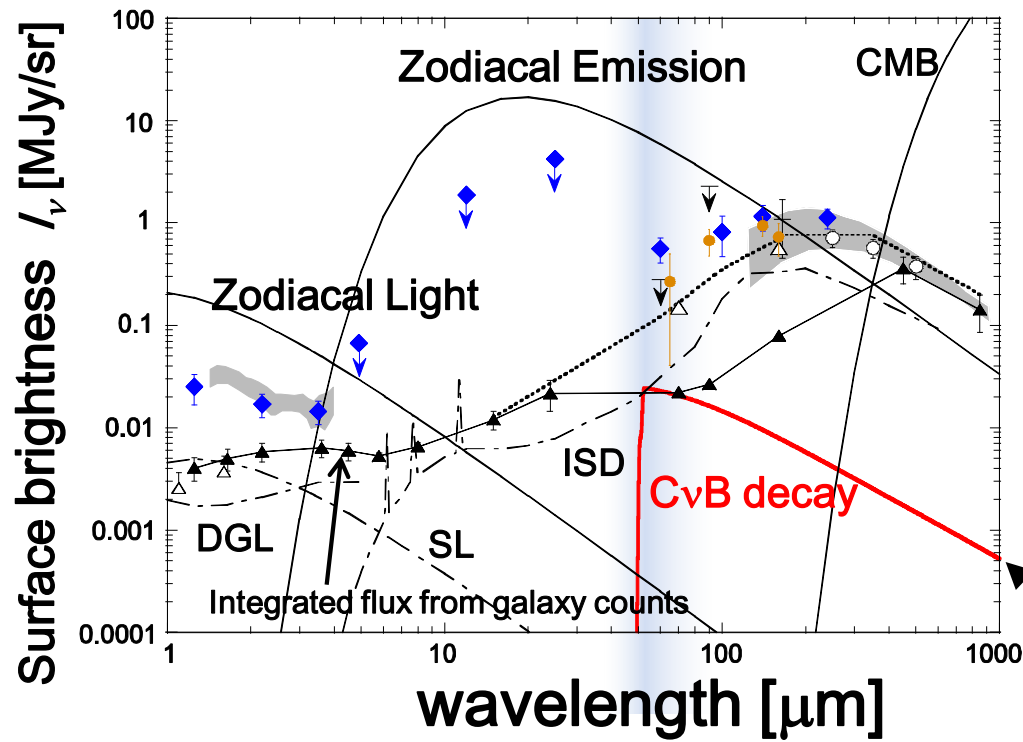
Proposed rocket experiment

with a diffraction grating and Nb/Al-STJ array combination

- 200-sec measurement at altitude of 200~300km
 - Telescope with **diameter of 15cm** and **focal length of 1m**
 - All optics (mirrors, filters, shutters and grating) will be cooled at $\sim 1.8\text{K}$
- At the focal point, diffraction grating covering **$\lambda=40\text{--}80\mu\text{m}$ (16-31meV)** and array of Nb/Al-STJ pixels of **50(in wavelength distribution) x 8(in spatial distribution)** are placed
 - Each Nb/Al-STJ pixel is used as **a single-photon counting detector** for FIR photon in $\lambda=40\text{--}80\mu\text{m}$ ($\Delta\lambda = 0.8\mu\text{m}$)
 - Sensitive area of **$100\mu\text{m} \times 100\mu\text{m}$** for each pixel (**$100\mu\text{rad} \times 100\mu\text{rad}$** in viewing angle)



Expected precision in the spectrum measurement



Telescope parameters

- Main mirror
 - $D=15\text{cm}$, $F=1\text{m}$
- detector
 - sensitive area
 $100\mu\text{m} \times 100\mu\text{m}$ / pixel
 - 50×8 array

$$\Delta\lambda = \frac{80\mu\text{m} - 40\mu\text{m}}{50} = 0.8\mu\text{m}$$

$$\tau = 1 \times 10^{14} \text{ yrs}$$

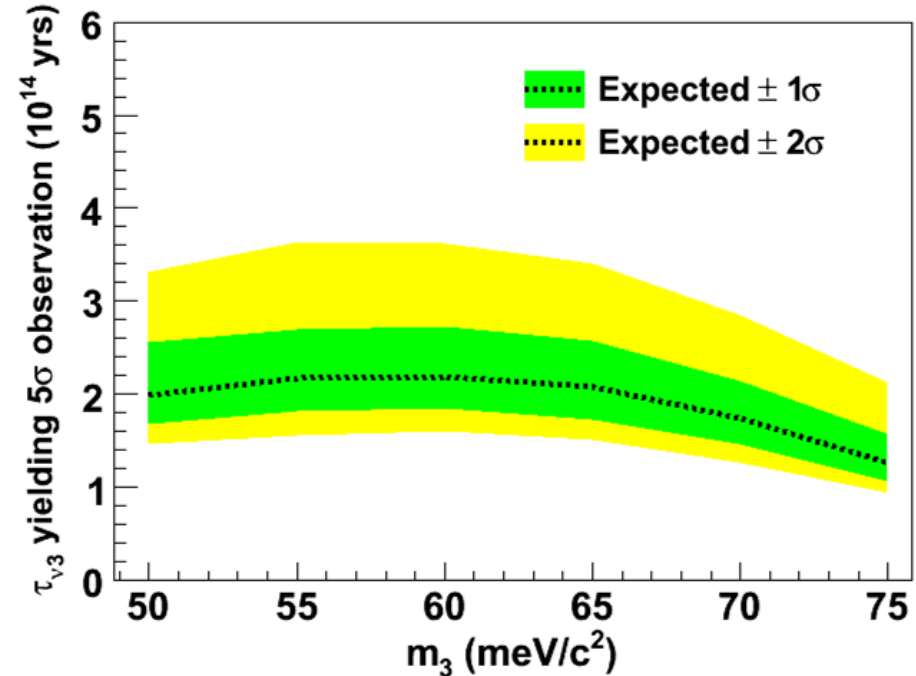
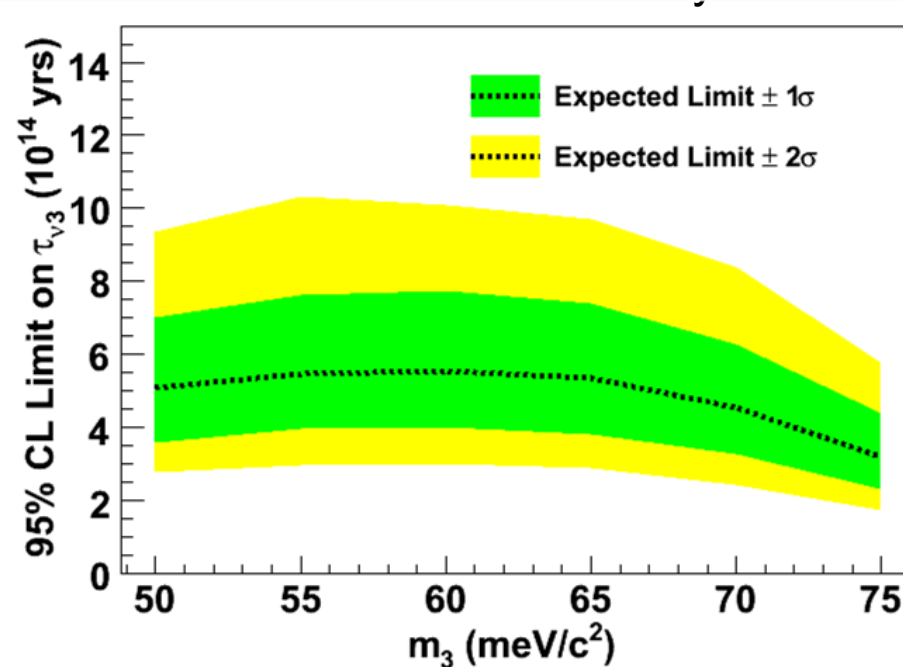
- Zodiacal emission \Rightarrow **343Hz / pixel**
 - 200sec measurement: 0.55M events / 8 pixels (at $\lambda = 50\mu\text{m}$)
 - **0.13%** accuracy measurement for each wavelength: **$\delta(I_\nu)=11\text{kJy/sr}$**
- Neutrino decay ($m_3 = 50 \text{ meV}$, $\tau_\nu = 1 \times 10^{14} \text{ yrs}$): **$I_\nu=25\text{kJy/sr}$**
 - **2.3σ** away from statistical fluctuation in ZE measurement

ν decay with $\tau_\nu = 10^{14} \text{ yrs}$ is possible to detect, or set lower limit!

Sensitivity to neutrino decay

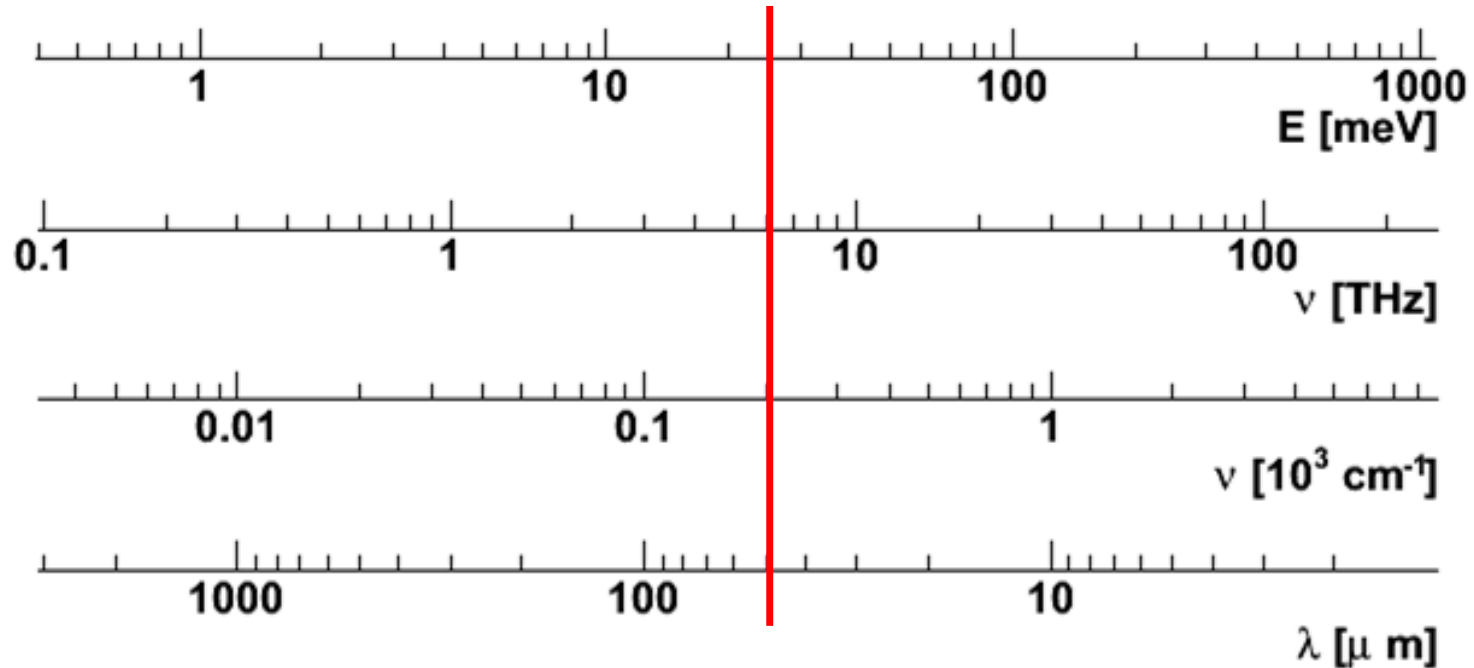
Parameters in the rocket experiment simulation

- telescope dia.: 15cm
- 50-column (λ : 40 μm – 80 μm) \times 8-row array
- Viewing angle per single pixel: 100 μrad \times 100 μrad
- Measurement time: 200 sec.
- Photon detection efficiency: 100%



- Can set lower limit on ν_3 lifetime at $4\text{--}6 \times 10^{14}$ yrs if no neutrino decay observed
- If ν_3 lifetime were 2×10^{14} yrs, the signal significance is at 5 σ level

Energy/Wavelength/Frequency



$$E_\gamma = 25 \text{ meV}$$

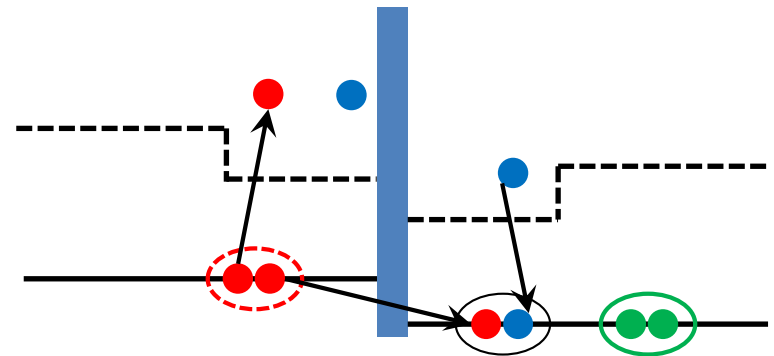
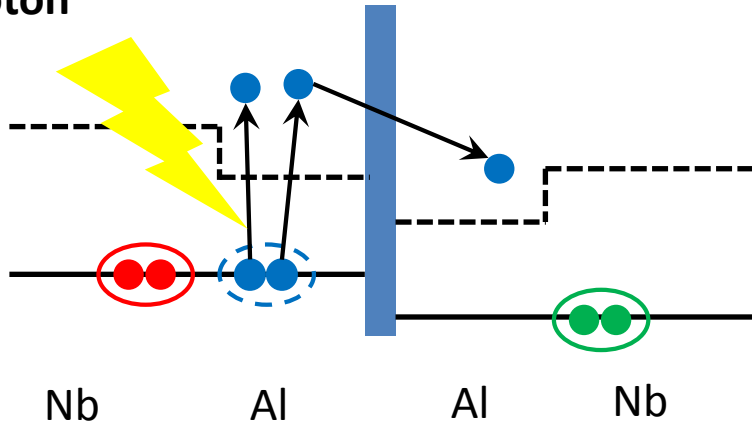
$$\nu = 6 \text{ THz}$$

$$\lambda = 50 \mu\text{m}$$

STJ back-tunneling effect

- Quasi-particles near the barrier can mediate Cooper pairs, resulting in true signal gain
 - Bi-layer fabricated with superconductors of different gaps $\Delta_{\text{Nb}} > \Delta_{\text{Al}}$ to enhance quasi-particle density near the barrier
 - Nb/Al-STJ Nb(200nm)/Al(10nm)/AlOx/Al(10nm)/Nb(100nm)
- Gain: $2 \sim 200$

Photon



STJ energy resolution

Statistical fluctuation in number of quasi-particles → energy resolution

→ Smaller superconducting gap energy Δ yields better energy resolution

$$\sigma_E = \sqrt{(1.7\Delta)FE}$$

Δ : Superconducting gap energy
F: fano factor
E: Photon energy

	Si	Nb	Al	Hf
T _c [K]		9.23	1.20	0.165
Δ [meV]	1100	1.550	0.172	0.020

T_c: SC critical temperature
Need ~1/10 T_c for practical operation

Nb

Well-established as Nb/Al-STJ
(back-tunneling gain from Al-layers)

$$N_{q.p.} = 25 \text{ meV} / 1.7 \Delta = 9.5$$

Poor energy resolution, but a single-photon detection is possible

Hf

Hf-STJ is not established as a practical photon detector yet

$$N_{q.p.} = 25 \text{ meV} / 1.7 \Delta = 735$$

2% energy resolution is achievable if Fano factor < 0.3 for a single-photon

検出器に要求されるNEP

Telescope parameters

- Main mirror: $D=15\text{cm}$, $F=1\text{m}$
- detector
 - 波長 $0.8\mu\text{m}$ ($= (80\mu\text{m}-40\mu\text{m})/50$, $\Delta\nu=c/50\mu\text{m}-c/80\mu\text{m}=94\text{GHz}$)あたり
 $100\mu\text{m} \times 100\mu\text{m} \times 8 \text{ pixels} \rightarrow$ 視野角 : $8 \times 10^{-8} \text{ sr}$
- Neutrino decay ($m_3 = 50 \text{ meV}$, $\tau_\nu = 1 \times 10^{14} \text{ yrs}$): $I_\nu=25\text{kJy/sr}$ @ $\lambda=50\mu\text{m}$
 $F_{\nu\gamma} = 25 \text{ kJy/sr} \times 8 \times 10^{-8} \text{ sr} \times \pi(15\text{cm}/2)^2 \times 94\text{GHz} = 3.3 \times 10^{-20} \text{ W}/8\text{pix}$
- Zodiacal emission: $I_\nu=8\text{MJy/sr}$ @ $\lambda=50\mu\text{m}$
 $F_{ZE} = 1.1 \times 10^{-17} \text{ W}/8\text{pix}$
- Δt 時間で F_{ZE} 積分した際の揺らぎ
エネルギー ϵ_γ のフォトン数揺らぎ起因 : $\epsilon_\gamma \sqrt{F_{ZE} \Delta t} / \epsilon_\gamma = \sqrt{\epsilon_\gamma F_{ZE} \Delta t}$
- 測定条件 , 検出器要件を決める不等式
$$NEP \times \sqrt{2\Delta t} < \sqrt{\epsilon_\gamma F_{ZE} \Delta t} < F_{\nu\gamma} \Delta t$$

 $\rightarrow \Delta t > 40\text{sec} (1\sigma), \Delta t > 200\text{sec} (2.2\sigma)$
 $\rightarrow NEP < 1.5 \times 10^{-19} \text{ W}/\sqrt{\text{Hz}}$ for $\Delta t=200\text{sec}$ with 8 pix