

(c) Transport for London



Mind the Gap on IceCube

Cosmic neutrino spectrum and muon anomalous magnetic moment

Toshihiko Ota



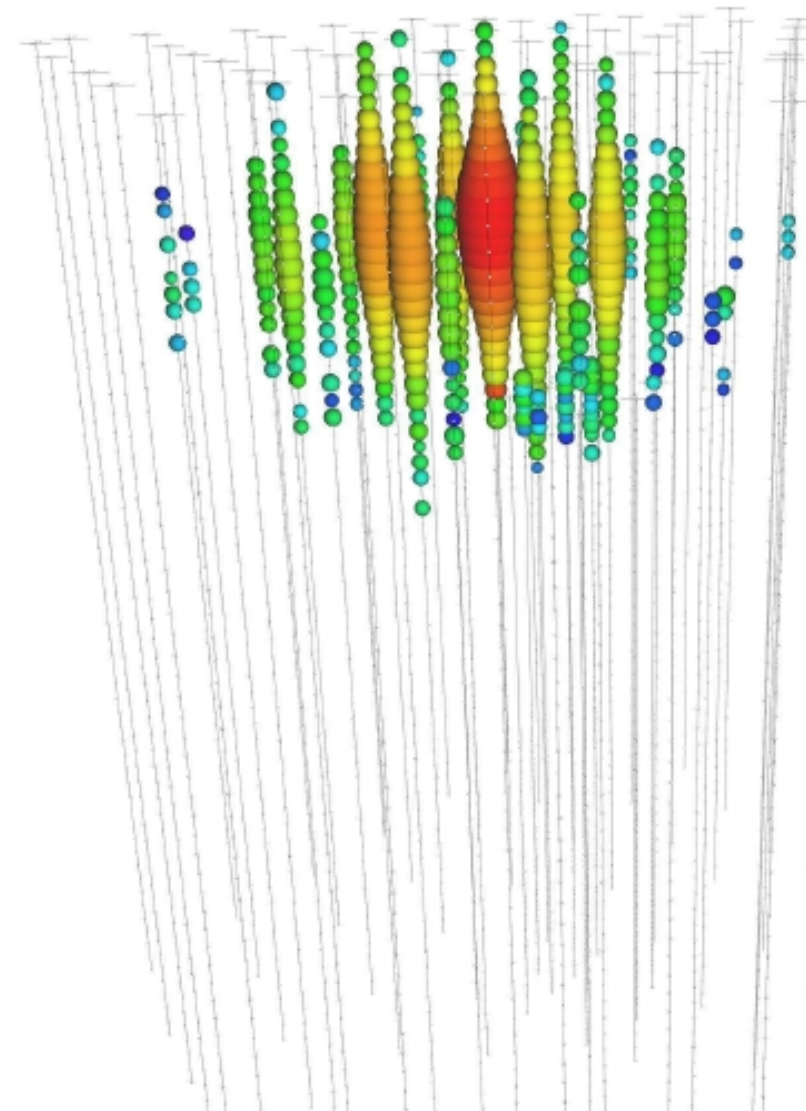
based on

T.Araki, Y.Konishi, F.Kaneko, TO, J.Sato, T.Shimomura

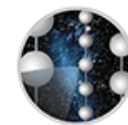
ArXiv.1409.4180v2

will be published in PRD

- PeV cosmic neutrino spectrum
IceCube collaboration PRL **113** (2014) 101101

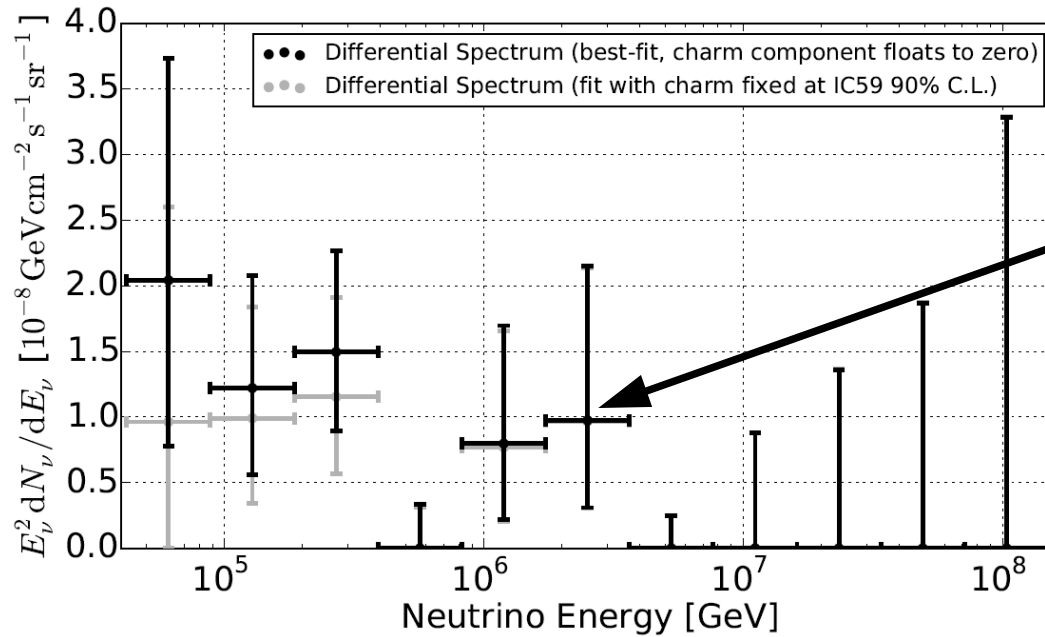


Event with the highest
deposit energy ~ 2 PeV

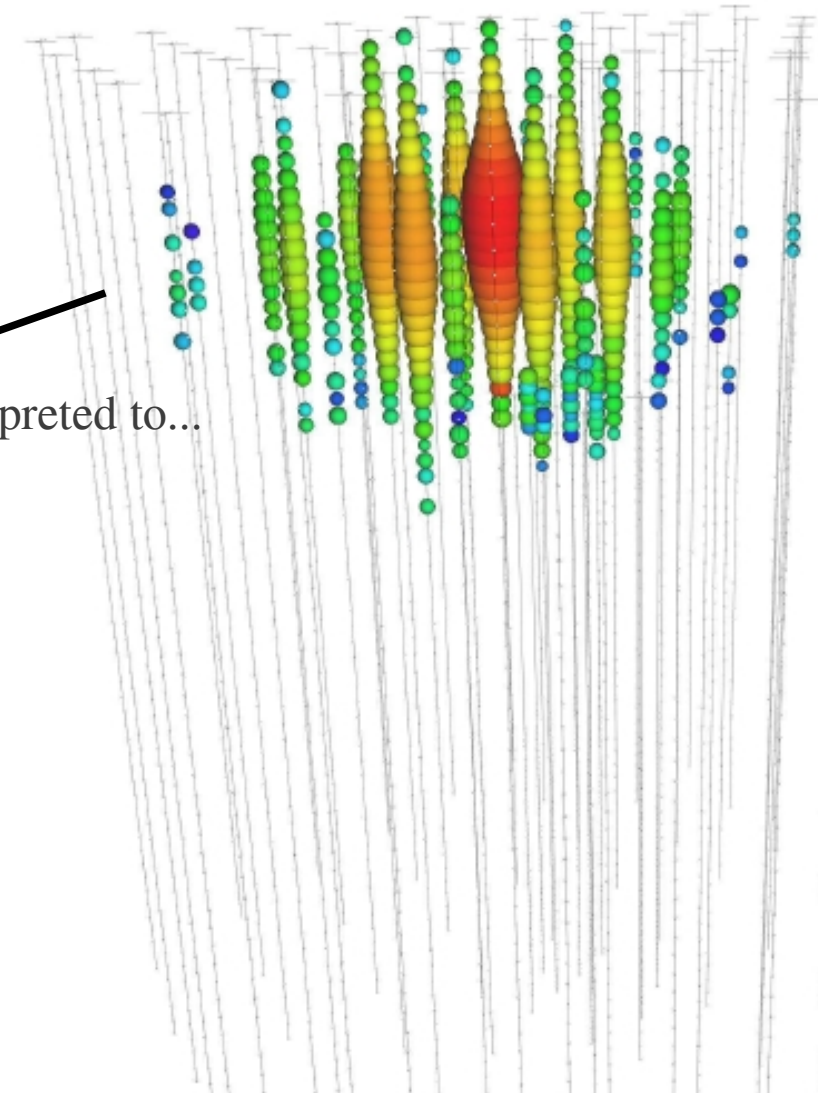


● PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101



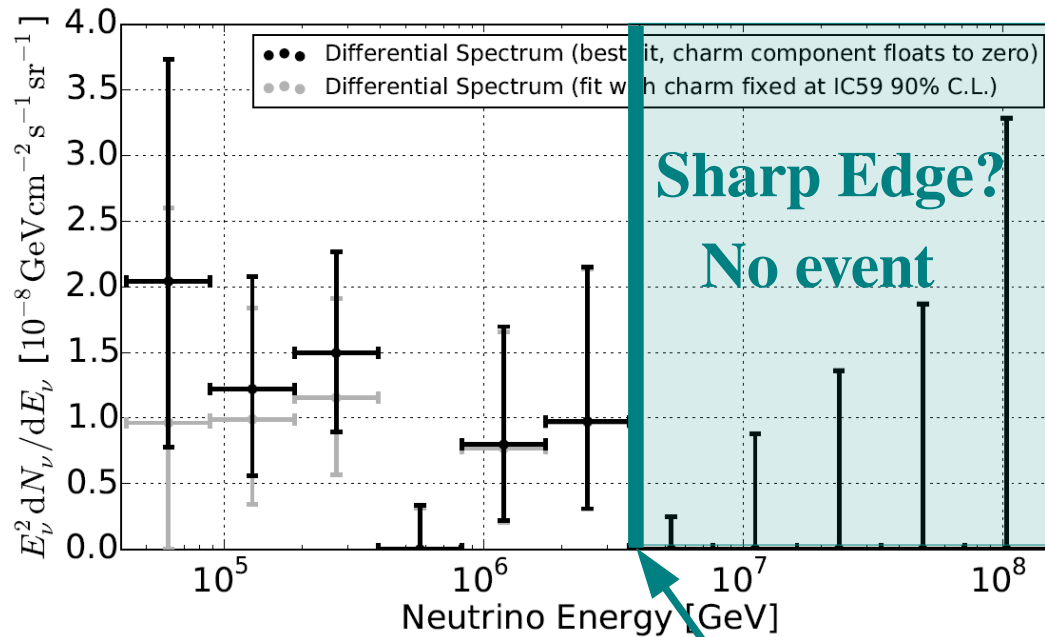
Interpreted to...



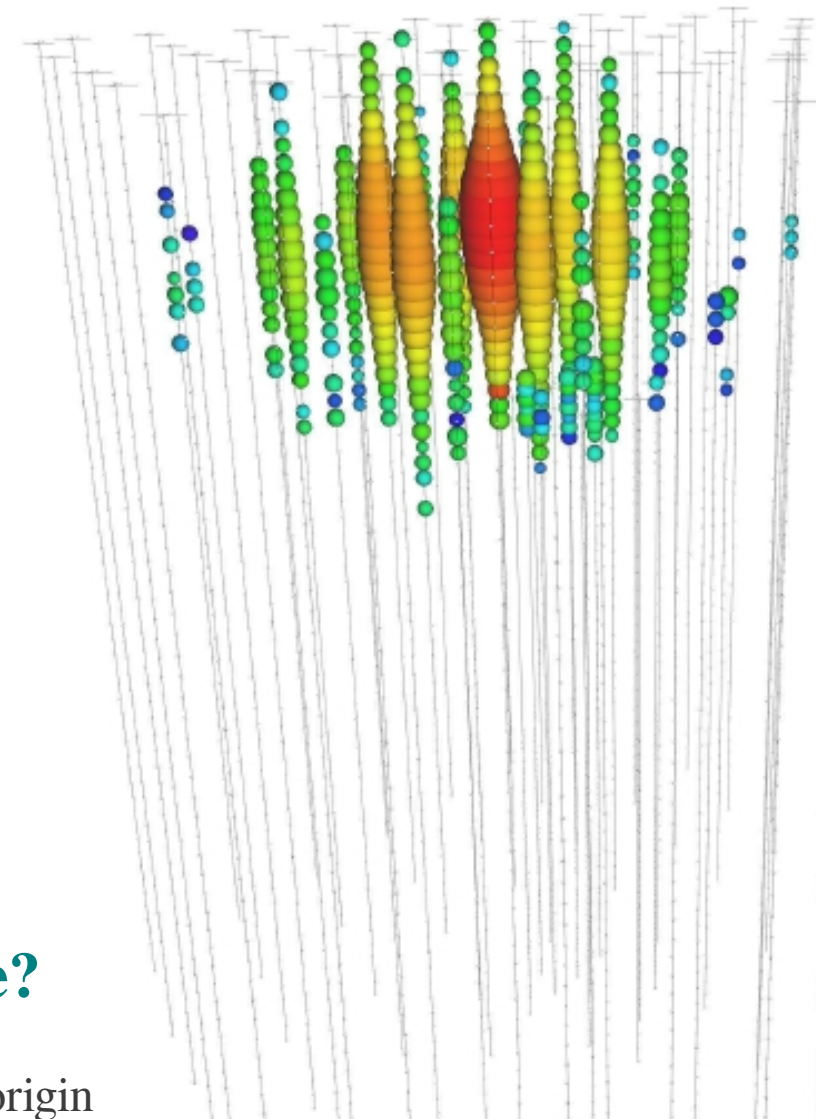
Event with the highest
deposit energy ~2 PeV



● PeV cosmic neutrino spectrum
IceCube collaboration PRL **113** (2014) 101101



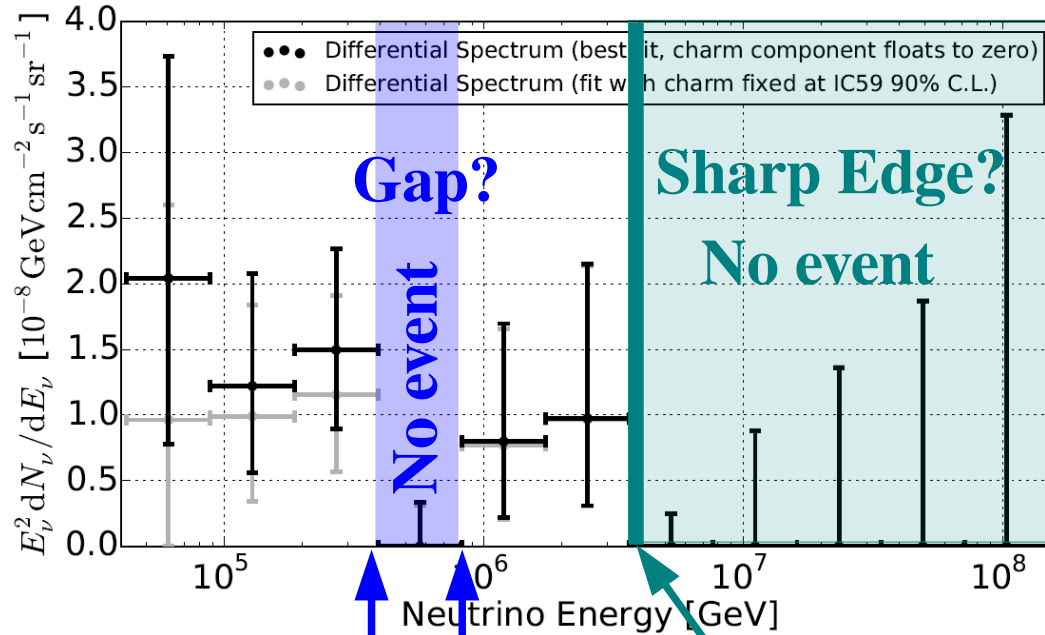
IceCube Edge?
at 3 PeV
may be astrophysical origin



Event with the highest
deposit energy ~2 PeV

● PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101

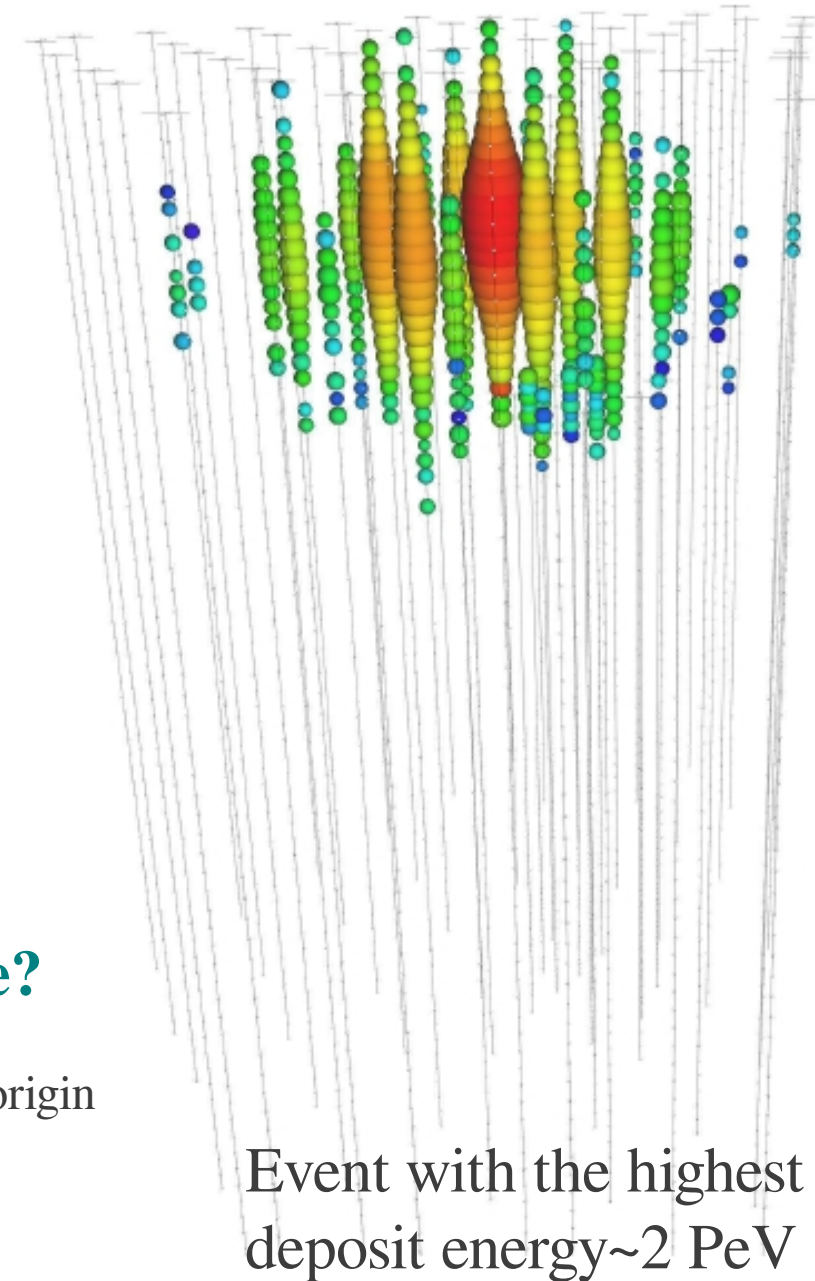


IceCube Gap

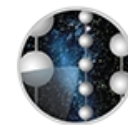
No event at 0.4-1 PeV

IceCube Edge?

at 3 PeV
may be astrophysical origin

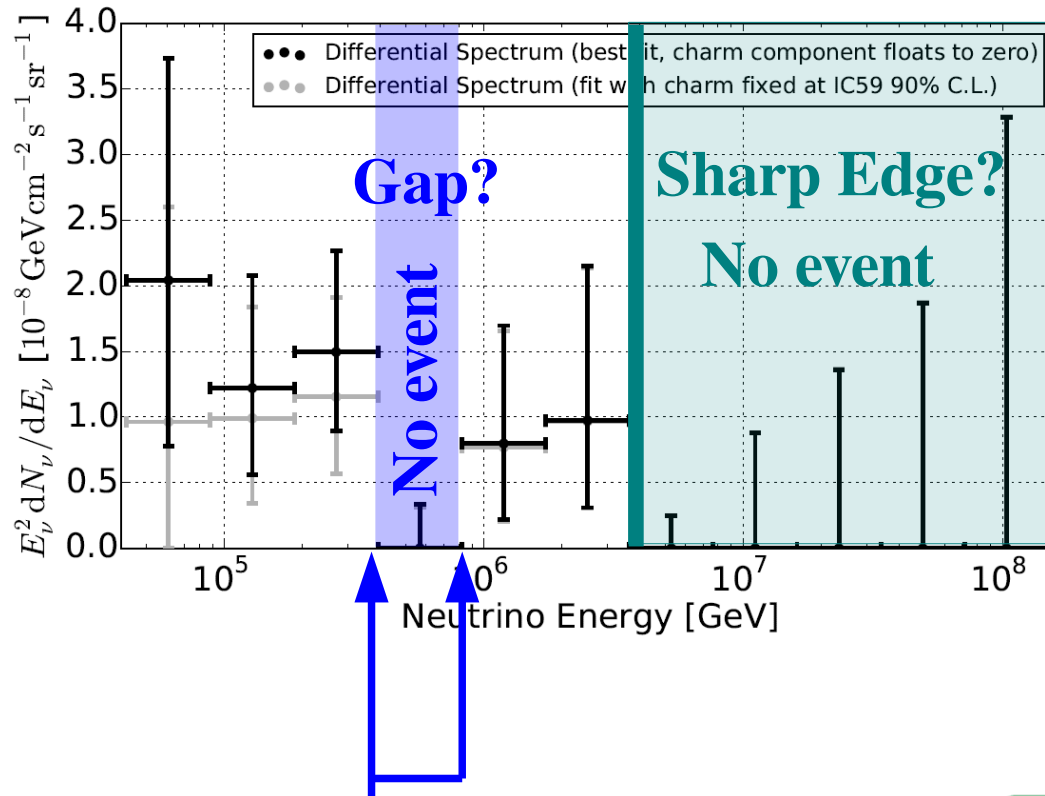


Event with the highest
deposit energy ~2 PeV



● PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101

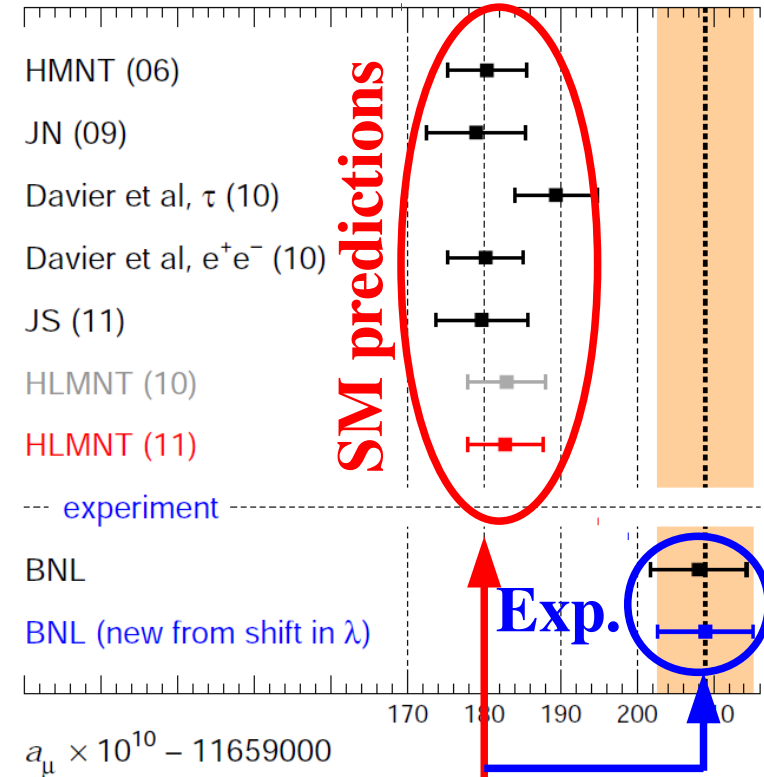


IceCube Gap

No event at 0.4-1 PeV

● Muon $g-2$

Hagiwara et al., J.Phys. **G38** (2011) 085003



$g_\mu - 2$ Gap

$$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}} = (26.1 \pm 8.0) \cdot 10^{-10} \quad (3.3\sigma)$$



New physics at the MeV scale
may explain both the gaps

1 IceCube gap

- Attenuation of cosmic neutrino by secret neutrino interaction
- Gauged leptonic force $L_\mu - L_\tau$ as secret interaction

2 Muon anomalous magnetic moment

- Gauged leptonic force as a contribution to $g-2$
- Constraints from colliders and neutrino trident process

3 A solution to the gaps

- Reproduction of IceCube gap → distance to the neutrino source
→ neutrino mass spectrum

If the **IceCube Gap** is explained by some **New Physics (NP)**...

● NP at Source: PeV Dark matter decay

Feldstein Kusenko Matsumoto Yanagida, PRD88 (2013) 015004. Zabala PRD89 (2014) 123514.

Ibarra Tran Weniger Int.J.Mod.Phys. A28 (2013) 1330040.

Esmaili Serpico JCAP 1311 (2013) 054, Esmaili Kang Serpico, 1410.5979.

Ema Jinno Moroi PLB733(2014) 120, JHEP 1410 (2014) 150. Rott Kohri Park 1408.3799.

Higaki Kitano Sato JHEP 1407(2014) 044. Fong Minakata Panes Zukanovich-Funchal 1411.5318.

➔ ● NP in Propagation: Scattering with CNB with a MeV mediator

As an effective int.: Ng Beacom PRD90 (2014) 065035, Ioka Murase PTEP 6 (2014) 061E01

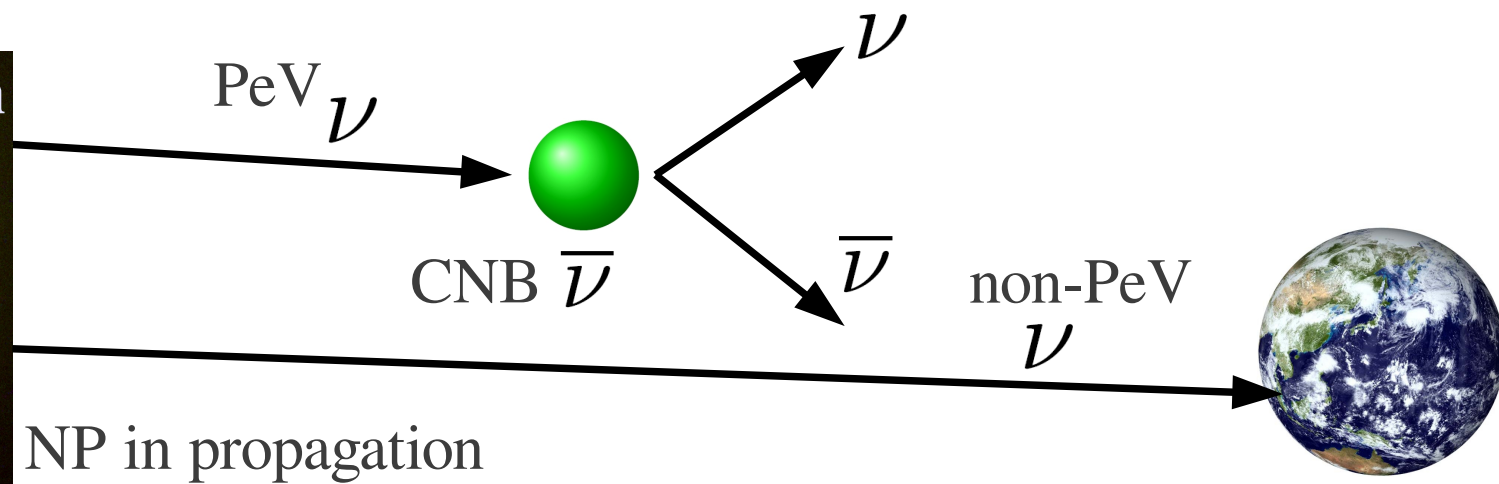
With neutrino mass model: Ibe Kaneta PRD90 (2014) 053011, Blum Hook Murase 1408.3799

● NP at Detection: CC int. mediated by a new TeV field

Barger Keung PLB727 (2013) 190...

NASA:Hubble heritage team

Continuous spectrum
@source



- In this talk, we pursue the possibility of

➔ **NP in propagation**, namely **Resonant scattering with CNB**

- We set **3 assumptions** for cosmic neutrino sources

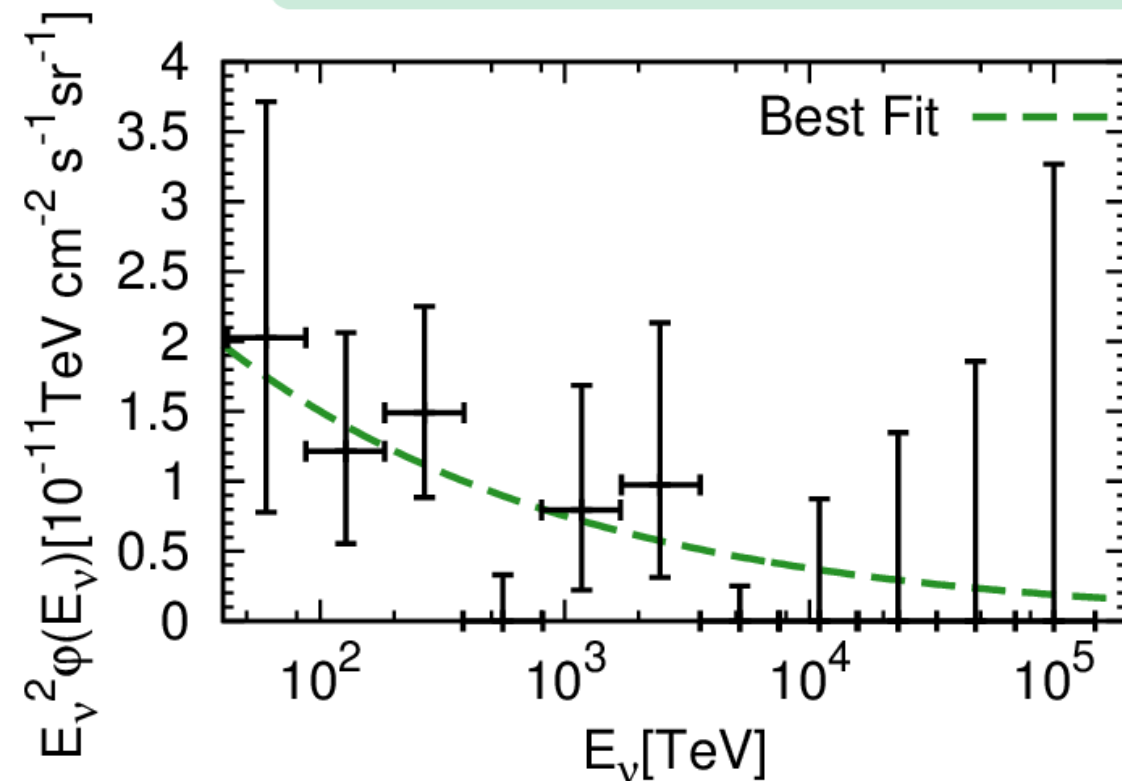


- In this talk, we pursue the possibility of

➔ **NP in propagation**, namely **Resonant scattering with CNB**

- We set **3 assumptions** for cosmic neutrino sources

1 Continuous (power-law) spectrum



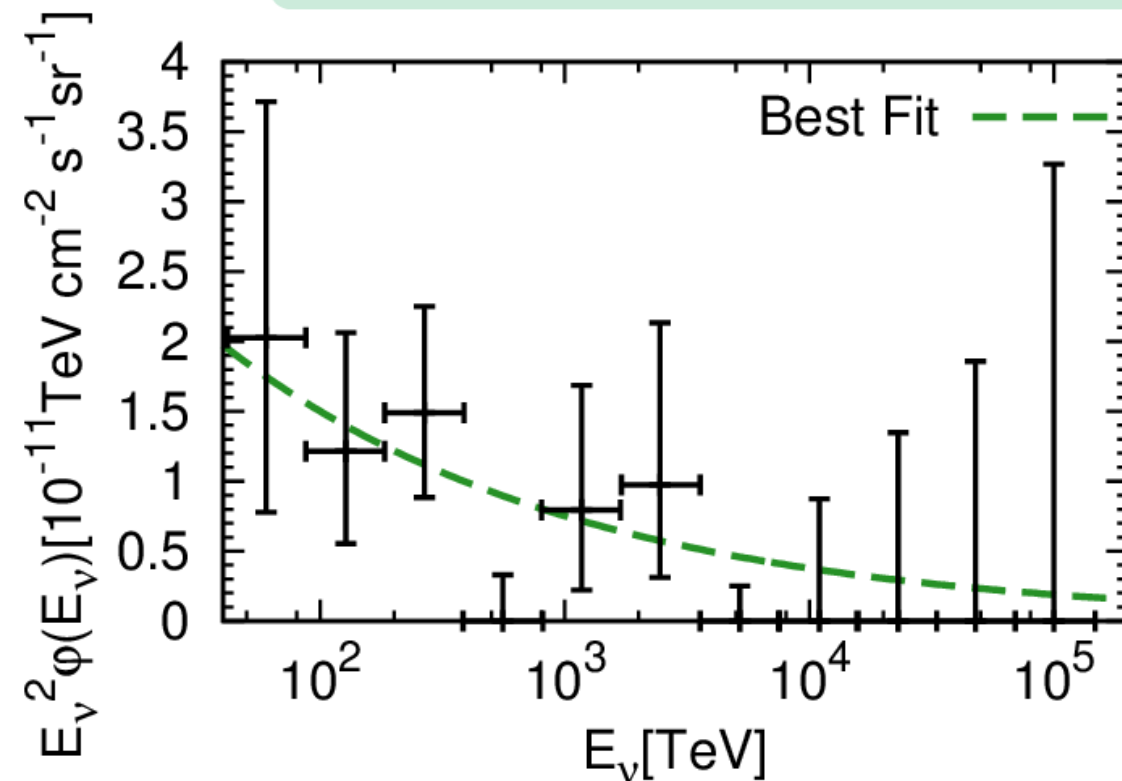
- 1 The spectrum shown with the **green curve** is reproduced, **if there is no NP**.

- In this talk, we pursue the possibility of

➔ **NP in propagation**, namely **Resonant scattering with CNB**

- We set **3 assumptions** for cosmic neutrino sources

- Continuous (power-law) spectrum
- Flavour ratio $\sim 1:1:1$ after leaving sources



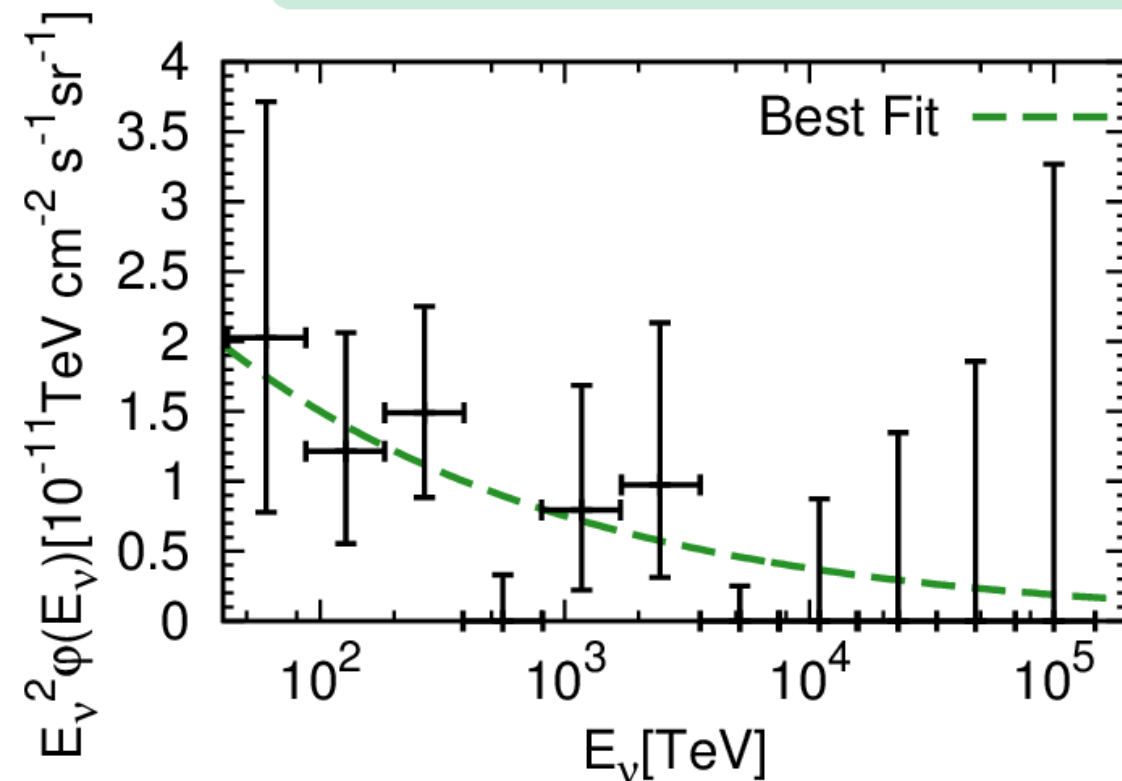
- The spectrum shown with the **green curve** is reproduced, **if there is no NP**.
- is not crucial. We will see...

- In this talk, we pursue the possibility of

➔ **NP in propagation**, namely **Resonant scattering with CNB**

- We set **3 assumptions** for cosmic neutrino sources

- Continuous (power-law) spectrum
- Flavour ratio $\sim 1:1:1$ after leaving sources
- Sources distribute around a particular redshift z_{source}



1 The spectrum shown with the **green curve** is reproduced, **if there is no NP.**

2 is not crucial. We will see...

3 for simplicity.

→ z -dependence of source distribution
 e.g., The star-formation rate has a peak at $z = 1 \sim 2$.

“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

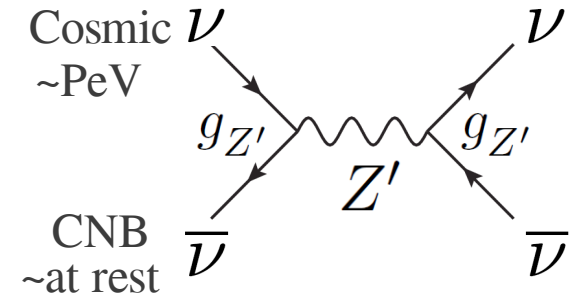


“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

● **Resonance condition**

$$s \simeq 2E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$



Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof

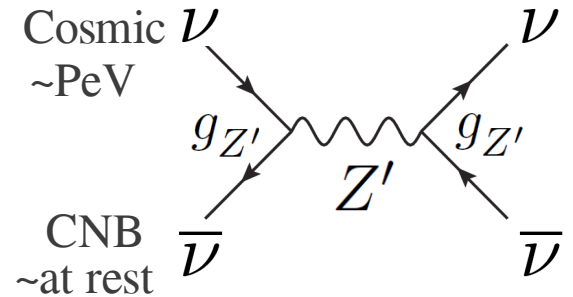
“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

● Resonance condition

$$s \simeq 2 E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$

~sub-PeV ~0.1eV



Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof

“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

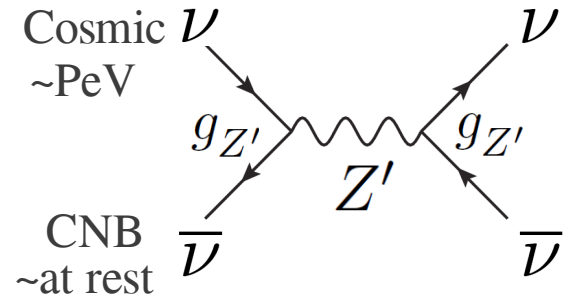
The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

● Resonance condition

$$s \simeq 2 E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$

\sim sub-PeV
 \sim 0.1eV

→ $M_{Z'} \sim$ MeV
NP @MeV scale



Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof

“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

● Resonance condition

$$s \simeq 2 E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$

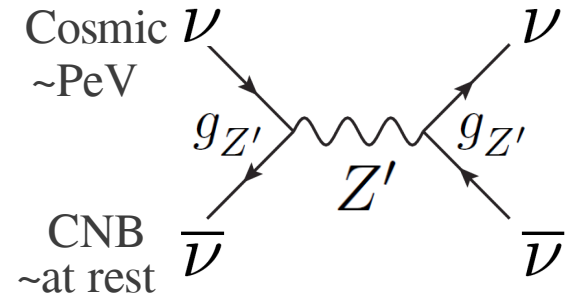
~sub-PeV ~0.1eV

→ $M_{Z'} \sim \text{MeV}$
NP @MeV scale

Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof

● How far can cosmic neutrinos travel in CNB? → Mean free path: λ

$$\lambda \simeq 1/n_{\text{CNB}} \sigma_{\text{@Res.}}$$



“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

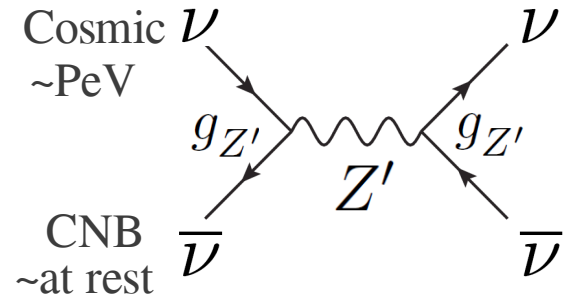
● Resonance condition

$$s \simeq 2 E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$

~sub-PeV ~0.1eV

→ $M_{Z'} \sim \text{MeV}$
NP @MeV scale

Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof



● How far can cosmic neutrinos travel in CNB? → Mean free path: λ

$$\lambda \simeq 1/n_{\text{CNB}} \sigma_{\text{@Res.}} \stackrel{!}{=} \text{Gpc}$$

Extra galactic source

“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

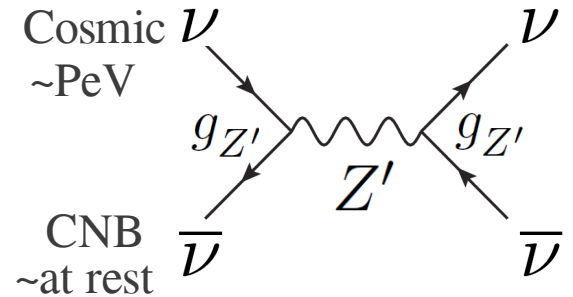
● Resonance condition

$$s \simeq 2 E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$

~sub-PeV ~0.1eV

→ $M_{Z'} \sim \text{MeV}$
NP @MeV scale

Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof



● How far can cosmic neutrinos travel in CNB? → Mean free path: λ

$$\lambda \simeq 1/n_{\text{CNB}} \sigma_{\text{@Res.}} \stackrel{!}{=} \text{Gpc}$$

Extra galactic source
 → $\sigma_{\text{@Res.}} \simeq 10^{-30} \text{ [cm}^2\text{]}$

“A sharp gap” → “Cosmic neutrino with a particular energy is scattered off”

The key idea is...“Resonant interaction with Cosmic Neutrino Background (CNB)”

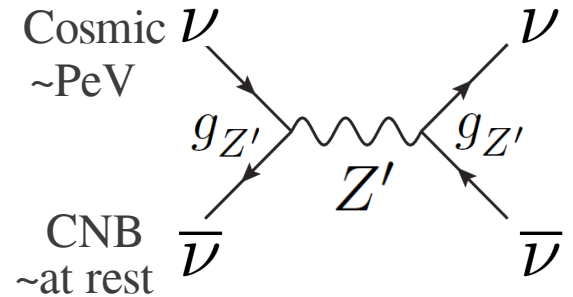
● Resonance condition

$$s \simeq 2 E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$

~sub-PeV ~0.1eV

→ $M_{Z'} \sim \text{MeV}$
NP @MeV scale

Why CNB? → $n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof



● How far can cosmic neutrinos travel in CNB? → Mean free path: λ

$$\lambda \simeq 1/n_{\text{CNB}} \sigma_{\text{@Res.}} \stackrel{!}{=} \text{Gpc}$$

Extra galactic source

→ $\sigma_{\text{@Res.}} \simeq 10^{-30} \text{ [cm}^2\text{]}$

IceCube Gap requires

$M_{Z'} \sim \text{MeV}, \quad \sigma_{\text{@Res.}} \gtrsim 10^{-30} \text{ [cm}^2\text{]}.$

Let us calculate the cross-section in a particular model...

- Gauged $U(1) L_\mu - L_\tau$ force as **a benchmark model**

Charge assignments $Y(L_\mu) = +1, Y(L_\tau) = -1,$
 $Y(\mu_R) = +1, Y(\tau_R) = -1, Y(\text{others}) = 0.$

$$\begin{aligned} \mathcal{L}_{L_\mu - L_\tau} = & g_{Z'} \bar{L}_\mu \gamma^\rho L_\mu Z'_\rho - g_{Z'} \bar{L}_\tau \gamma^\rho L_\tau Z'_\rho \\ & + g_{Z'} \bar{\mu}_R \gamma^\rho \mu_R Z'_\rho - g_{Z'} \bar{\tau}_R \gamma^\rho \tau_R Z'_\rho \end{aligned}$$

- Gauged $U(1) L_\mu - L_\tau$ force as **a benchmark model**

Charge assignments $Y(L_\mu) = +1, Y(L_\tau) = -1,$
 $Y(\mu_R) = +1, Y(\tau_R) = -1, Y(\text{others}) = 0.$

$$\begin{aligned} \mathcal{L}_{L_\mu - L_\tau} &= g_{Z'} \bar{L}_\mu \gamma^\rho L_\mu Z'_\rho - g_{Z'} \bar{L}_\tau \gamma^\rho L_\tau Z'_\rho \\ &\quad + g_{Z'} \bar{\mu}_R \gamma^\rho \mu_R Z'_\rho - g_{Z'} \bar{\tau}_R \gamma^\rho \tau_R Z'_\rho \\ &= g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho + g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{l}_\alpha \gamma^\rho l_\beta Z'_\rho \end{aligned}$$

- Gauged $U(1) L_\mu - L_\tau$ force as **a benchmark model**

Charge assignments $Y(L_\mu) = +1, Y(L_\tau) = -1,$
 $Y(\mu_R) = +1, Y(\tau_R) = -1, Y(\text{others}) = 0.$

$$\begin{aligned} \mathcal{L}_{L_\mu - L_\tau} &= g_{Z'} \bar{L}_\mu \gamma^\rho L_\mu Z'_\rho - g_{Z'} \bar{L}_\tau \gamma^\rho L_\tau Z'_\rho \\ &\quad + g_{Z'} \bar{\mu}_R \gamma^\rho \mu_R Z'_\rho - g_{Z'} \bar{\tau}_R \gamma^\rho \tau_R Z'_\rho \\ &= \underline{g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho} + \underline{g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{\ell}_\alpha \gamma^\rho \ell_\beta Z'_\rho} \end{aligned}$$

Neutrino secret int.

Coupling in mass eigenbasis

$$g_{ij} = g_{Z'} (U_{\text{PMNS}}^\dagger)_{i\alpha} \text{diag}(0, 1, -1)_{\alpha\beta} (U_{\text{PMNS}})_{\beta j}$$

* Cosmic neutrino is produced as a flavour eigenstate = a coherent sum of mass eigenstates.
 But the coherence is lost in its travel.

Constrained! but...

Contribute to muon g-2

We discuss it in **Sec. 2**

- Gauged $U(1) L_\mu - L_\tau$ force as **a benchmark model**

Charge assignments $Y(L_\mu) = +1, Y(L_\tau) = -1,$
 $Y(\mu_R) = +1, Y(\tau_R) = -1, Y(\text{others}) = 0.$

$$\begin{aligned} \mathcal{L}_{L_\mu - L_\tau} &= g_{Z'} \bar{L}_\mu \gamma^\rho L_\mu Z'_\rho - g_{Z'} \bar{L}_\tau \gamma^\rho L_\tau Z'_\rho \\ &\quad + g_{Z'} \bar{\mu}_R \gamma^\rho \mu_R Z'_\rho - g_{Z'} \bar{\tau}_R \gamma^\rho \tau_R Z'_\rho \\ &= \underline{g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho} + \underline{g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{\ell}_\alpha \gamma^\rho \ell_\beta Z'_\rho} \end{aligned}$$

Neutrino secret int.

Constrained! but...

Contribute to muon g-2

Coupling in mass eigenbasis

$$g_{ij} = g_{Z'} (U_{\text{PMNS}}^\dagger)_{i\alpha} \text{diag}(0, 1, -1)_{\alpha\beta} (U_{\text{PMNS}})_{\beta j}$$

We discuss it in **Sec. 2**

* Cosmic neutrino is produced as a flavour eigenstate = a coherent sum of mass eigenstates.
 But the coherence is lost in its travel.

- Motivated from...**

- (almost) Maximal mixing Choubey Rodejohann Eur.Phys.J C40 (2005) 259

- Gauge anomaly free Foot Mod.Phys.A6 (1991) 527, He et al., PRD43 (1990) R22

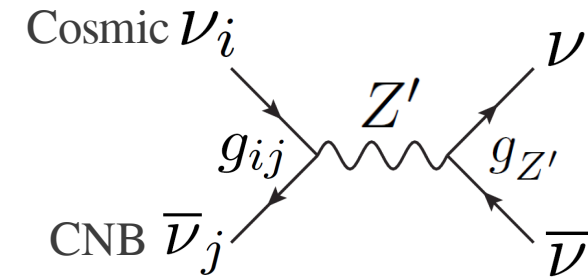
- In this talk, we do not go into the details of the spontaneous breaking of the $L_\mu - L_\tau$ sym. \longrightarrow

Model parameters

$g_{Z'}$ and $M_{Z'}$

● Cross-section of the neutrino scattering proc.

$$\sigma(\nu_i \bar{\nu}_j \rightarrow \nu \bar{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$



● Decay rate

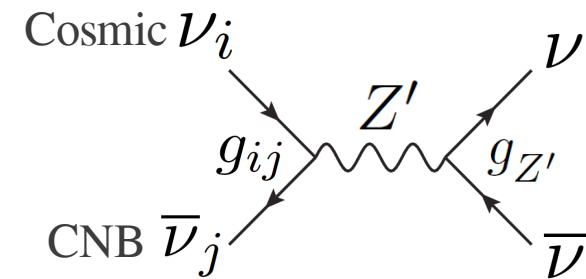
$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi}$$

● Cross-section @ Resonance

$$\sigma_{@Res.} = \frac{4\pi |g_{ij}|^2}{M_{Z'}^2} \delta \left(1 - \frac{M_{Z'}^2}{s} \right)$$

● Cross-section of the neutrino scattering proc.

$$\sigma(\nu_i \bar{\nu}_j \rightarrow \nu \bar{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$



● Decay rate

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi}$$

● Cross-section @ Resonance

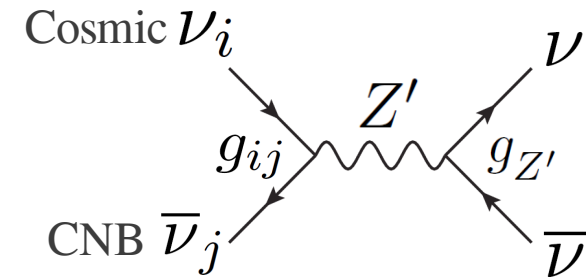
For IceCube Gap

$$\sigma_{@Res.} = \frac{4\pi |g_{ij}|^2}{M_{Z'}^2} \delta\left(1 - \frac{M_{Z'}^2}{s}\right) \stackrel{!}{=} 10^{-30} [\text{cm}^2]$$

$M_{Z'} \sim \text{MeV}$

● Cross-section of the neutrino scattering proc.

$$\sigma(\nu_i \bar{\nu}_j \rightarrow \nu \bar{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$



● Decay rate

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi}$$

● Cross-section @ Resonance

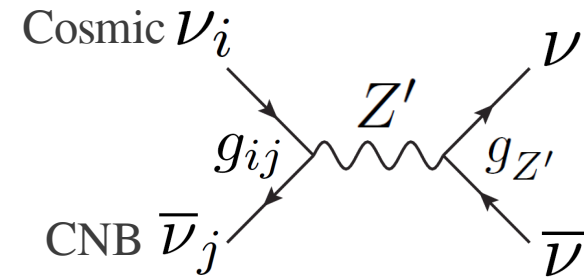
For IceCube Gap

$$\sigma_{@Res.} = \frac{4\pi |g_{ij}|^2}{M_{Z'}^2} \delta\left(1 - \frac{M_{Z'}^2}{s}\right) \stackrel{!}{=} 10^{-30} [\text{cm}^2]$$

$M_{Z'} \sim \text{MeV} \rightarrow g_{Z'} \simeq 10^{-3}$

● Cross-section of the neutrino scattering proc.

$$\sigma(\nu_i \bar{\nu}_j \rightarrow \nu \bar{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$



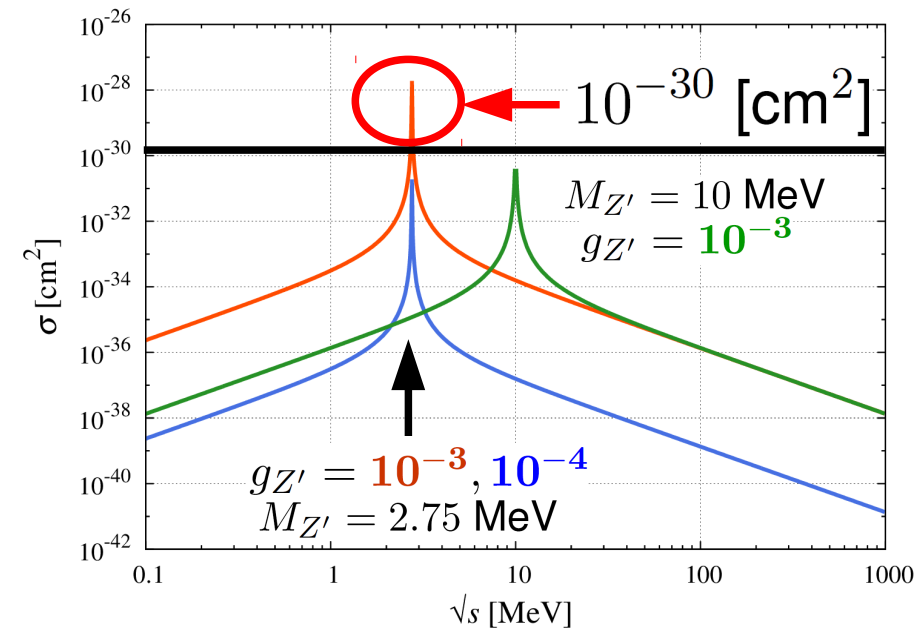
● Decay rate

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi}$$

● Cross-section @ Resonance

$$\sigma_{@Res.} = \frac{4\pi |g_{ij}|^2}{M_{Z'}^2} \delta\left(1 - \frac{M_{Z'}^2}{s}\right) \stackrel{!}{=} 10^{-30} [\text{cm}^2]$$

For IceCube Gap
 $M_{Z'} \sim \text{MeV} \rightarrow g_{Z'} \simeq 10^{-3}$



IceCube Gap requires

$$M_{Z'} \sim \text{MeV}, \quad g_{Z'} \gtrsim 10^{-4}.$$

- The width might be **too narrow** for the **IceCube Gap** (0.4-1PeV).
- We can ask the help to m_ν and z

→ **Sec. 3**

Before going into the details of the cosmic neutrino spectrum, let's check muon g-2.

1 IceCube gap

- Attenuation of cosmic neutrino by secret neutrino interaction
- Gauged leptonic force $L_\mu - L_\tau$ as secret interaction

2 Muon anomalous magnetic moment

- Gauged leptonic force as a contribution to $g-2$
- Constraints from colliders and neutrino trident process

3 A solution to the gaps

- Reproduction of IceCube gap → distance to the neutrino source
→ neutrino mass spectrum

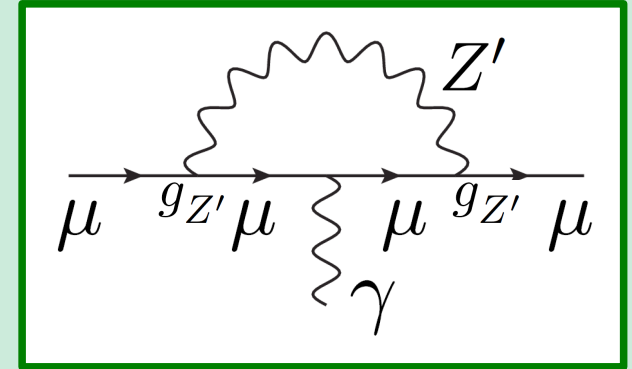
Z' contribution to $g_\mu - 2$

$$\mathcal{L}_{L_\mu - L_\tau} = \underbrace{g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho}_{\text{Neutrino secret int.}} + \underbrace{g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{\ell}_\alpha \gamma^\rho \ell_\beta Z'_\rho}_{\text{Contribute to muon } g-2}$$

Z' contribution to $g_\mu - 2$

$$\mathcal{L}_{L_\mu - L_\tau} = g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho + \underline{g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{l}_\alpha \gamma^\rho l_\beta Z'_\rho}$$

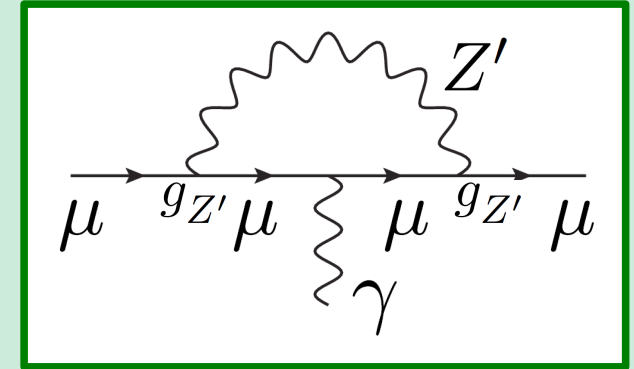
- $M_{Z'} \gg m_\mu \rightarrow \Delta a_\mu^{Z'} = \frac{g_{Z'}^2}{12\pi^2} \frac{m_\mu^2}{M_{Z'}^2}$
- $M_{Z'} \ll m_\mu \rightarrow \Delta a_\mu^{Z'} = \frac{g_{Z'}^2}{8\pi^2}$



Z' contribution to $g_\mu - 2$

$$\mathcal{L}_{L_\mu - L_\tau} = g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho + \underline{g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{\ell}_\alpha \gamma^\rho \ell_\beta Z'_\rho}$$

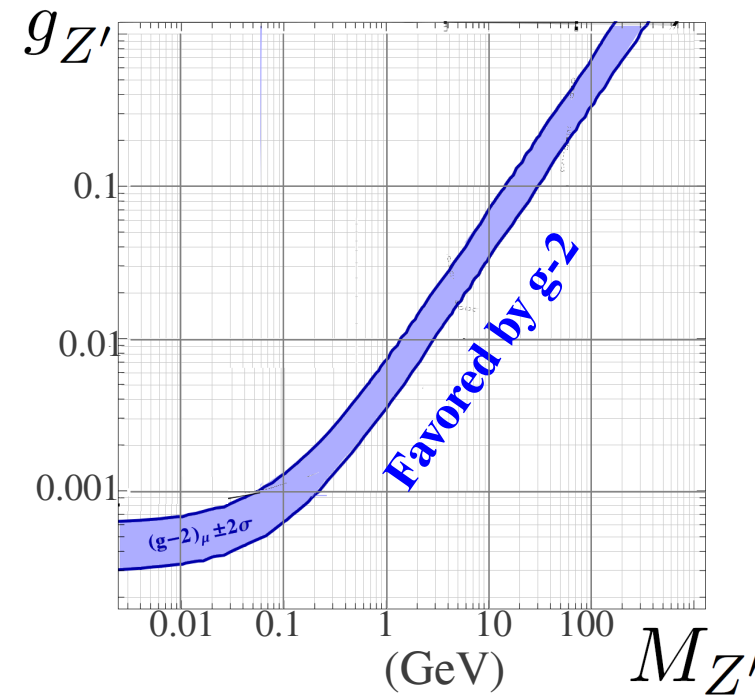
- $M_{Z'} \gg m_\mu \rightarrow \Delta a_\mu^{Z'} = \frac{g_{Z'}^2}{12\pi^2} \frac{m_\mu^2}{M_{Z'}^2}$
- $M_{Z'} \ll m_\mu \rightarrow \Delta a_\mu^{Z'} = \frac{g_{Z'}^2}{8\pi^2}$



$g_\mu - 2$ Gap

$$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}} = (26.1 \pm 8.0) \cdot 10^{-10} \quad (3.3\sigma)$$

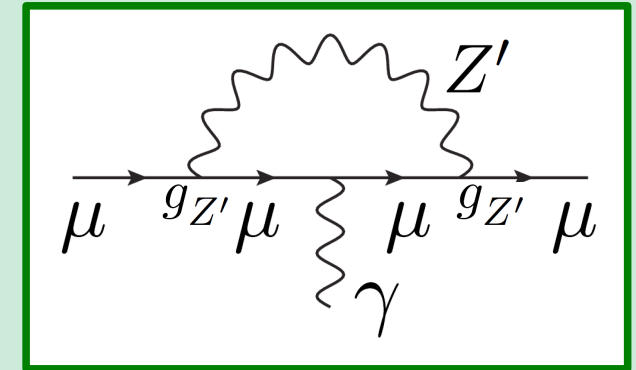
\rightarrow We need $\Delta a_\mu^{\text{NP}} \simeq (20-30) \cdot 10^{-10}$



Z' contribution to $g_\mu - 2$

$$\mathcal{L}_{L_\mu - L_\tau} = g_{ij} \bar{\nu}_i \gamma^\rho P_L \nu_j Z'_\rho + \underline{g_{Z'} \text{diag}(0, 1, -1)_{\alpha\beta} \bar{\ell}_\alpha \gamma^\rho \ell_\beta Z'_\rho}$$

- $M_{Z'} \gg m_\mu \rightarrow \Delta a_\mu^{Z'} = \frac{g_{Z'}^2}{12\pi^2} \frac{m_\mu^2}{M_{Z'}^2}$
- $M_{Z'} \ll m_\mu \rightarrow \Delta a_\mu^{Z'} = \frac{g_{Z'}^2}{8\pi^2}$



$g_\mu - 2$ Gap

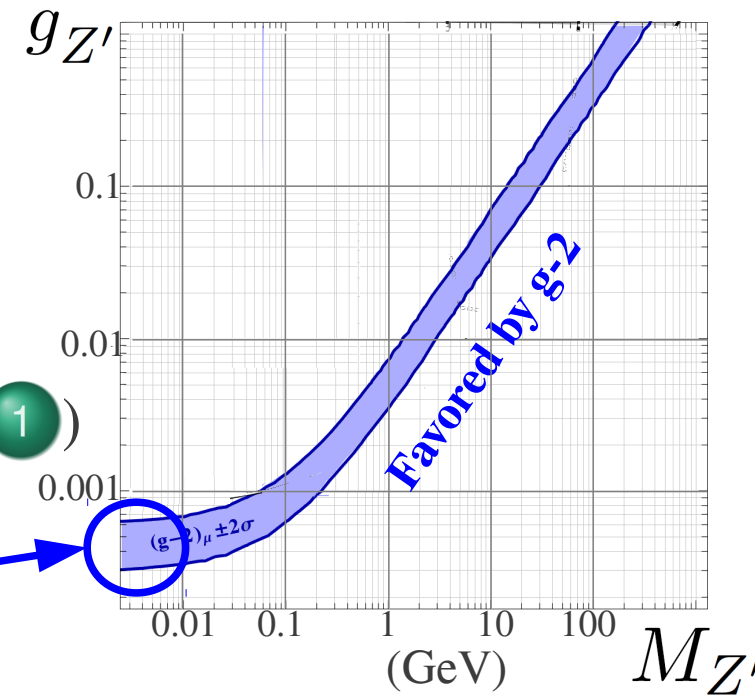
$$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}} = (26.1 \pm 8.0) \cdot 10^{-10} \quad (3.3\sigma)$$

→ We need $\Delta a_\mu^{\text{NP}} \simeq (20-30) \cdot 10^{-10}$

- Let me remind (back-of-the envelope calc. in **Sec. 1**)

IceCube Gap requires

$$M_{Z'} \sim \text{MeV}, \quad g_{Z'} \gtrsim 10^{-4}.$$



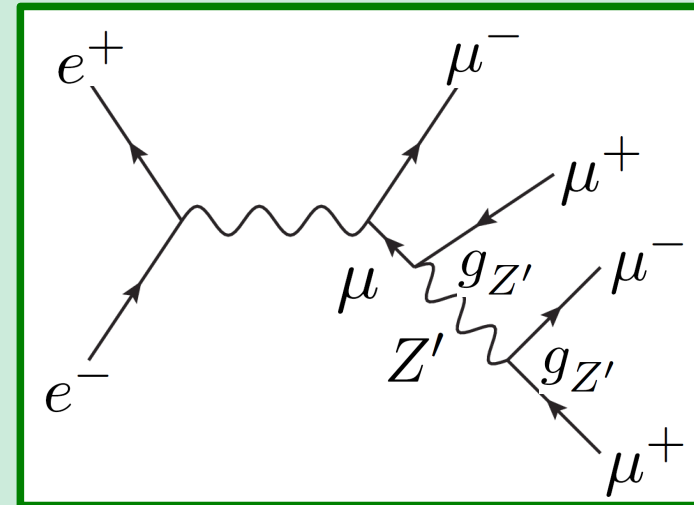
Collider bounds Harigaya et al., JHEP 1403 (2014) 105.

● Process: $e^+e^- \rightarrow 4\mu$

$$PP(P\bar{P}) \rightarrow 4\mu/2\mu2\tau$$

only constrain relatively heavy Z'

→ LEP, LHC: $g_{Z'} \lesssim 0.1$ at $M_{Z'} \simeq 100$ GeV



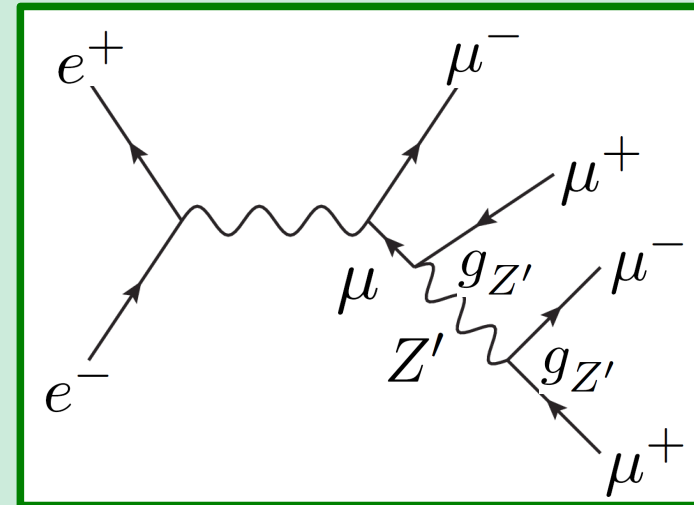
Collider bounds Harigaya et al., JHEP 1403 (2014) 105.

● Process: $e^+e^- \rightarrow 4\mu$

$$PP(P\bar{P}) \rightarrow 4\mu/2\mu2\tau$$

only constrain relatively heavy Z'

→ LEP, LHC: $g_{Z'} \lesssim 0.1$ at $M_{Z'} \simeq 100$ GeV



Rare meson decays Lessa and Peres, PRD75 (2007) 094001

● Process: $\pi^+ / K^+ \rightarrow \mu^+ \nu_\mu Z'$

Bound from Kaon decay → $g_{Z'} \lesssim 0.01$ at $M_{Z'} \sim \text{MeV}$

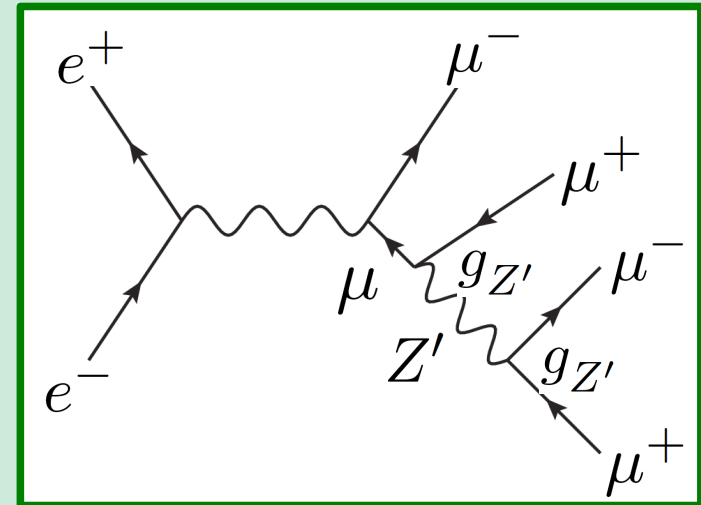
Collider bounds Harigaya et al., JHEP 1403 (2014) 105.

● Process: $e^+e^- \rightarrow 4\mu$

$$PP(P\bar{P}) \rightarrow 4\mu/2\mu 2\tau$$

only constrain relatively heavy Z'

→ LEP, LHC: $g_{Z'} \lesssim 0.1$ at $M_{Z'} \simeq 100$ GeV



Rare meson decays Lessa and Peres, PRD75 (2007) 094001

● Process: $\pi^+ / K^+ \rightarrow \mu^+ \nu_\mu Z'$

Bound from Kaon decay → $g_{Z'} \lesssim 0.01$ at $M_{Z'} \sim \text{MeV}$

● The most relevant bound from lab. experiments is

Neutrino trident process in neutrino-nucleon scattering

Altmannshofer Gori Pospelov Yavin, PRL 113 (2014) 091801

● Bounds from CMB, BBN, and also from SN1987A → References in Ng Beacom

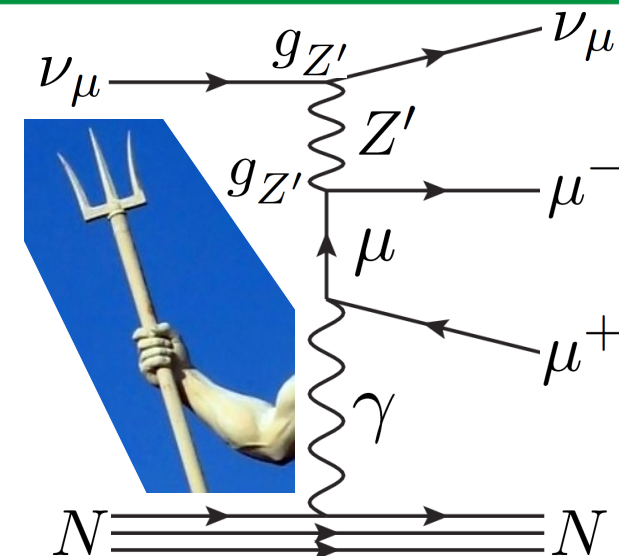
- **Neutrino trident process**

in neutrino-nucleon scattering events

- Available **data** reported by CCFR in 1991!

37 events (± 12.4)

CCFR collaboration, PRL **66** (1991) 3117
excavated recently (only cited 18 times)*



Altmannshofer et al., PRL **113** (2014) 091801

*The trident process must be recorded on the hard disks of the near detectors in modern oscillation experiments. They should be opened!

● Neutrino trident process

in neutrino-nucleon scattering events

- Available **data** reported by CCFR in 1991!

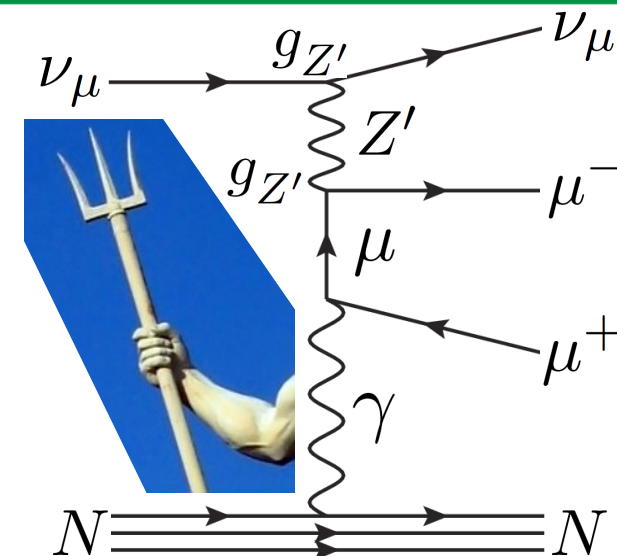
37 events (± 12.4)

CCFR collaboration, PRL **66** (1991) 3117
excavated recently (only cited 18 times)*

- Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

→ **Consistent** → constrains $g_{Z'}$ and $M_{Z'}$



Altmannshofer et al., PRL **113** (2014) 091801

*The trident process must be recorded on the hard disks of the near detectors in modern oscillation experiments. They should be opened!

● **Neutrino trident process**

in neutrino-nucleon scattering events

- Available **data** reported by CCFR in 1991!

37 events (± 12.4)

CCFR collaboration, PRL **66** (1991) 3117
excavated recently (only cited 18 times)*

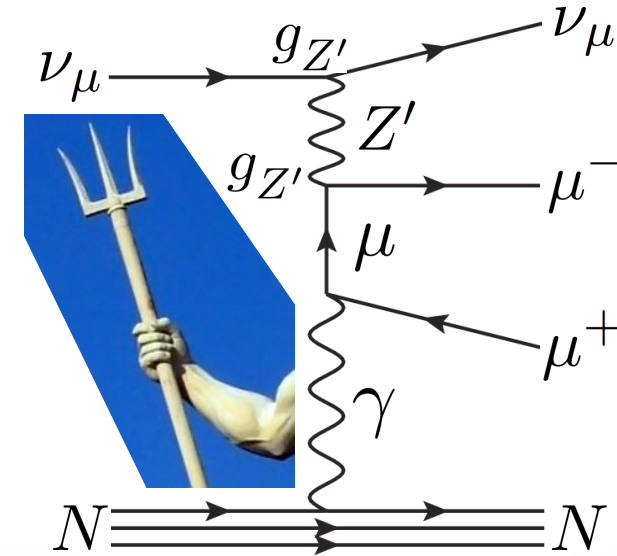
- Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

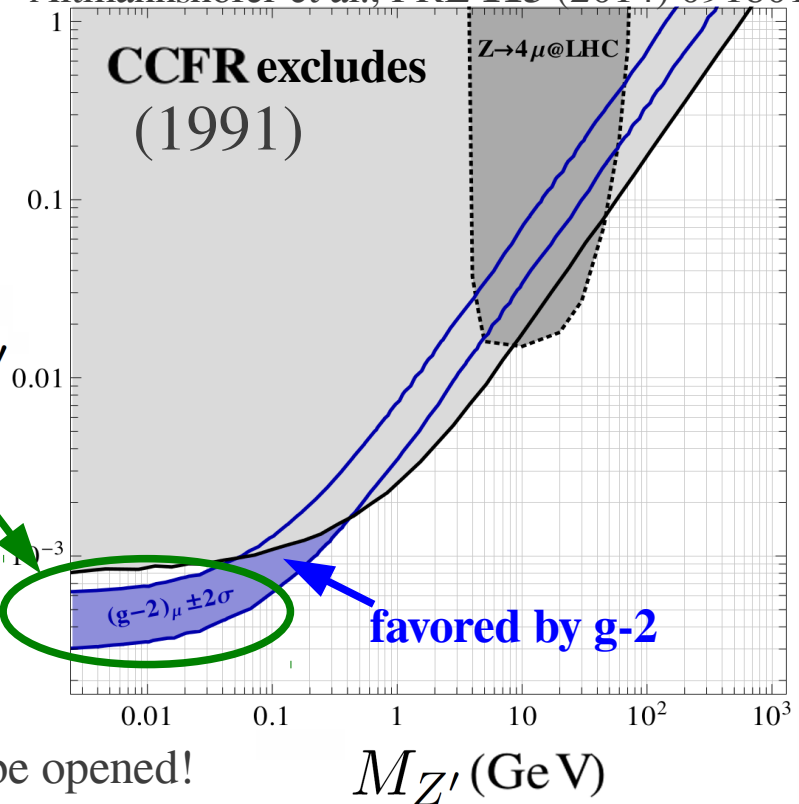
Consistent → constrains $g_{Z'}$ and $M_{Z'}$

$g_{\mu} - 2$ favored - Trident excl.

$M_{Z'} \lesssim 100$ MeV, $g_{Z'} \simeq \text{several} \cdot 10^{-4}$.



Altmannshofer et al., PRL **113** (2014) 091801



*The trident process must be recorded on the hard disks of the near detectors in modern oscillation experiments. They should be opened!

● **Neutrino trident process**

in neutrino-nucleon scattering events

- Available **data** reported by CCFR in 1991!

37 events (± 12.4)

CCFR collaboration, PRL **66** (1991) 3117
excavated recently (only cited 18 times)*

- Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

→ **Consistent** → constrains $g_{Z'}$ and $M_{Z'}$

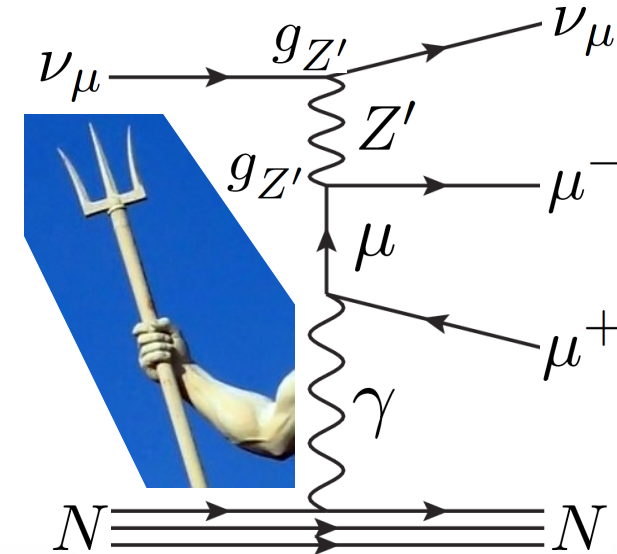
$g_{\mu} - 2$ favored - Trident excl.

$M_{Z'} \lesssim 100 \text{ MeV}, \quad g_{Z'} \simeq \text{several} \cdot 10^{-4}.$

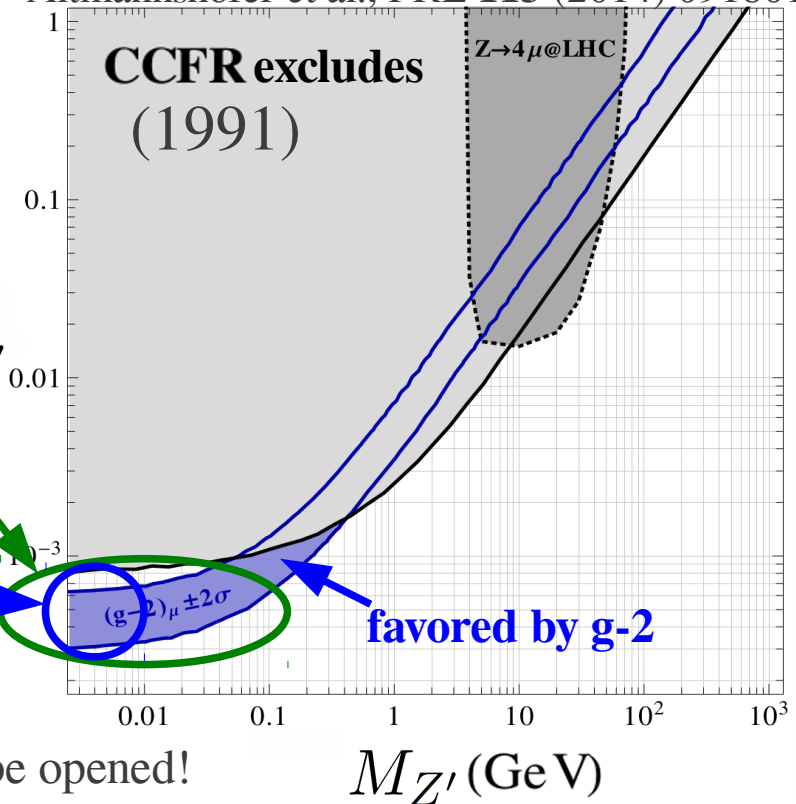
↕ **Coincide!**

IceCube Gap requires

$M_{Z'} \sim \text{MeV}, \quad g_{Z'} \gtrsim 10^{-4}.$



Altmannshofer et al., PRL **113** (2014) 091801



*The trident process must be recorded on the hard disks of the near detectors in modern oscillation experiments. They should be opened!

● Neutrino trident process

in neutrino-nucleon scattering events

- Available **data** reported by CCFR in 1991!

37 events (± 12.4)

CCFR collaboration, PRL **66** (1991) 3117
excavated recently (only cited 18 times)*

- Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

→ **Consistent** → constrains $g_{Z'}$ and $M_{Z'}$

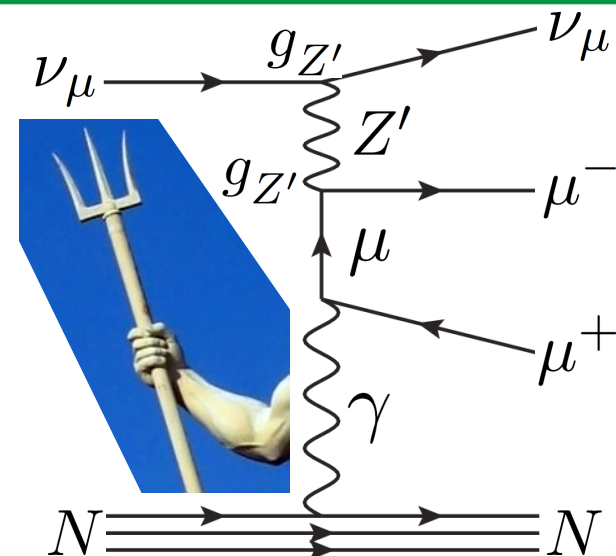
$g_\mu - 2$ favored - Trident excl.

$M_{Z'} \lesssim 100 \text{ MeV}$, $g_{Z'} \simeq \text{several} \cdot 10^{-4}$.

↕ **Coincide!**

IceCube Gap requires

$M_{Z'} \sim \text{MeV}$, $g_{Z'} \gtrsim 10^{-4}$.



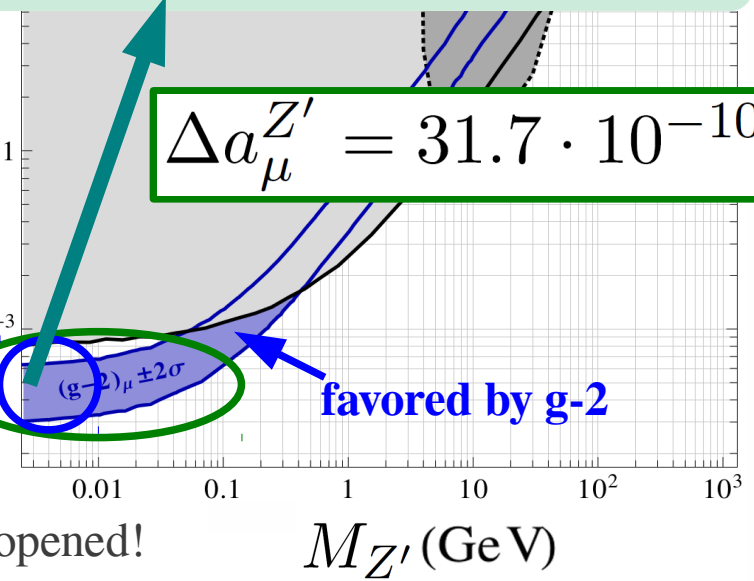
Altmannshofer et al., PRL **113** (2014) 091801

Reference values

$M_{Z'} = 2.75 \text{ MeV}$ and $g_{Z'} = 5.0 \cdot 10^{-4}$

$g_{Z',0.01}$

$\Delta a_\mu^{Z'} = 31.7 \cdot 10^{-10}$



favored by g-2

*The trident process must be recorded on the hard disks of the near detectors in modern oscillation experiments. They should be opened!

1 IceCube gap

- Attenuation of cosmic neutrino by secret neutrino interaction
- Gauged leptonic force $L_\mu - L_\tau$ as secret interaction

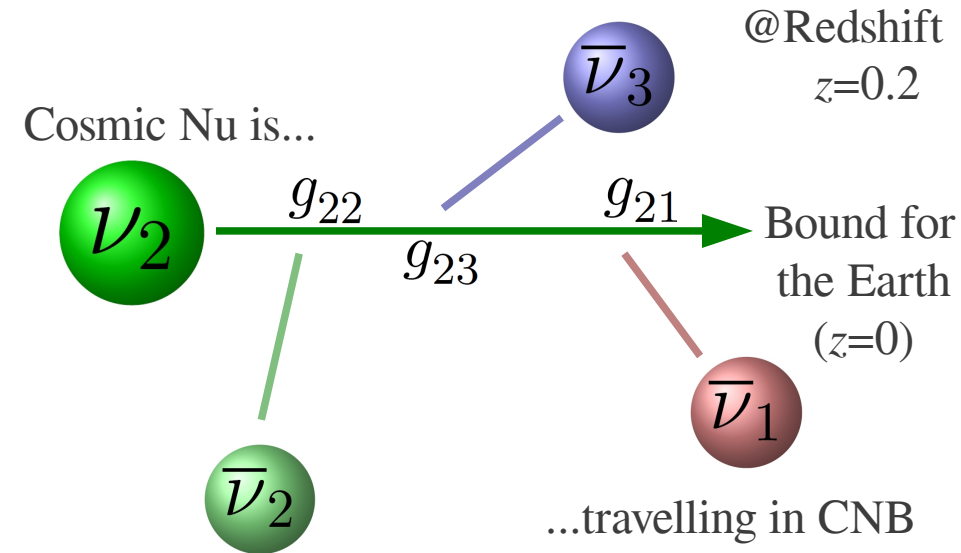
2 Muon anomalous magnetic moment

- Gauged leptonic force as a contribution to $g-2$
- Constraints from colliders and neutrino trident process

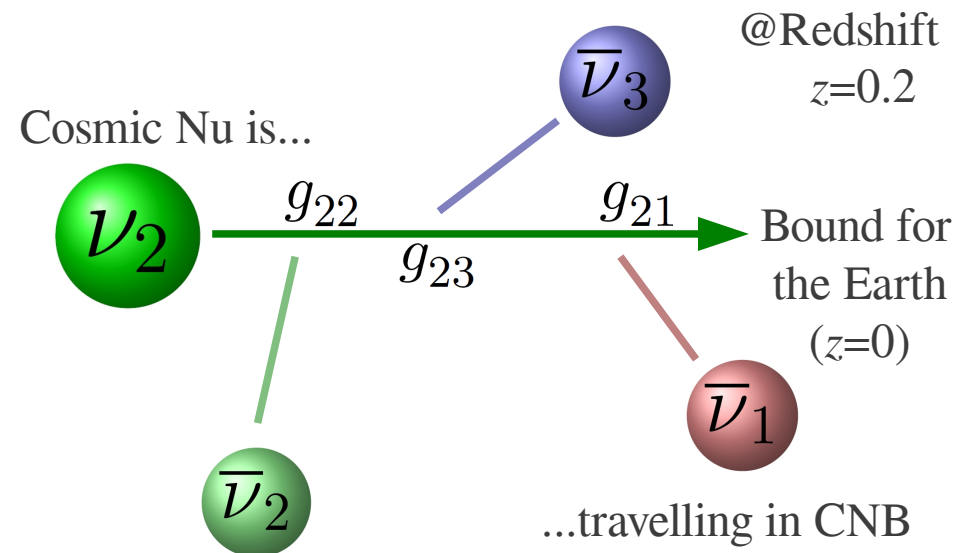
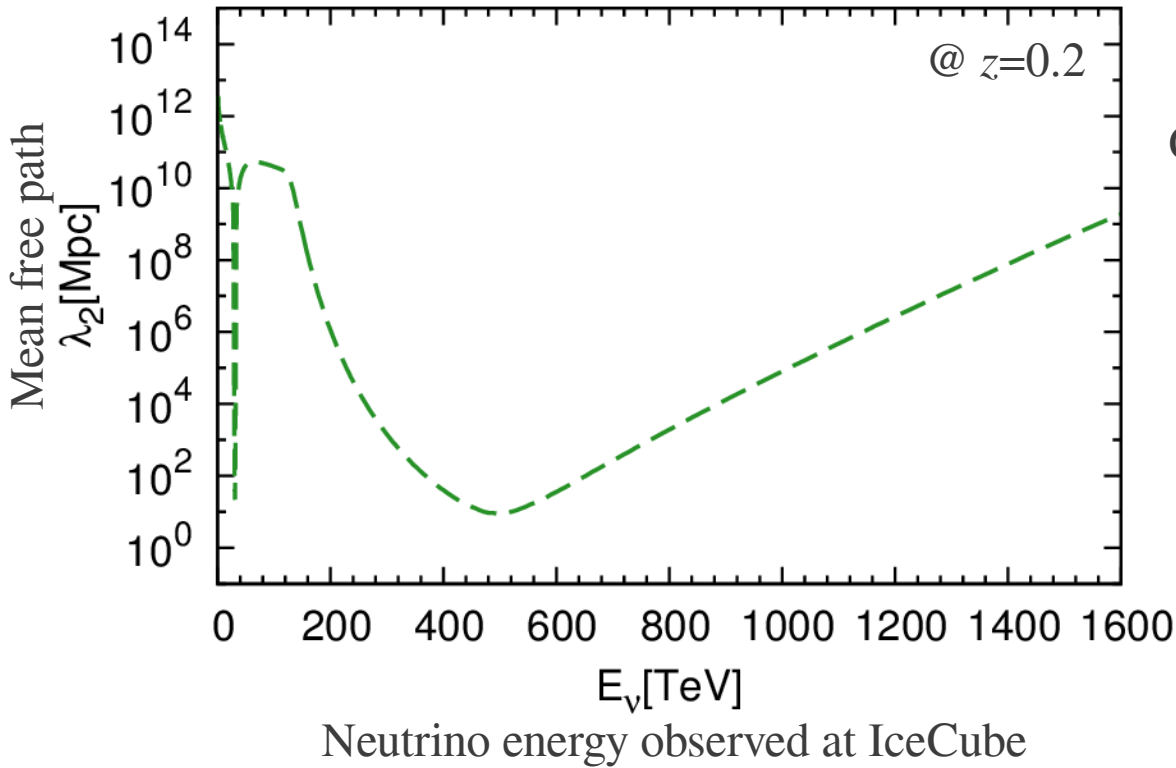
3 A solution to the gaps

- Reproduction of IceCube gap → distance to the neutrino source
→ neutrino mass spectrum

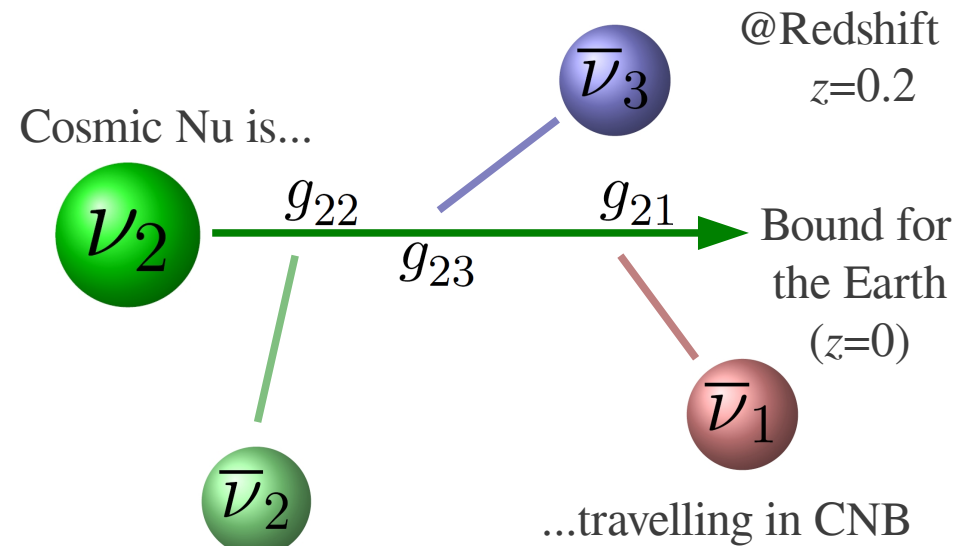
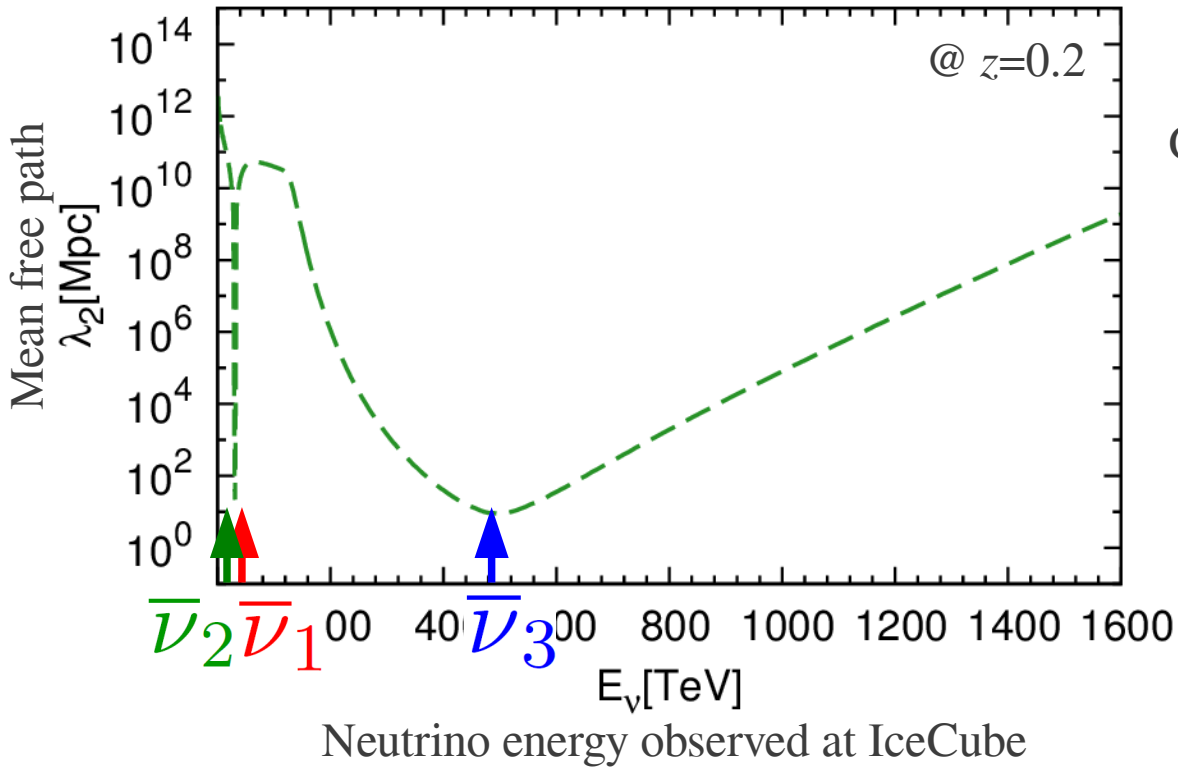
- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4} \rightarrow$ Favored by **$g-2$** and allowed by **Trident**
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us calculate the mean free path (for 2nd neutrino) at $z=0.2$.



- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4} \rightarrow$ Favored by **$g-2$** and allowed by **Trident**
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us calculate the mean free path (for 2nd neutrino) at $z=0.2$.



- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4} \rightarrow$ Favored by **$g-2$** and allowed by **Trident**
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us calculate the mean free path (for 2nd neutrino) at $z=0.2$.

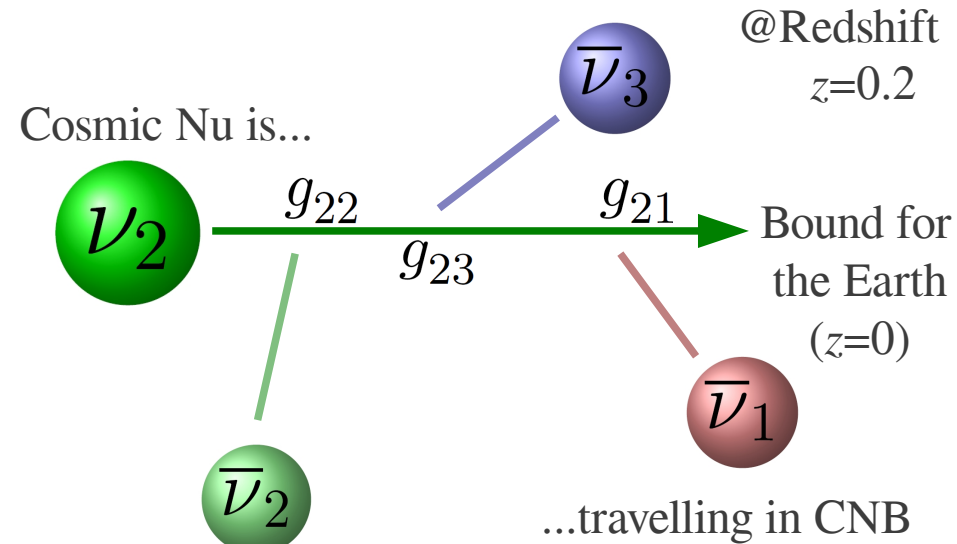
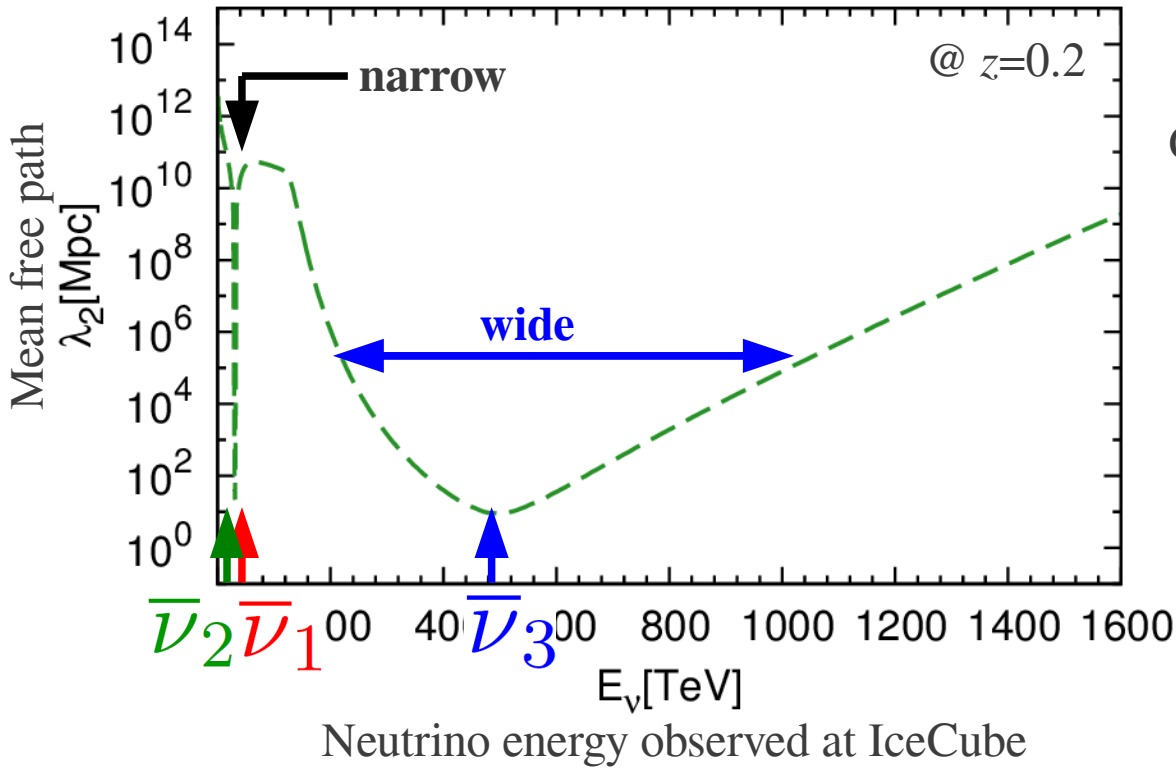


● Resonant condition w. CNB distribution

$$s \simeq \underbrace{2E_{\nu_{i=2}}(1+z)}_{\text{Neutrino energy @ } z} \left[\sqrt{|\mathbf{p}|^2 + m_{\nu_j}^2} - |\mathbf{p}| \cos \theta \right] \stackrel{!}{=} M_{Z'}^2$$

$|\mathbf{p}|$: CNB momentum follows Fermi-Dirac dist. $\lesssim (1+z)T_{\nu 0} \sim 2.0 \cdot 10^{-4} \text{ [eV]} @ z = 0.2$

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4} \rightarrow$ Favored by **$g-2$** and allowed by **Trident**
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us calculate the mean free path (for 2nd neutrino) at $z=0.2$.



$$m_{\nu_3} \sim |\mathbf{p}|$$

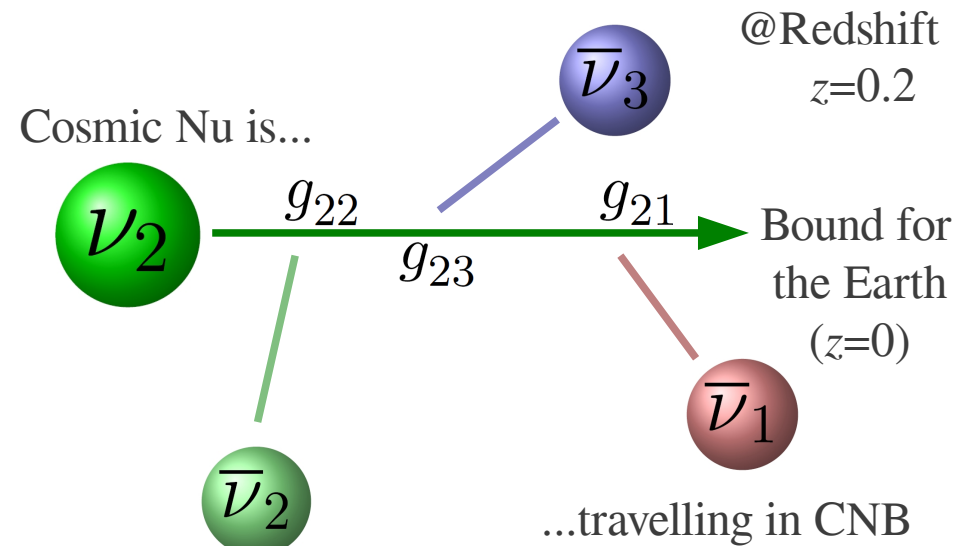
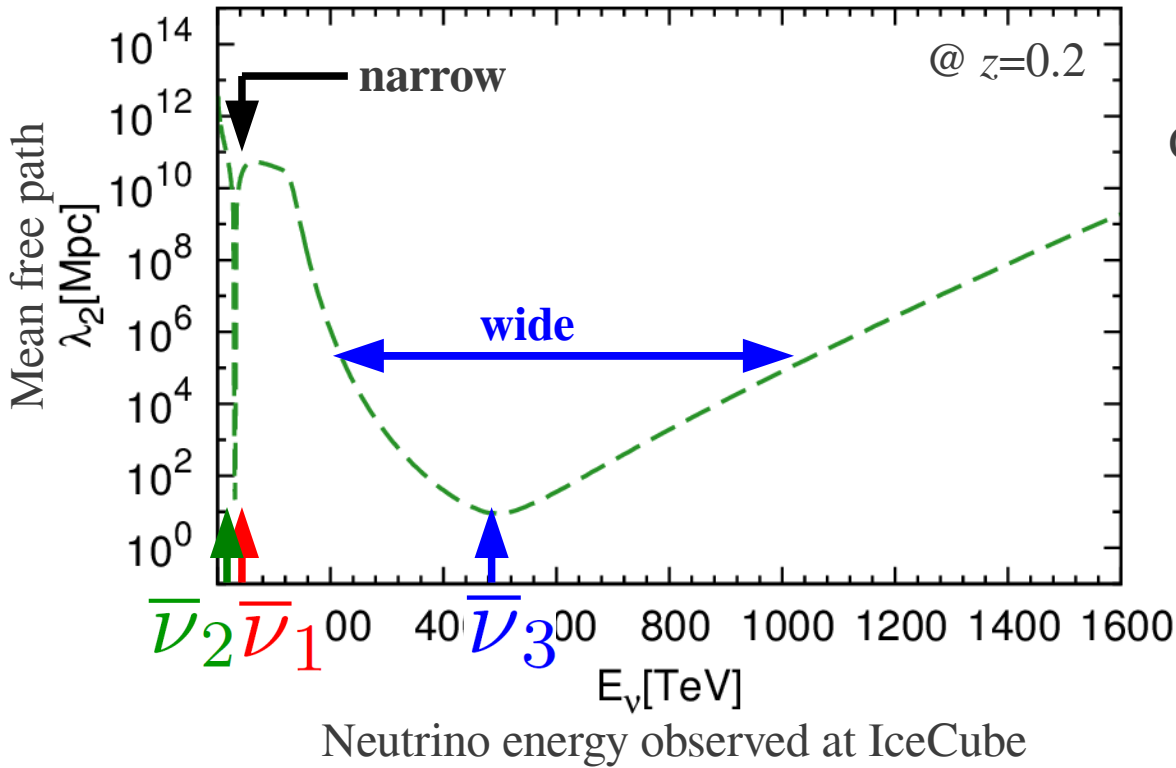
$$m_{\nu_{1,2}} \gg |\mathbf{p}|$$

- Resonant condition w. CNB distribution

$$s \simeq \underbrace{2E_{\nu_{i=2}}(1+z)}_{\text{Neutrino energy @ } z} \left[\sqrt{|\mathbf{p}|^2 + m_{\nu_i}^2} - |\mathbf{p}| \cos \theta \right] \stackrel{!}{=} M_{Z'}^2$$

$|\mathbf{p}|$: CNB momentum follows Fermi-Dirac dist. $\lesssim (1+z)T_{\nu 0} \sim 2.0 \cdot 10^{-4} \text{ [eV]} @ z = 0.2$

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4} \rightarrow$ Favored by **$g-2$** and allowed by **Trident**
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us calculate the mean free path (for 2nd neutrino) at $z=0.2$.



$$m_{\nu_3} \sim |\mathbf{p}|$$

$$m_{\nu_{1,2}} \gg |\mathbf{p}|$$

- Resonant condition w. CNB distribution

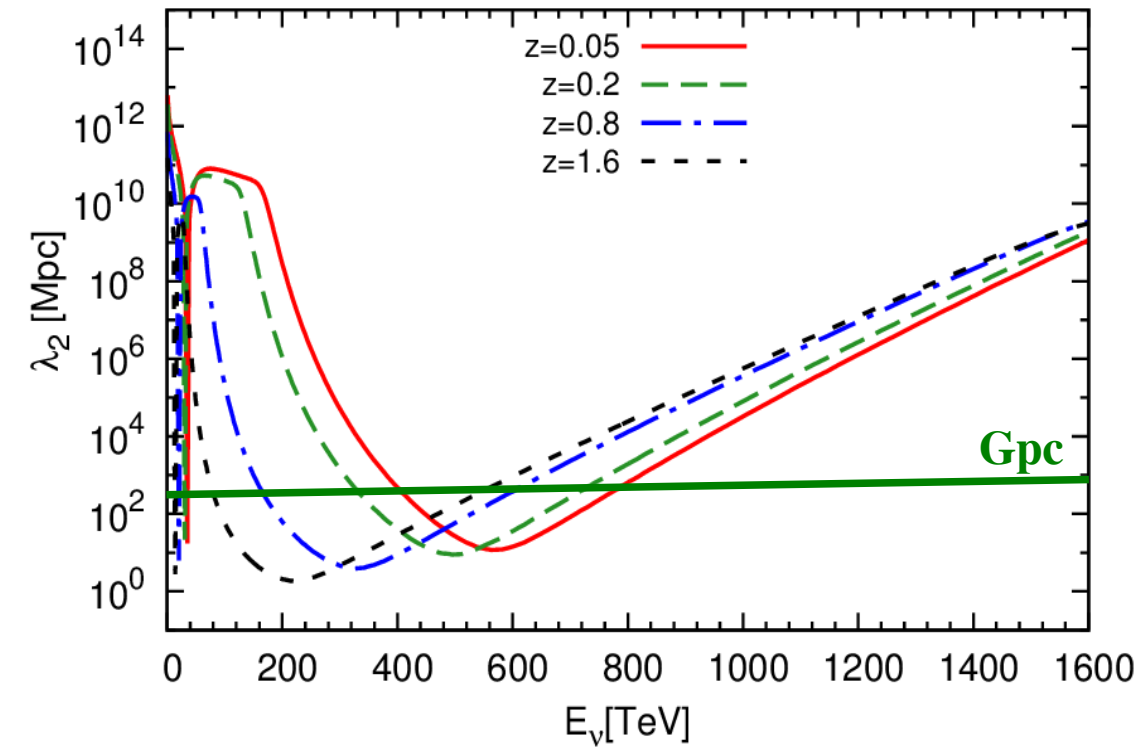
$$s \simeq \underbrace{2E_{\nu_{i=2}}(1+z)}_{\text{Neutrino energy @ } z} \left[\sqrt{|\mathbf{p}|^2 + m_{\nu_i}^2} - |\mathbf{p}| \cos \theta \right] \stackrel{!}{=} M_{Z'}^2$$

z shifts resonant E
 Small $m_{\text{Nu}} \rightarrow$ wide width
 Large $z \rightarrow$ wide width

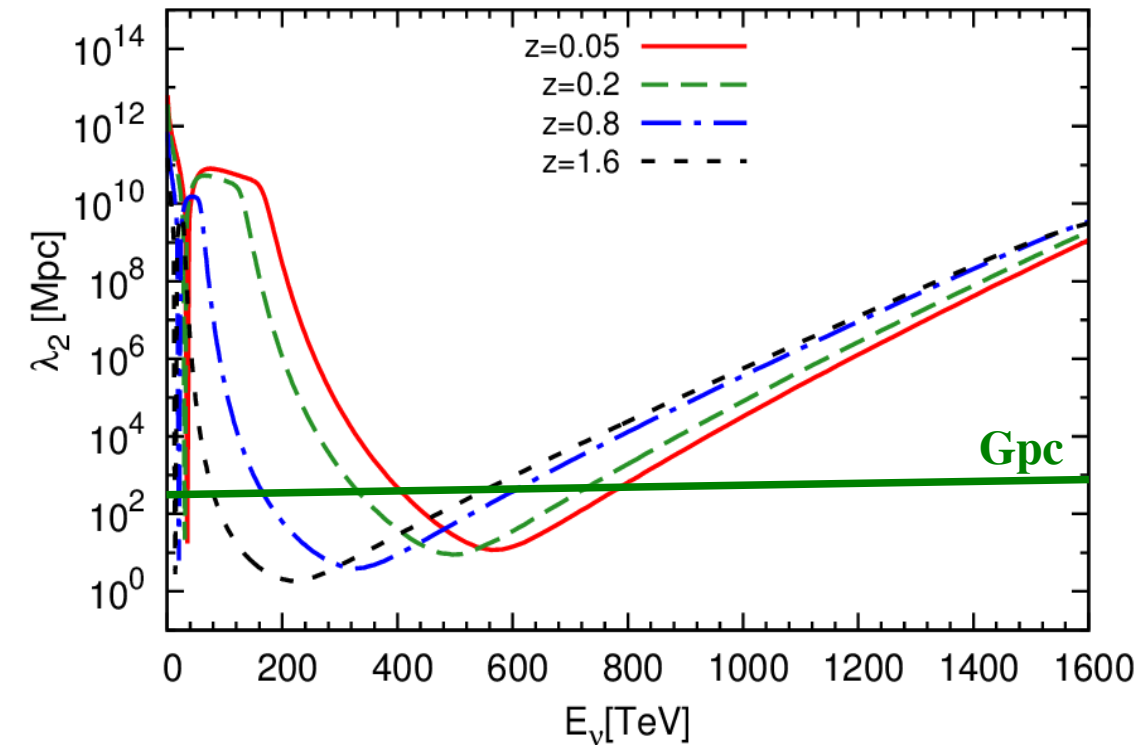
$|\mathbf{p}|$: CNB momentum follows Fermi-Dirac dist. $\lesssim (1+z) T_{\nu 0} \sim 2.0 \cdot 10^{-4} \text{ [eV]} @ z = 0.2$

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4}$
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us have a closer look at **z dependence** of MFP

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4}$
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us have a closer look at **z dependence** of MFP



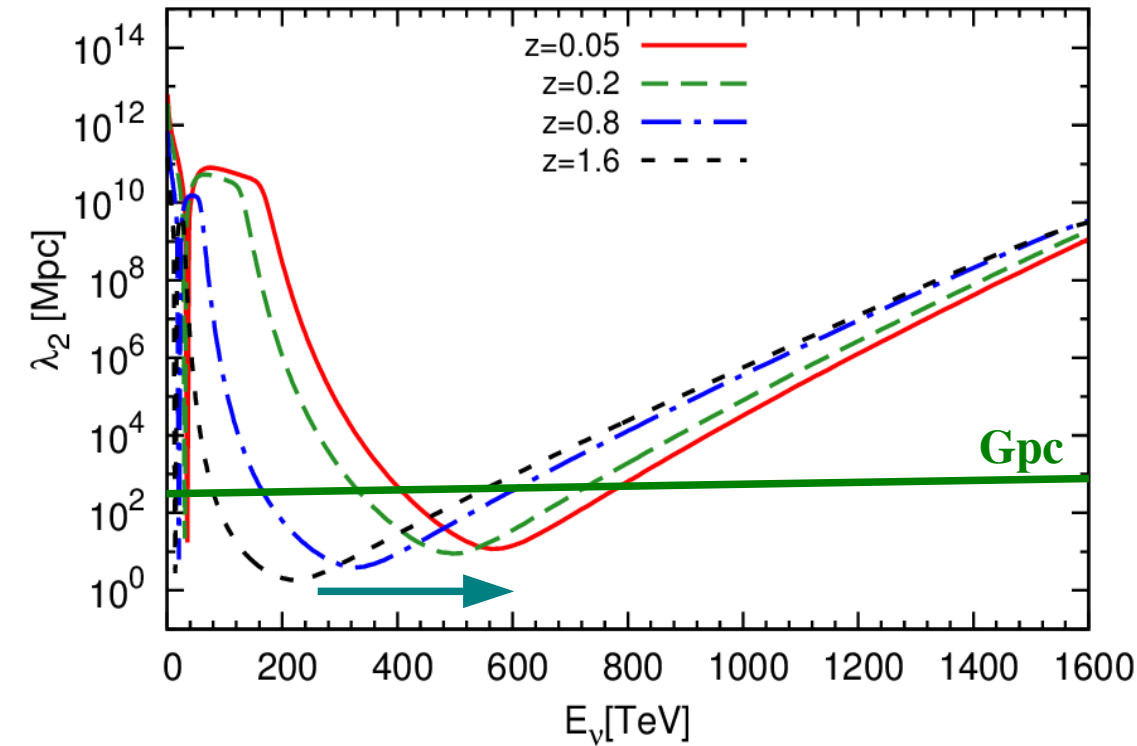
- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4}$
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us have a closer look at **z dependence** of MFP



- Cosmic neutrinos travel from z_{source} to $z = 0$ (Earth)
- The resonance energy shifts along the travel path.

To keep the width of the gap appropriate, the source should not be so distant from the Earth.

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4}$
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us have a closer look at **z dependence** of MFP



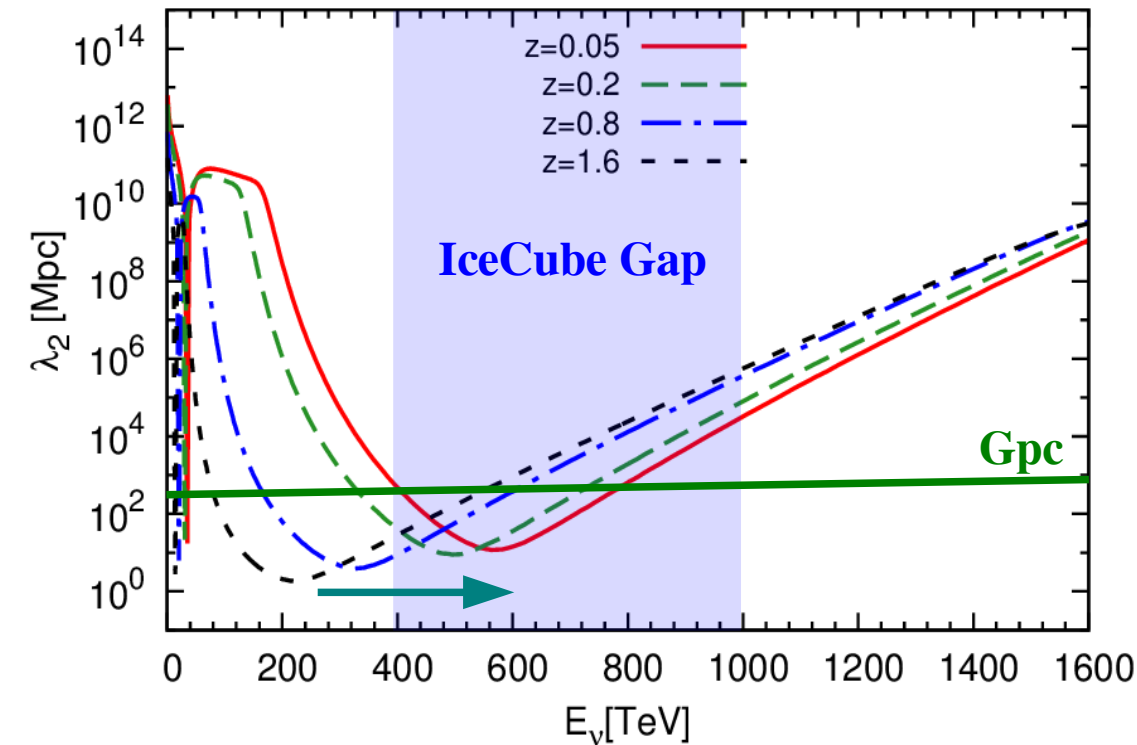
Peak position moves

$z_{\text{source}} \longrightarrow z = 0$

- Cosmic neutrinos travel from z_{source} to $z = 0$ (Earth)
- The resonance energy shifts along the travel path.

To keep the width of the gap appropriate, the source should not be so distant from the Earth.

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4}$
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV}$ and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us have a closer look at **z dependence** of MFP



Peak position moves

$z_{\text{source}} \longrightarrow z = 0$

- Cosmic neutrinos travel from z_{source} to $z = 0$ (Earth)
- The resonance energy shifts along the travel path.

To keep the width of the gap appropriate, the source should not be so distant from the Earth.

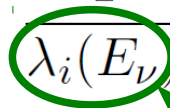
- We set $z_{\text{source}}=0.2$ so that the IceCube Gap is reproduced.

In reality, sources of cosmic neutrinos are distributed following some distribution function (e.g., the star formation rate)

● Mean free path → Spectrum

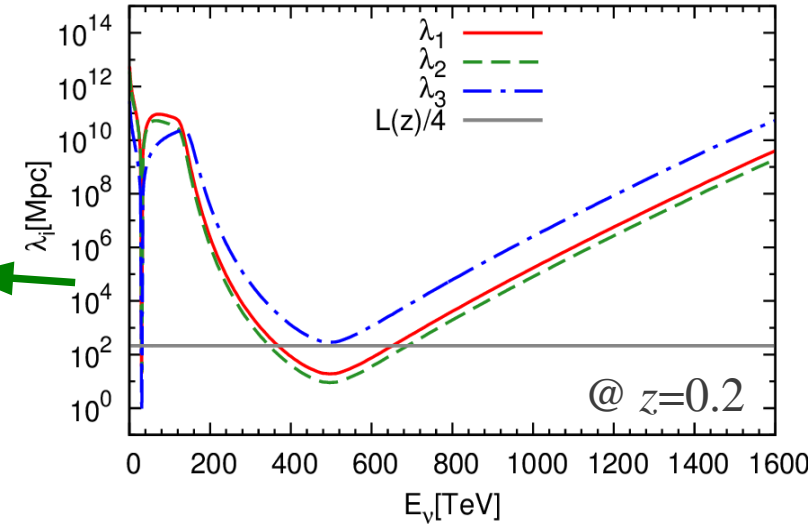
Following the approximation adopted in Ibe Kaneta PRD...

$$\varphi_i(E_\nu) = \varphi_i^{\text{original}}(E_\nu) \exp \left[- \int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_\nu)} \frac{dL}{dz} dz \right]$$



MFP

Same for 3 cosmic Nu's...



The resulting gap does not depend on the initial flavour composition.

● Mean free path → Spectrum

Following the approximation adopted in Ibe Kaneta PRD...

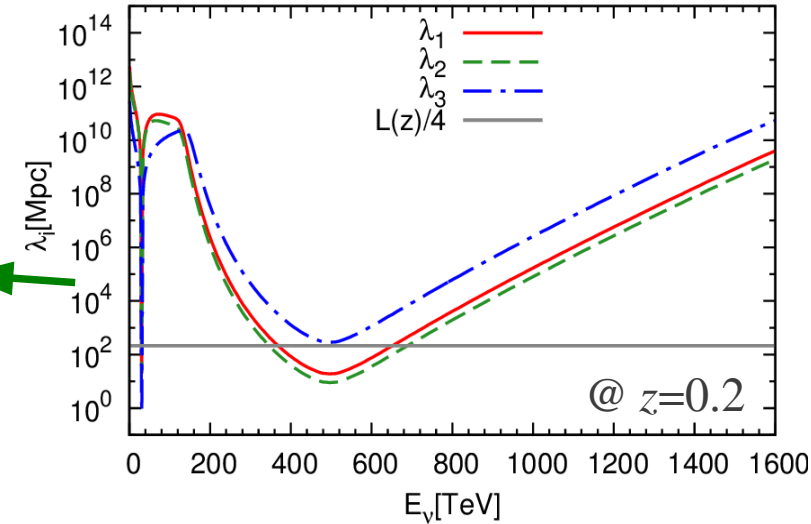
$$\varphi_i(E_\nu) = \varphi_i^{\text{original}}(E_\nu) \exp \left[- \int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_\nu)} \frac{dL}{dz} dz \right]$$

Resulting spectrum

Continuous (power-law) spectrum

MFP

Same for 3 cosmic Nu's...



The resulting gap does not depend on the initial flavour composition.

● Mean free path → Spectrum

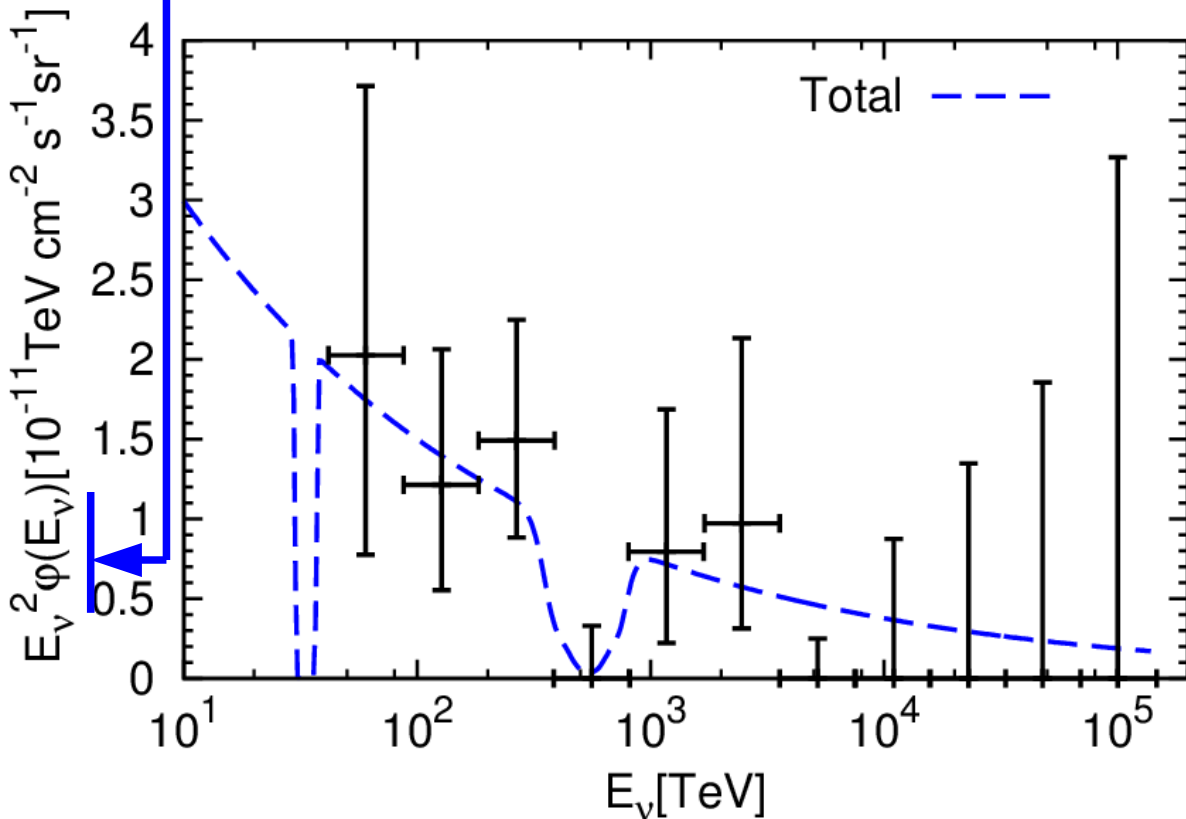
Following the approximation adopted in Ibe Kaneta PRD...

$$\varphi_i(E_\nu) = \varphi_i^{\text{original}}(E_\nu) \exp \left[- \int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_\nu)} \frac{dL}{dz} dz \right]$$

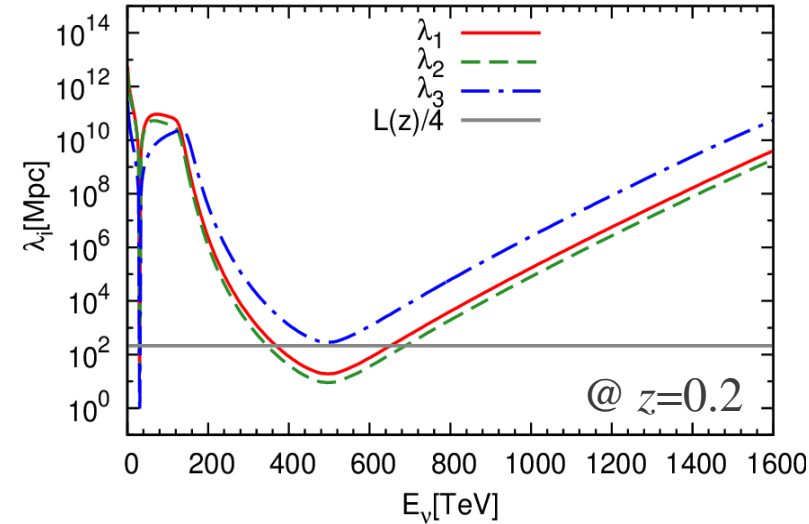
Resulting spectrum

$$\varphi(E_\nu) = \sum_i \varphi_i(E_\nu)$$

assuming flavour universal $\varphi_i^{\text{original}}(E_\nu)$



Same for 3 cosmic Nu's...



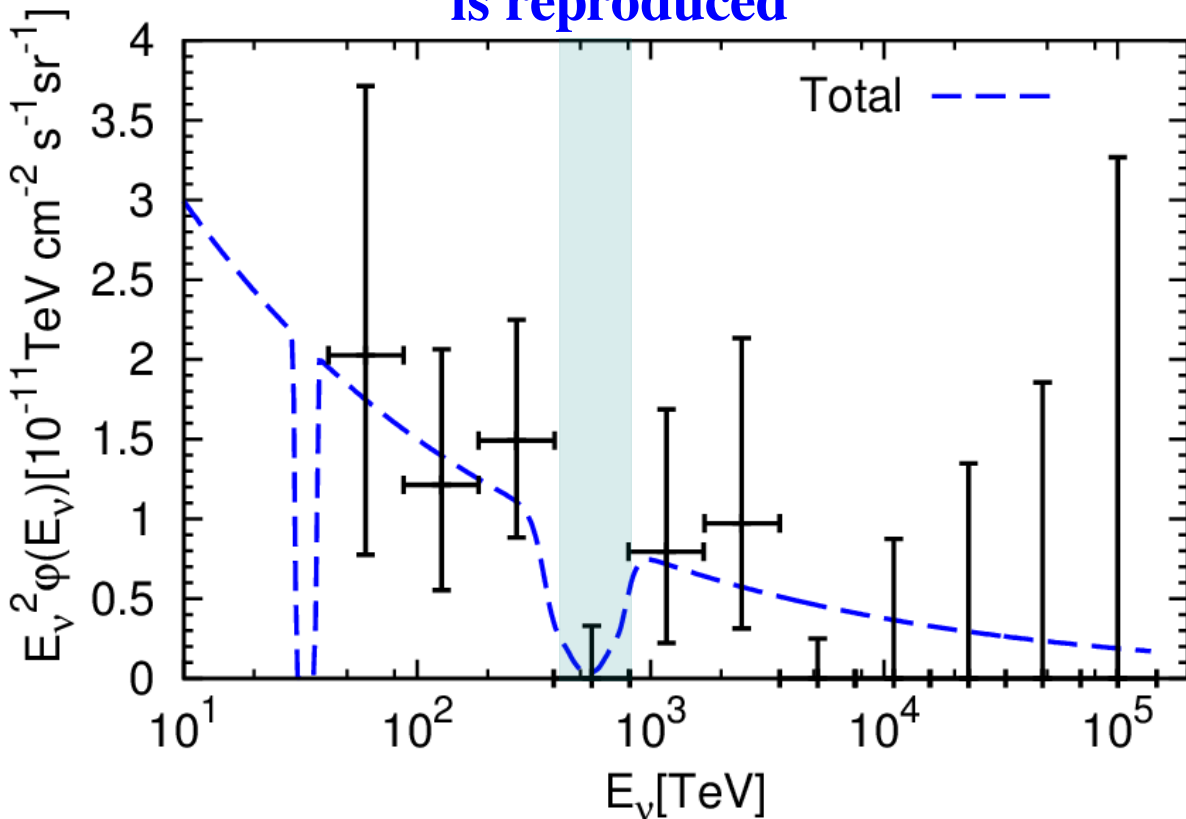
The resulting gap does not depend on the initial flavour composition.

● Mean free path → Spectrum

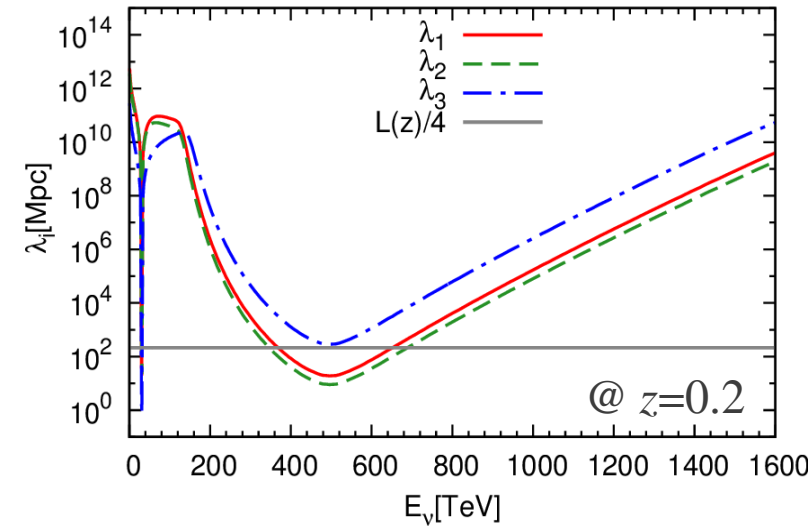
Following the approximation adopted in Ibe Kaneta PRD...

$$\varphi_i(E_\nu) = \varphi_i^{\text{original}}(E_\nu) \exp \left[- \int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_\nu)} \frac{dL}{dz} dz \right]$$

IceCube Gap
is reproduced



Same for 3 cosmic Nu's...



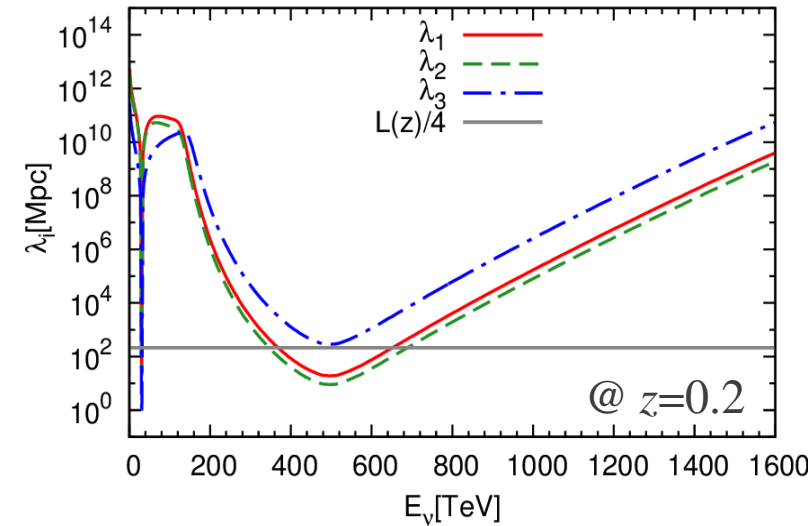
The resulting gap does not depend on the initial flavour composition.

Mean free path → Spectrum

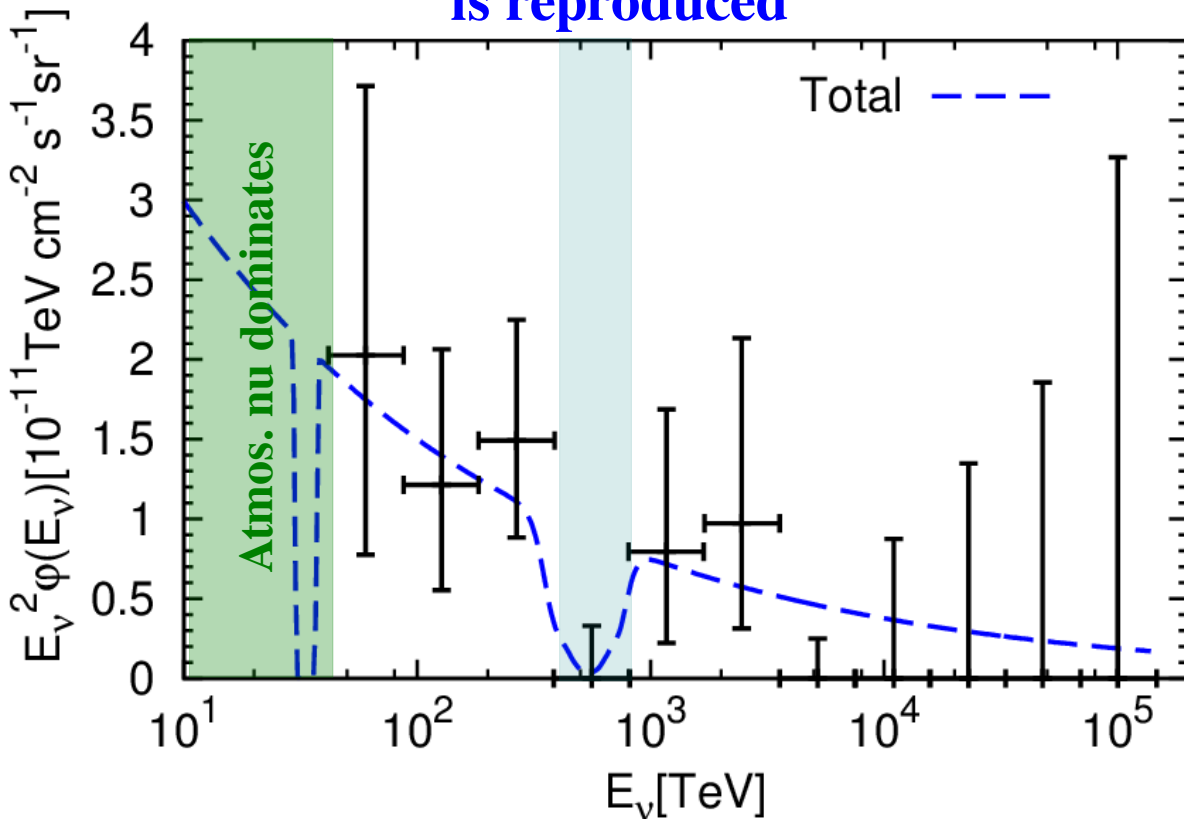
Following the approximation adopted in Ibe Kaneta PRD...

$$\varphi_i(E_\nu) = \varphi_i^{\text{original}}(E_\nu) \exp \left[- \int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_\nu)} \frac{dL}{dz} dz \right]$$

Same for 3 cosmic Nu's...



**IceCube Gap
is reproduced**



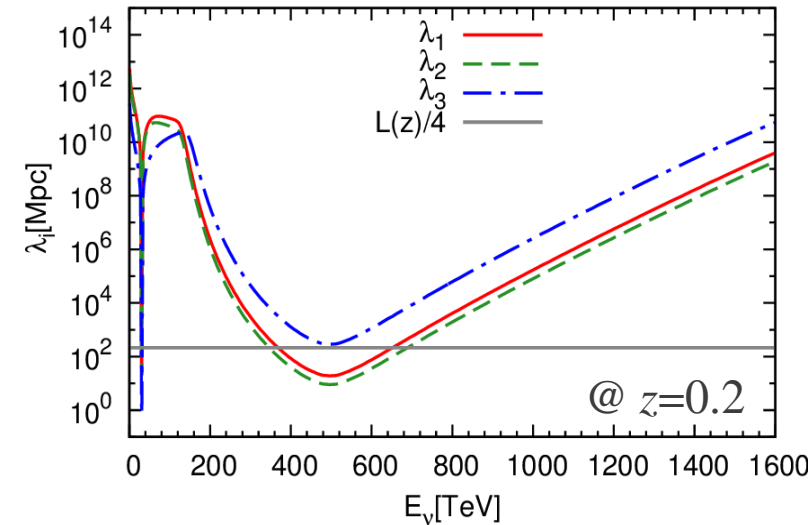
The resulting gap does not depend on the initial flavour composition.

Mean free path → Spectrum

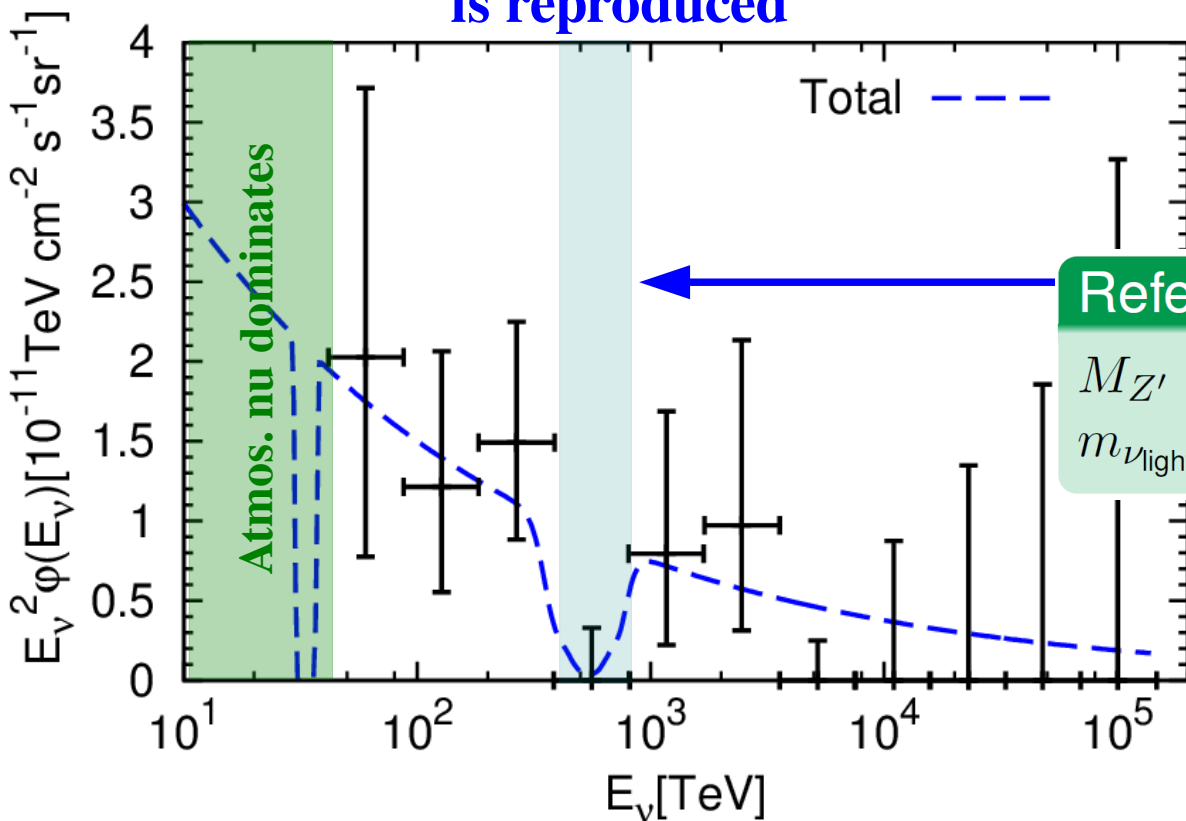
Following the approximation adopted in Ibe Kaneta PRD...

$$\varphi_i(E_\nu) = \varphi_i^{\text{original}}(E_\nu) \exp \left[- \int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_\nu)} \frac{dL}{dz} dz \right]$$

Same for 3 cosmic Nu's...



IceCube Gap
is reproduced



The resulting gap does not depend on the initial flavour composition.

Reference values

$$M_{Z'} = 2.75 \text{ MeV}, \quad g_{Z'} = 5.0 \cdot 10^{-4},$$

$$m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV (IH) and } z_{\text{source}} = 0.2.$$

Z' contribution to muon g-2

$$\Delta a_\mu^{Z'} = 31.7 \cdot 10^{-10}$$

g-2 Gap is filled

We dig the cosmic neutrino spectrum to make a gap and swing around the surplus soil to fill the gap in muon $g-2$.

Reference values

$$M_{Z'} = 2.75 \text{ MeV}, \quad g_{Z'} = 5.0 \cdot 10^{-4},$$

$$m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3} \text{ eV (IH) and } z_{\text{source}} = 0.2.$$

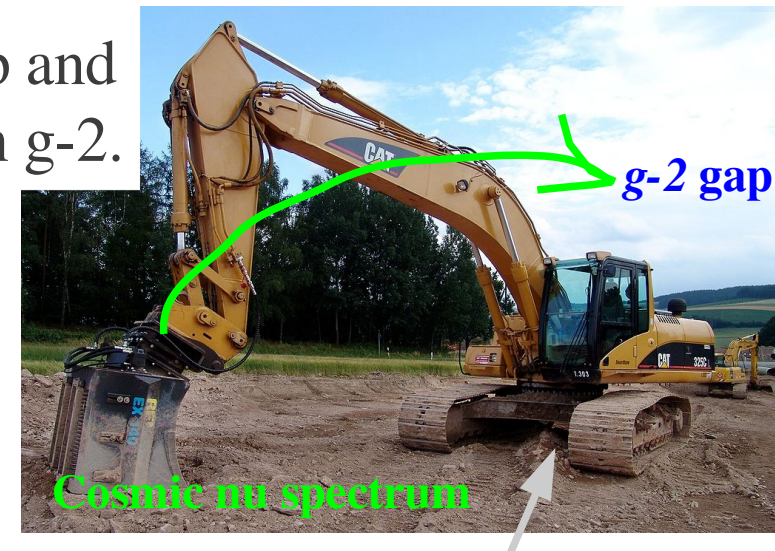
- IceCube Gap is reproduced.
- $\Delta a_{\mu}^{Z'} = 31.7 \cdot 10^{-10}$

But we did not...

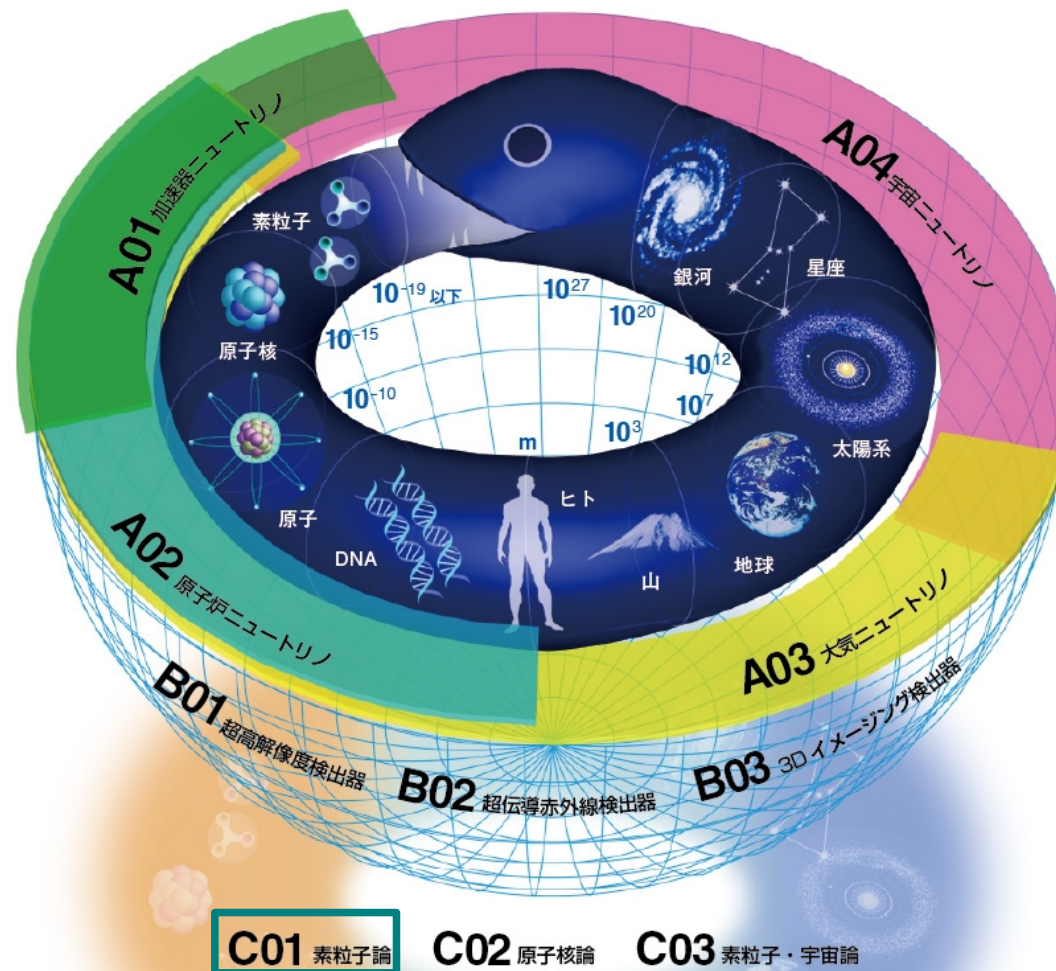
- ...take into account distribution of neutrino sources.
- ...also take into account secondary neutrino effect.
- ...discuss details of the model.

This small try shows that the idea works!

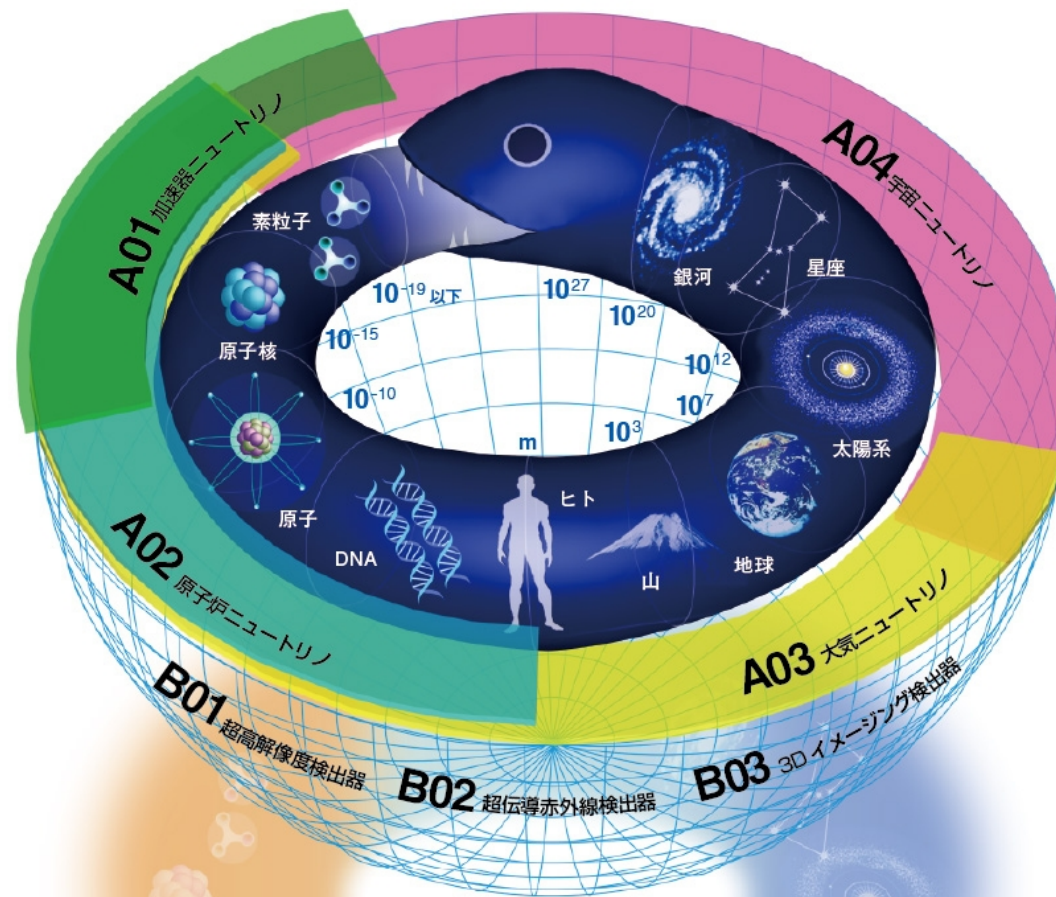
More precise, detailed, and sophisticated study may be worth to be done.



This tool is called as
“ $U(1)$ leptonic force L_{μ} - L_{τ} ”



I belong to this corner.



C01 素粒子論

C02 原子核論

C03 素粒子・宇宙論

I belong to this corner.

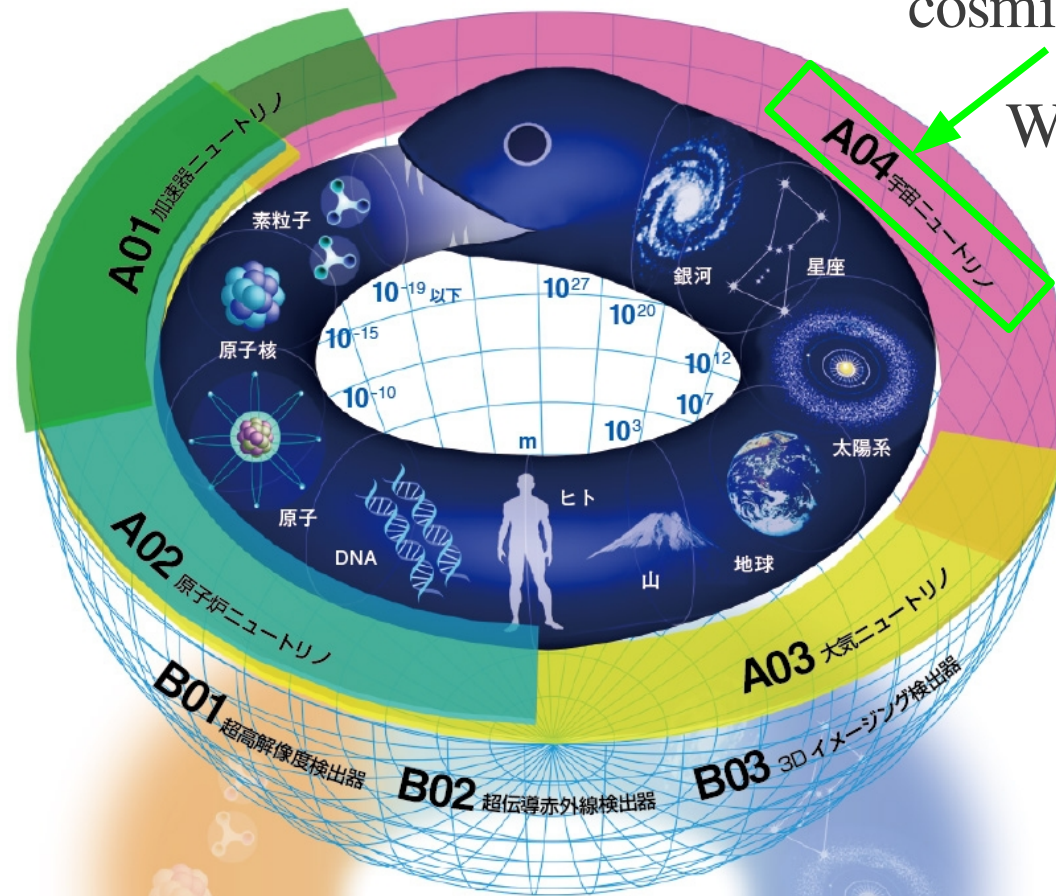
We share interest and ideas with...

Special thanks to Yoshida-san

We **mind the gap** on
cosmic neutrino spectrum

We are also motivated
from **muon g-2**

Charged Lepton



C01 素粒子論

C02 原子核論

C03 素粒子・宇宙論

I belong to this corner.

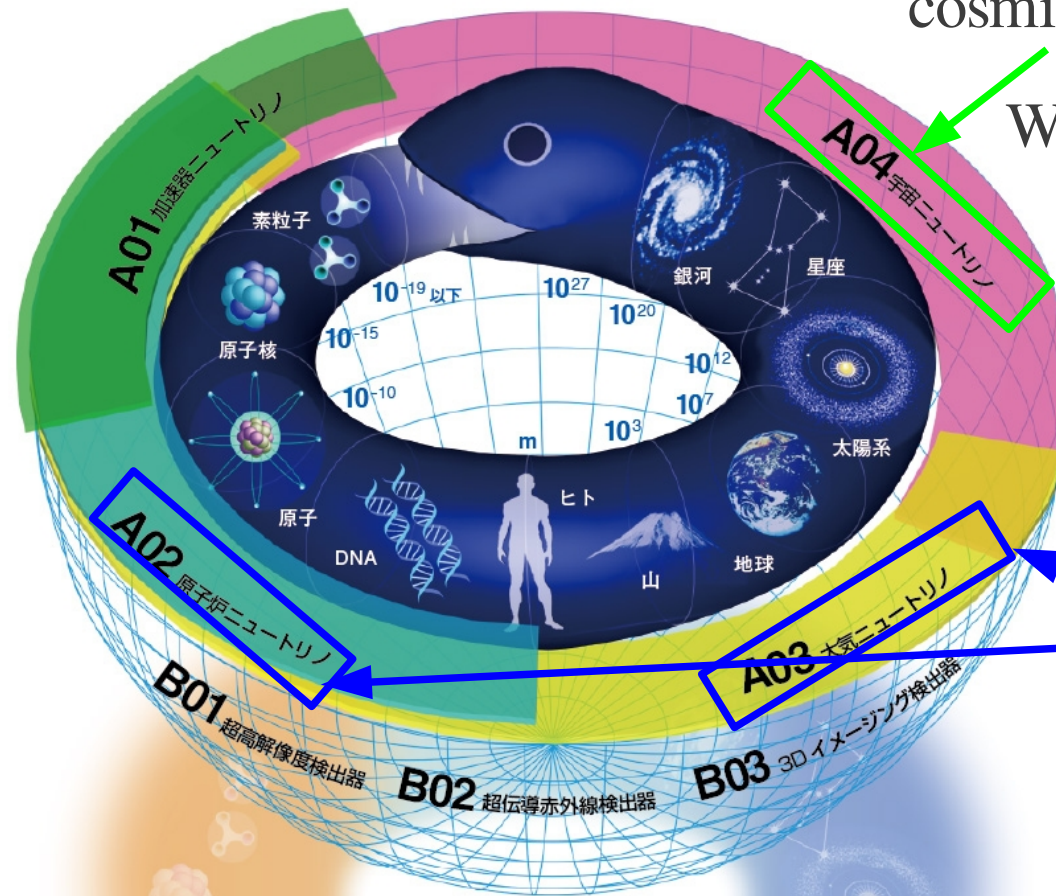
We share interest
and ideas with...

Special thanks to Yoshida-san

We **mind the gap** on
cosmic neutrino spectrum

We are also motivated
from **muon g-2**

Charged Lepton



The model is
inspired by...

C01 素粒子論

C02 原子核論

C03 素粒子・宇宙論

I belong to this corner.

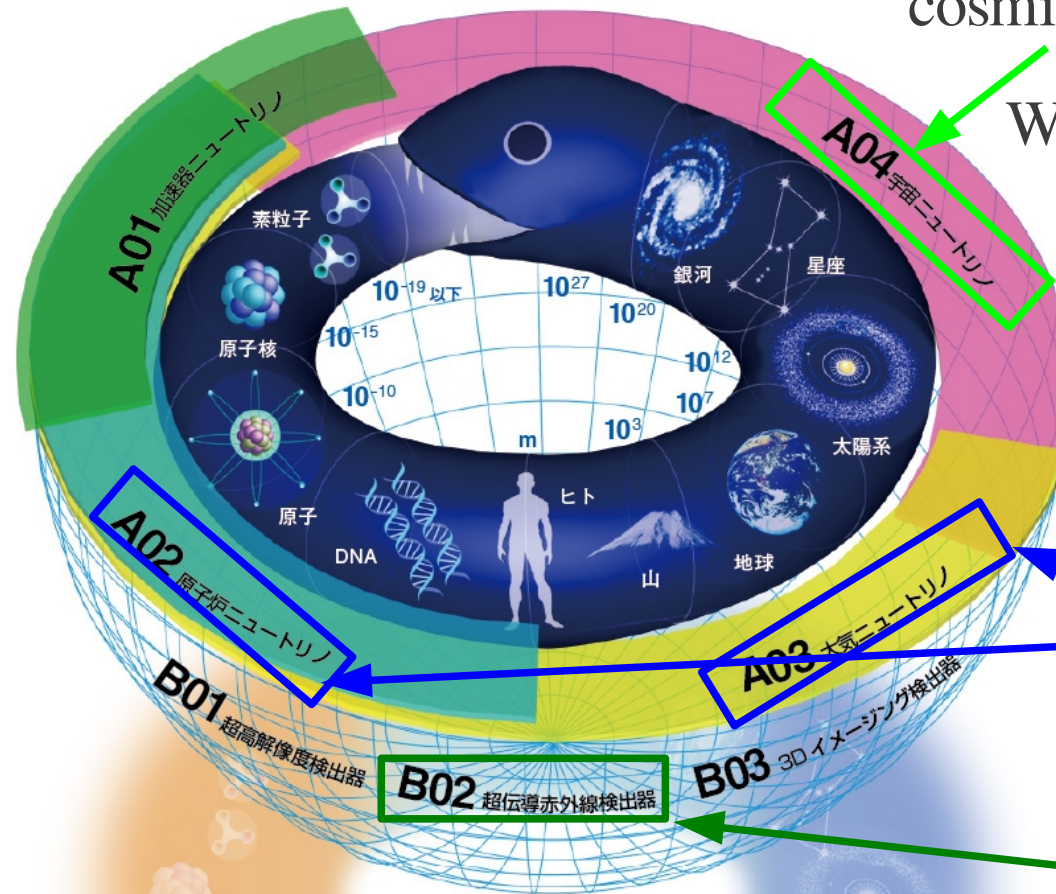
We share interest
and ideas with...

Special thanks to Yoshida-san

We **mind the gap** on
cosmic neutrino spectrum

We are also motivated
from **muon g-2**

Charged Lepton



The model is
inspired by...

CNB int.

C01 素粒子論

C02 原子核論

C03 素粒子・宇宙論

I belong to this corner.

We share interest
and ideas with...

Precision measurement of neutrino int.

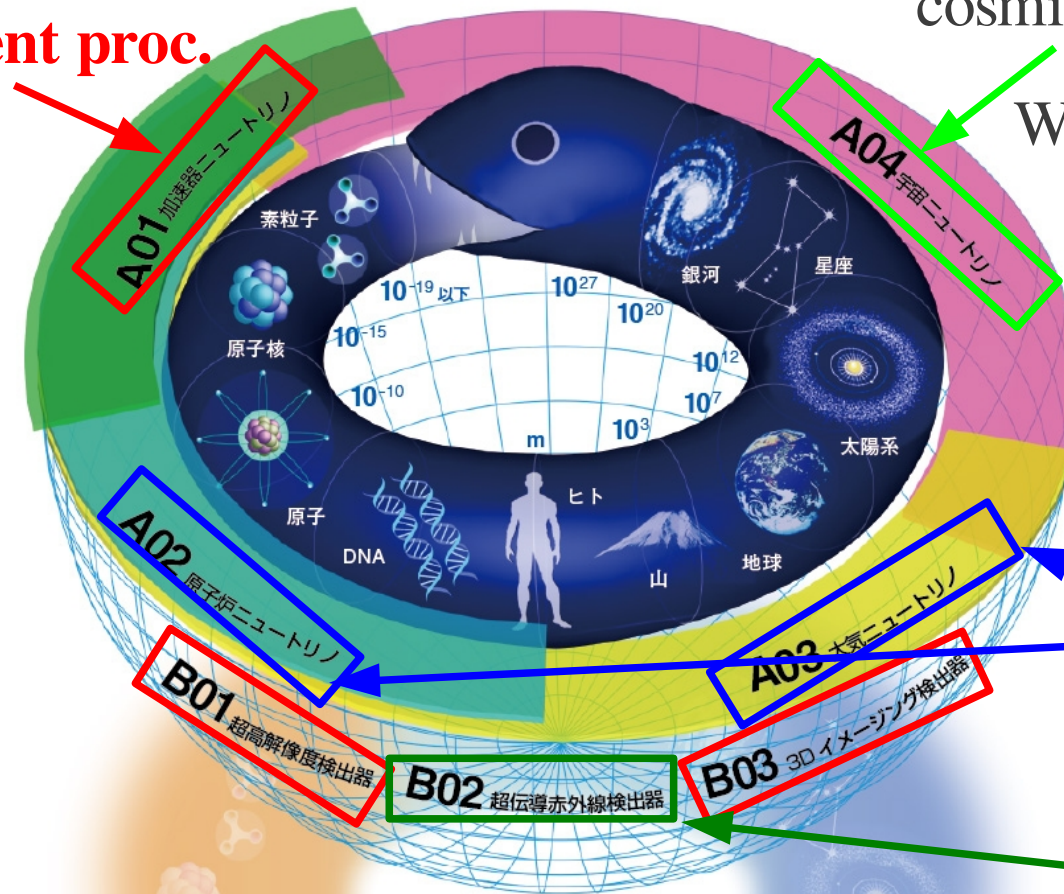
The most relevant bound is **neutrino trident proc.**

Special thanks to Yoshida-san

We **mind the gap** on cosmic neutrino spectrum

We are also motivated from **muon g-2**

Charged Lepton



The model is inspired by...

CNB int.

I belong to this corner.

We share interest and ideas with...

Precision measurement of neutrino int.

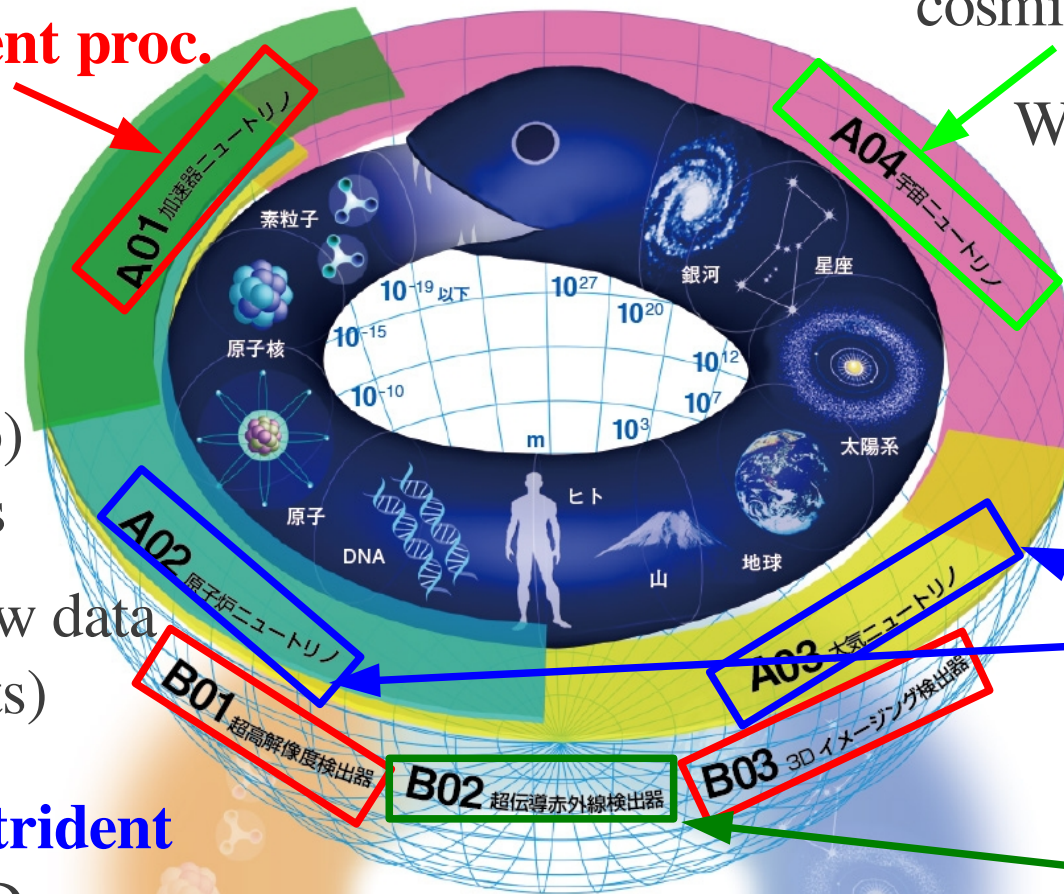
The most relevant bound is **neutrino trident proc.**

Special thanks to Yoshida-san

We **mind the gap** on cosmic neutrino spectrum

We are also motivated from **muon g-2**

Charged Lepton



Ask a favor (or two) to experimentalists

- Please provide new data (to feed theorists)

- Please check the **trident proc.** in your HDs.

The model is inspired by...

CNB int.

I belong to this corner.

C01 素粒子論 C02 原子核論 C03 素粒子・宇宙論

We share interest and ideas with...

Back up slides

