
News from Universal Extra Dimensions

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work with Gustavo Burdman and Eduardo Ponton

Bosons in compact spatial dimensions

4D flat spacetime \perp one dimension of size πR :



Boundary conditions : $\frac{\partial}{\partial y}\phi(x, 0) = \frac{\partial}{\partial y}\phi(x, \pi R) = 0$

KK decomposition : $\phi(x, y) = \frac{1}{\sqrt{\pi R}} \left[\phi^0(x) + \sqrt{2} \sum_{j \geq 1} \phi^j(x) \cos \left(\frac{jy}{R} \right) \right]$

Zero-mode: ϕ^0 - wave function is flat along the extra dimension.

Kaluza-Klein modes, $\phi^j(x)$:

particles with momentum in extra dimensions,

or 4D point of view: a tower of massive particles:

$$m_j^2 = m_0^2 + \frac{j^2}{R^2}$$

Gauge bosons in 5D:

$A_\mu(x^\nu, y)$, $\mu, \nu = 0, 1, 2, 3$, and

$A_y(x^\nu, y)$ – polarization along the extra dimension.

4D point of view: $A_y(x^\nu, y)$ is a tower of spin-0 KK modes.

Dirichlet B.C : $A_y(x^\nu, 0) = A_y(x^\nu, \pi R) = 0$

KK decomposition : $A_y(x^\nu, y) = \sqrt{\frac{2}{L}} \sum_{j \geq 1} A_y^j(x^\nu) \sin\left(\frac{jy}{R}\right)$

→ $A_y(x^\nu, y)$ does not have a 0-mode! (Odd field)

$A_y^j(x^\nu)$ becomes the longitudinal degree of freedom of
the spin-1 KK mode $A_\mu^j(x^\nu)$.

Fermions in a compact dimension

Lorentz group in 5D \Rightarrow vector-like fermions:

$$\chi = \chi_L + \chi_R$$

Chiral boundary conditions:

$$\chi_L(x^\mu, 0) = \chi_L(x^\mu, \pi R) = 0$$

$$\frac{\partial}{\partial y} \chi_R(x^\mu, 0) = \frac{\partial}{\partial y} \chi_R(x^\mu, \pi R) = 0$$

Kaluza-Klein decomposition:

$$\chi = \frac{1}{\sqrt{\pi R}} \left\{ \chi_R^0(x^\mu) + \sqrt{2} \sum_{j \geq 1} \left[\chi_R^j(x^\mu) \cos \left(\frac{\pi j y}{L} \right) + \chi_L^j(x^\mu) \sin \left(\frac{\pi j y}{L} \right) \right] \right\}$$

Universal Extra Dimensions

T. Appelquist, H.-C. Cheng, B. Dobrescu, Phys.Rev.D64 (2001)

All Standard Model particles propagate in $D \geq 5$ dimensions.

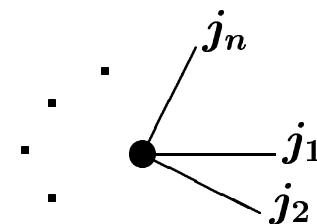
Kaluza-Klein modes are states of definite momentum along the compact dimensions.

Momentum conservation \rightarrow KK-number conservation

$$\mathcal{L}_{4D} = \int dy \mathcal{L}_{5D}$$

At each interaction vertex:

$$j_1 \pm j_2 \pm \dots \pm j_n = 0 \text{ for a certain choice of } \pm$$

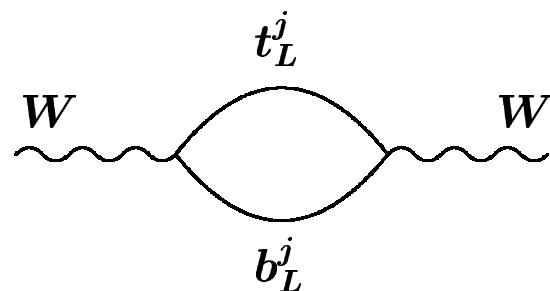


In particular: $0 \pm \dots \pm 0 \neq 1$

\Rightarrow tree-level exchange of KK modes does not contribute to currently measurable quantities

\Rightarrow no single KK 1-mode production at colliders

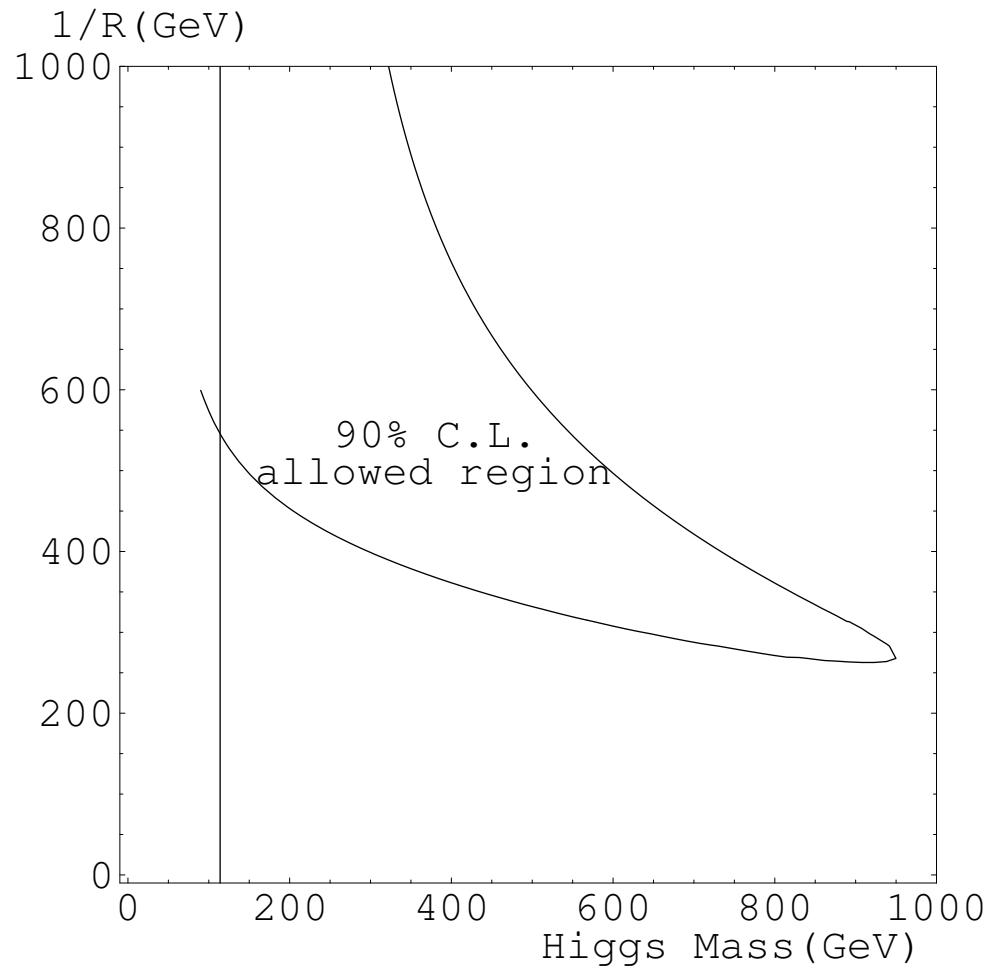
Bounds from one-loop shifts in W and Z masses, and other observables:



$$\frac{1}{R} \gtrsim 300 - 500 \text{ GeV}$$

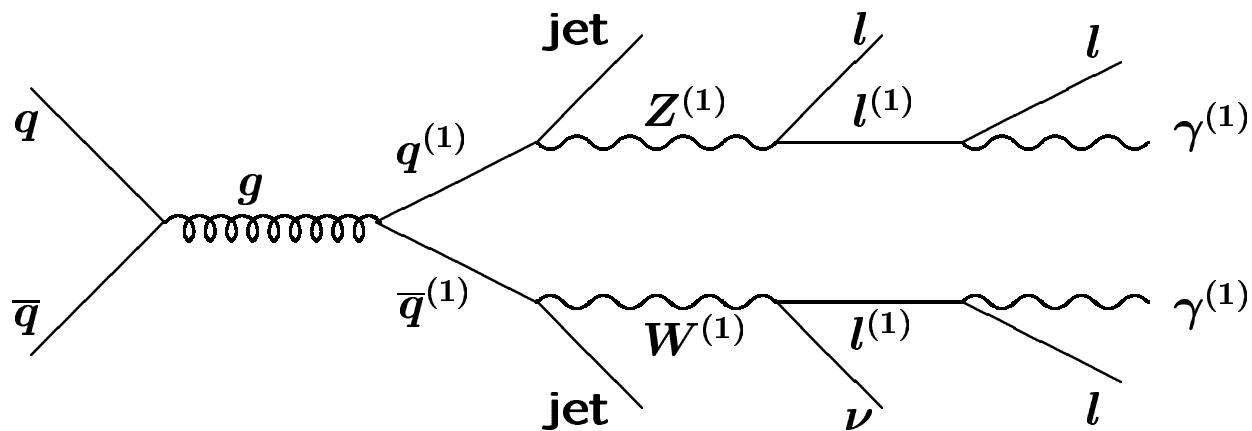
Contributions to the T parameter from top KK modes may compensate for the effect of a heavy Higgs boson on the electroweak fits.

*Appelquist, Yee,
hep-ph/0211023*



- Pair production of KK 1-modes at colliders:
cascade decays to $4l + \cancel{E}_T$ (soft leptons).
Could be discovered soon!
(Cheng, Matchev, Schmaltz, hep-ph/0205314)

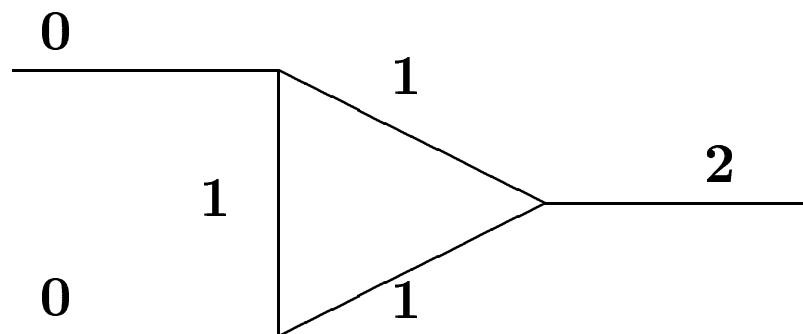
CDF analysis of $3l + \cancel{E}_T$ (soft leptons): $1/R > 280$ GeV (Run I)



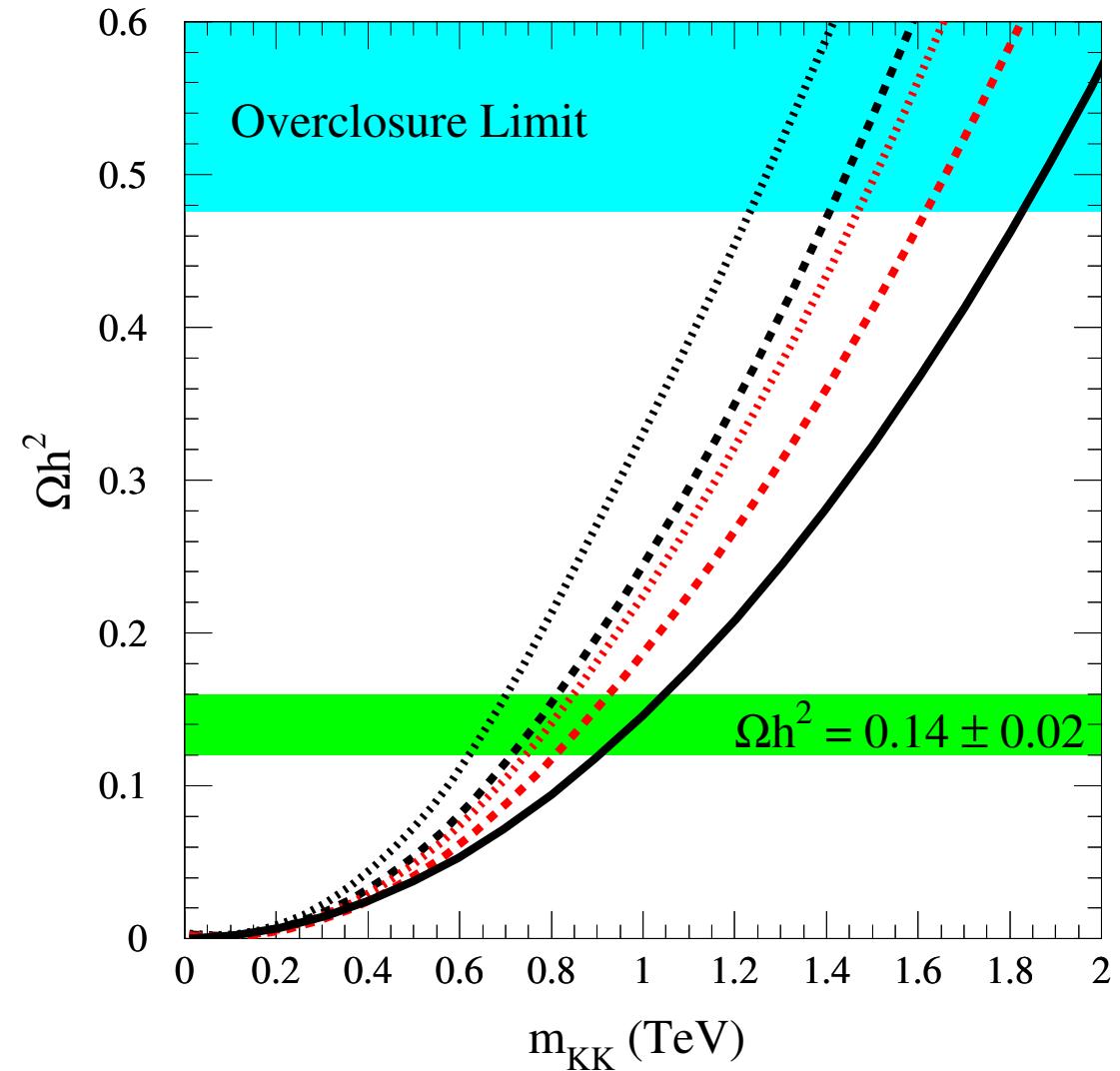
At one-loop level: $j_1 \pm j_2 \pm \dots \pm j_n = \text{even}$

KK parity is conserved: $(-1)^j \Rightarrow$ **Lightest KK Particle is a viable dark matter candidate.**

At colliders: ***s*-channel production of the 2-modes**

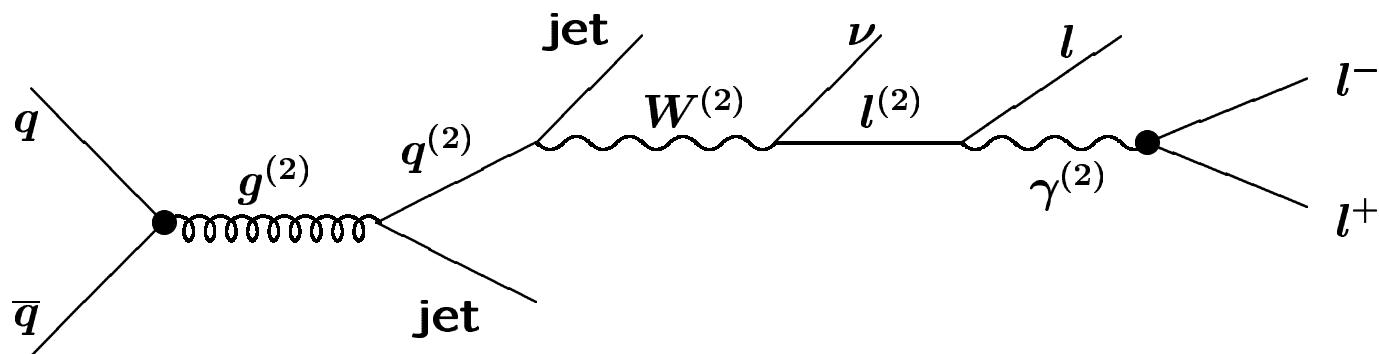


**Lightest KK particle
is stable in UED:
 $\gamma^{(1)}$ is a viable dark
matter candidate**
 (from Servant, Tait,
[hep-ph/0206071](#))



Second-level masses: $\sim 2/R$.

Cascade decay of the 2-mode is followed by $\gamma^{(2)}$ decay into hard leptons:



Particularly useful at the LHC (K. Kong, K. Matchev, hep-ph/0509246)

→ would allow discrimination of UED & MSSM.

6-Dimensional Standard Model

Motivation

- 3 generations of quarks and leptons are required for global $SU(2)_W$ anomaly cancellation
- proton is long-lived due to 6D Lorentz invariance
- neutrinos are interesting (Dirac masses)

Two possible chirality assignments for the 6D quarks and leptons:

$$Q_+, \, u_-, \, d_-, \, \left\{ \begin{array}{l} \mathcal{L}_+, \, \mathcal{E}_-, \, \mathcal{N}_- \\ \text{or} \\ \mathcal{L}_-, \, \mathcal{E}_+, \, \mathcal{N}_+ \end{array} \right.$$

Chiral boundary conditions on a square

(Dobrescu, Ponton, hep-ph/0401032; G. Burdman and E. Ponton, hep-ph/0506334)

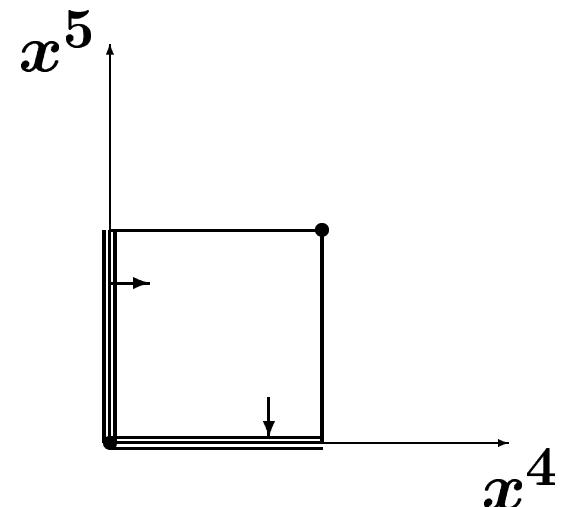
Identify pairs of adjacent sides:

$$\mathcal{L}(x^\mu, y, 0) = \mathcal{L}(x^\mu, 0, y)$$

$$\mathcal{L}(x^\mu, y, L) = \mathcal{L}(x^\mu, L, y)$$

$$\Phi(y, 0) = e^{in\pi/2} \Phi(0, y), \dots$$

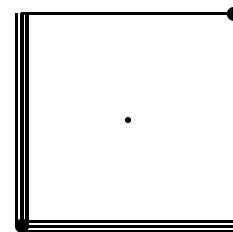
$$\partial_5 \Phi|_{(x^4, x^5) = (y, 0)} = -e^{in\pi/2} \partial_4 \Phi|_{(x^4, x^5) = (0, y)}$$



Spectrum of KK modes:

(j, k)	$(1, 0)$	$(1, 1)$	$(2, 0)$	$(2, 1)$ $(1, 2)$	$(2, 2)$	$(3, 0)$	$(3, 1)$ $(1, 3)$
$M_{j,k}R$	1	$\sqrt{2}$	2	$\sqrt{5}$	$2\sqrt{2}$	3	$\sqrt{10}$

Kaluza-Klein parity



Reflections about the center of the square $(L/2, L/2)$,

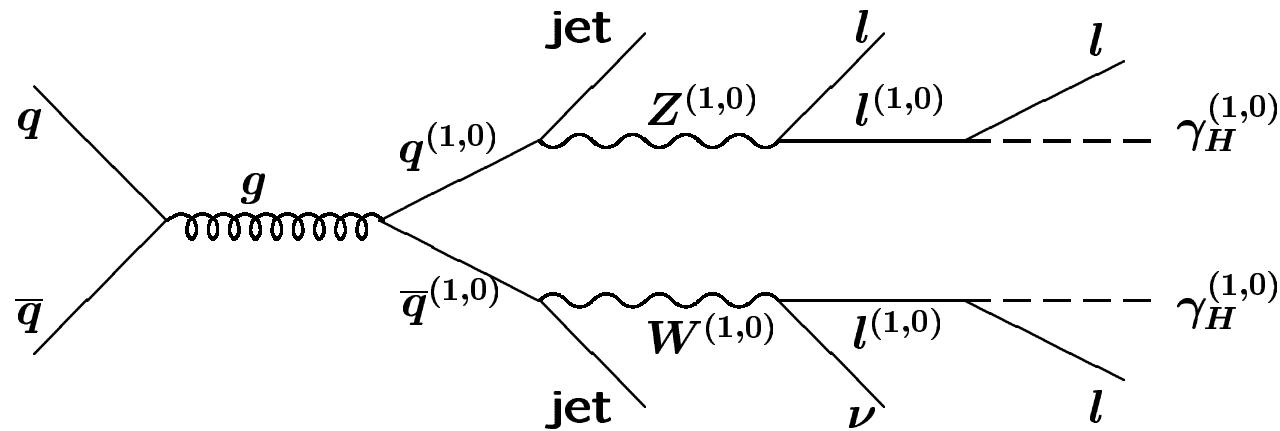
$$(x^4, x^5) \mapsto (L - x^4, L - x^5)$$

\Rightarrow invariance under Z_2 transformation

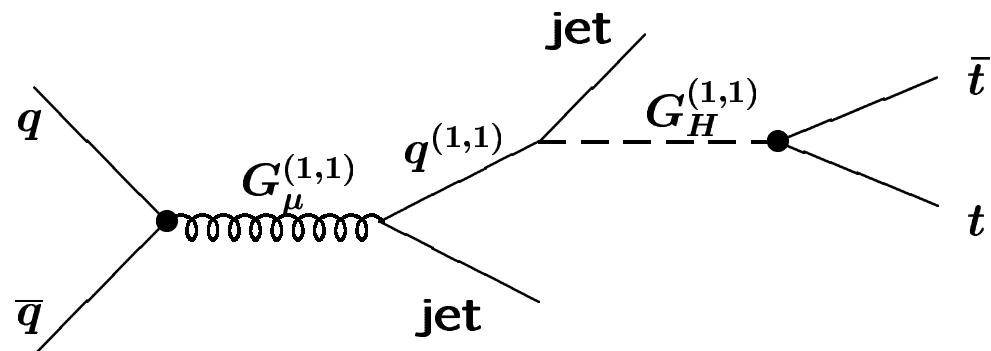
$$\Phi^{(j,k)}(x^\mu) \mapsto (-1)^{j+k} \Phi^{(j,k)}(x^\mu)$$

Signals at colliders:

Pair production of (1,0) modes: $3l + \cancel{E}_T$ (soft leptons)

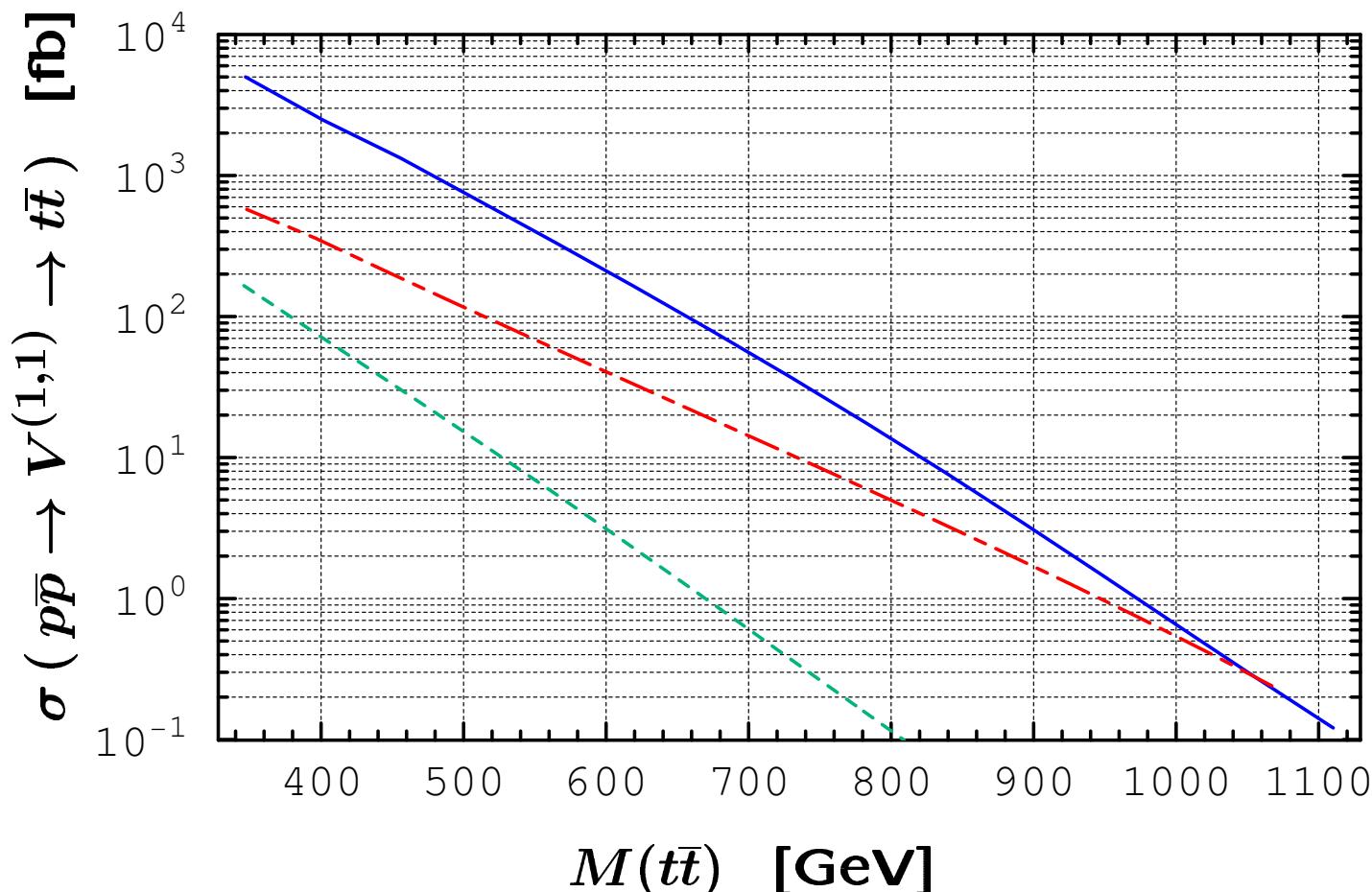


**s-channel production of a (1,1) mode of mass $\sqrt{2}/R$
→ $t\bar{t}$ resonance**

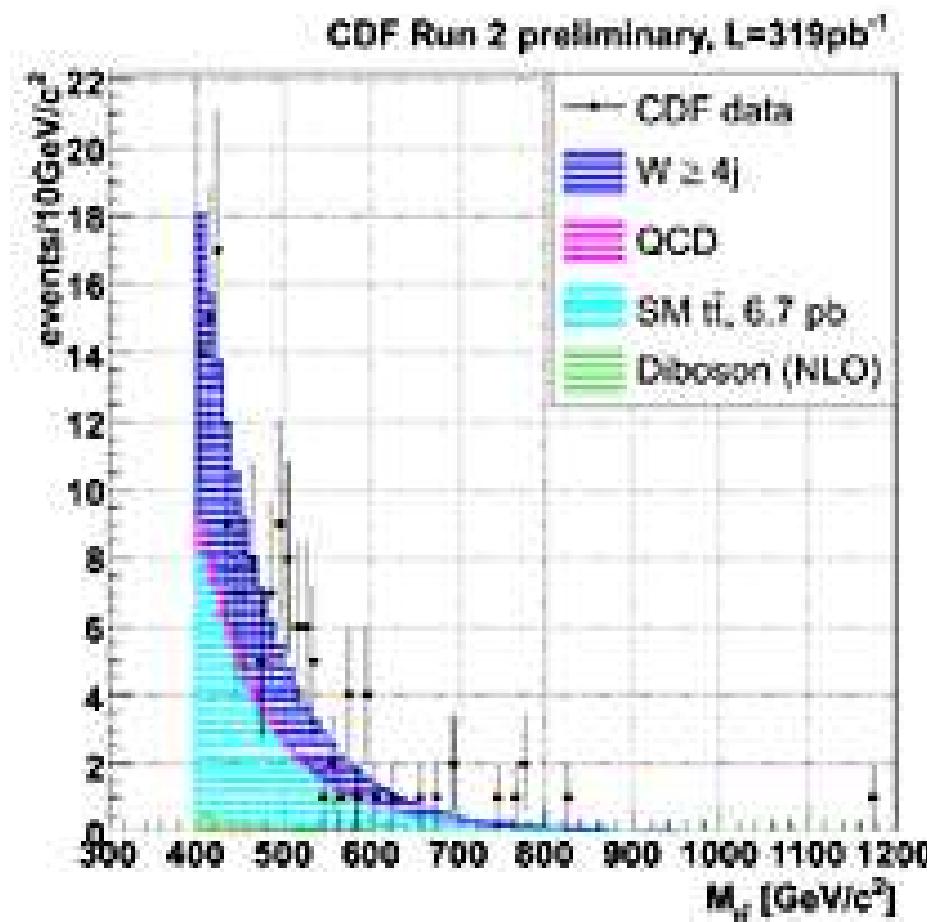
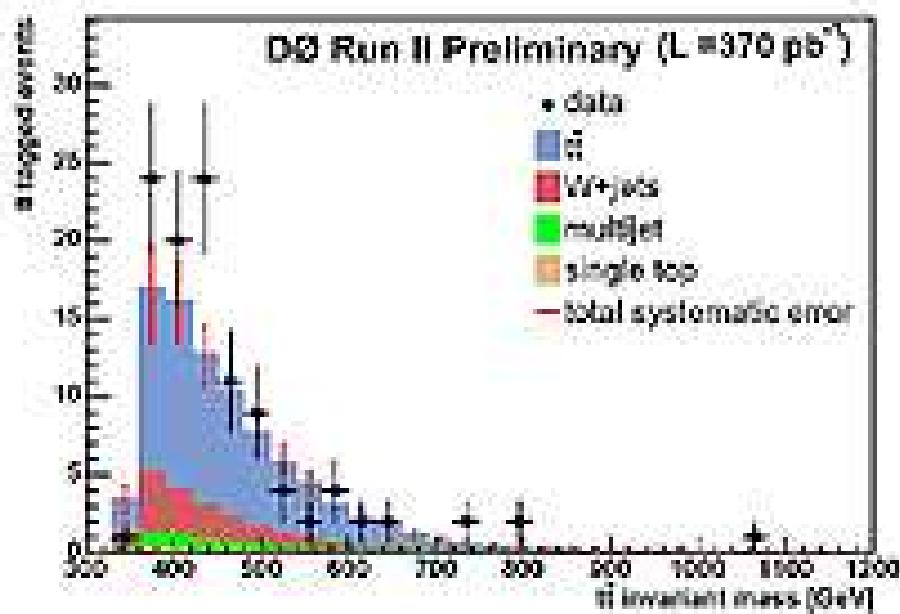


Production of $t\bar{t}$ pairs at the Tevatron from mass peaks at:

- $G_H^{(1,1)} + W_\mu^{(1,1)3}$ ————— $M_{t\bar{t}} \simeq 1.10 \sqrt{2}/R$
- $W_H^{(1,1)3} + B_\mu^{(1,1)}$ - - - - - $M_{t\bar{t}} \simeq 0.96 \sqrt{2}/R$
- $B_H^{(1,1)}$ - - - - - $M_{t\bar{t}} \simeq 0.87 \sqrt{2}/R$



D0 and CDF searches



Conclusions

- **Universal Extra Dimensions**

- compactification scale can be as low as ~ 300 GeV.
- lightest KK mode is a dark matter candidate

- **6-Dimensional Standard Model**

- 3 generations of quarks and leptons are required for global $SU(2)_W$ anomaly cancellation
- proton is long-lived due to 6D Lorentz invariance

- **Look for Kaluza-Klein modes at colliders:**

- 3 soft leptons + jets + E_T
- series of narrow $t\bar{t}$ resonances (or $\ell^+\ell^-$ resonances in the 5D case)