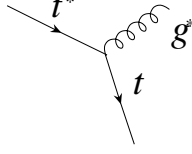


# Spin flip probability of Top Quark by Gluon Radiation

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Assume off-shell gluon radiation via tree diagram from off-shell top quark  $t^* \rightarrow t + g^*$



$$\mathfrak{M} \propto \varepsilon_\mu^* \bar{u}_t \gamma^\mu u_{t^*}$$

In the  $t^*$  rest frame, assume  $t$  flight direction to be  $z$  direction.

$p$ : magnitude of momenta of  $t$  and  $g^*$

$m_g$ : virtual mass of  $g^*$

We suppose a matrix element with index( $i, j$ , and  $k$ ) for spin states of  $t$ ,  $t^*$ , and  $g^*$

$$\mathfrak{M}_{ijk} \propto \varepsilon_\mu^*(k) \bar{u}_t(i) \gamma^\mu u_{t^*}(j)$$

$$u_t(i) = \sqrt{E_t + m_t} \begin{pmatrix} \phi_i \\ \frac{\sigma_3 p}{E_t + m_t} \phi_i \end{pmatrix} \quad E_t^2 = m_t^2 + p^2$$

$$u_{t^*}(j) = \sqrt{m_{t^*}} \begin{pmatrix} \psi_j \\ 0 \end{pmatrix} \quad m_{t^*} = \sqrt{m_t^2 + p^2} + \sqrt{m_{g^*}^2 + p^2}$$

$$\begin{aligned} \mathfrak{M}_{ijk} &\propto \varepsilon_\mu^*(k) \bar{u}_t(i) \gamma^\mu u_{t^*}(j) \\ &\propto \varepsilon_\mu^*(k) \begin{pmatrix} \phi_i^\dagger & -\phi_i^\dagger \frac{\sigma_3 p}{E_t + m_t} \end{pmatrix} \gamma^\mu \begin{pmatrix} \psi_j \\ 0 \end{pmatrix} \\ &\propto \varepsilon_\mu^*(k) \left\{ \phi_i^\dagger \psi_j; \phi_i^\dagger \frac{\sigma_3 p}{E_t + m_t} \sigma_{[123]} \psi_j \right\} \\ &\propto \varepsilon_\mu^*(k) \left\{ \phi_i^\dagger \psi_j; \frac{p}{E_t + m_t} \phi_i^\dagger \left[ \begin{pmatrix} & 1 \\ -1 & \end{pmatrix}, \begin{pmatrix} & -i \\ -i & \end{pmatrix}, 1 \right] \psi_j \right\} \end{aligned}$$

Breakdown as for spin states of  $t$  and  $t^*$

$$\mathfrak{M}_{\uparrow\uparrow k} \propto \varepsilon_\mu^*(k) \left\{ 1; \frac{p}{E_t + m_t} [0, 0, 1] \right\}$$

$$\mathfrak{M}_{\uparrow\downarrow k} \propto \varepsilon_\mu^*(k) \left\{ 0; \frac{p}{E_t + m_t} [1, -i, 0] \right\}$$

$$\mathfrak{M}_{\downarrow\uparrow k} \propto \varepsilon_\mu^*(k) \left\{ 0; \frac{p}{E_t + m_t} [-1, -i, 0] \right\}$$

$$\mathfrak{M}_{\downarrow\downarrow k} \propto \varepsilon_\mu^*(k) \left\{ 1; \frac{p}{E_t + m_t} [0, 0, 1] \right\}$$

Remind polarized vectors ( $\lambda$ : helicity) for spin-1 particles

$$\varepsilon^\mu(p, \lambda = \pm) = (0; \mp 1/\sqrt{2}, -i/\sqrt{2}, 0)$$

$$\varepsilon^\mu(p, \lambda = 0) = (p; 0, 0, E_{g^*})/m_{g^*}$$

Therefore

$$\varepsilon_\mu^*(-p, \lambda = \pm) = (0; \pm 1/\sqrt{2}, -i/\sqrt{2}, 0)$$

$$\varepsilon_\mu^*(-p, \lambda = 0) = (-p; 0, 0, -E_{g^*})/m_{g^*}$$

Finally we obtain the complete set of matrix elements for each spin state.

$$\begin{aligned}
\mathfrak{M}_{\uparrow\uparrow\pm} &= \mathfrak{M}_{\downarrow\downarrow\pm} = 0 & \mathfrak{M}_{\uparrow\uparrow 0} &= \mathfrak{M}_{\downarrow\downarrow 0} = -\left(1 + \frac{E_{g^*}}{E_t + m_t}\right)p/m_{g^*} \\
\mathfrak{M}_{\uparrow\downarrow\pm} &= \frac{p}{E_t + m_t}(\pm 1/\sqrt{2} + 1/\sqrt{2}) & \mathfrak{M}_{\uparrow\downarrow 0} &= 0 \\
\mathfrak{M}_{\downarrow\uparrow\pm} &= \frac{p}{E_t + m_t}(\mp 1/\sqrt{2} + 1/\sqrt{2}) & \mathfrak{M}_{\downarrow\uparrow 0} &= 0
\end{aligned}$$

Only the followings survive among above:

$$\begin{aligned}
\mathfrak{M}_{\uparrow\uparrow 0} &= \mathfrak{M}_{\downarrow\downarrow 0} = -\left(1 + \frac{E_{g^*}}{E_t + m_t}\right)p/m_{g^*} = -\left(\frac{E_t + E_{g^*} + m_t}{E_t + m_t}\right)p/m_{g^*} \\
&= -\frac{p}{E_t + m_t} \cdot \frac{m_{t^*} + m_t}{m_{g^*}} \\
\mathfrak{M}_{\uparrow\downarrow+} &= \mathfrak{M}_{\downarrow\uparrow-} = \frac{\sqrt{2}p}{E_t + m_t}
\end{aligned}$$

The spin flip probability of top quark by gluon radiation is given as

$$\begin{aligned}
P_{\text{flip}} &= \frac{|\mathfrak{M}_{\uparrow\downarrow+}|^2 + |\mathfrak{M}_{\downarrow\uparrow-}|^2}{|\mathfrak{M}_{\uparrow\downarrow+}|^2 + |\mathfrak{M}_{\downarrow\uparrow-}|^2 + |\mathfrak{M}_{\uparrow\uparrow 0}|^2 + |\mathfrak{M}_{\downarrow\downarrow 0}|^2} \\
&= \frac{2}{2 + \left(\frac{m_{t^*} + m_t}{m_{g^*}}\right)^2} \\
&= \frac{2m_{g^*}^2}{2m_{g^*}^2 + (m_{t^*} + m_t)^2}
\end{aligned}$$

[EOF]